

BRIGHTEST CLUSTER GALAXIES AS STANDARD CANDLES

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Received 1994 June 7; accepted 1994 August 15

ABSTRACT

We investigate the use of brightest cluster galaxies (BCGs) as standard candles for measuring galaxy peculiar velocities on large scales. We have obtained precise large-format CCD surface photometry and redshifts for an all-sky, volume-limited ($z \leq 0.05$) sample of 119 BCG. We reinvestigate the Hoessel (1980) relationship between the metric luminosity, L_m , within the central $10 h^{-1}$ kpc of the BCGs and the logarithmic slope of the surface brightness profile, α . The L_m - α relationship reduces the cosmic scatter in L_m from 0.327 mag to 0.244 mag, yielding a typical distance accuracy of 17% per BCG. Residuals about the L_m - α relationship are independent of BCG luminosity, BCG $B-R_c$ color, BCG location within the host cluster, and richness of the host cluster. The metric luminosity is independent of cluster richness even before correcting for its dependence on α , which provides further evidence for the unique nature of the BCG luminosity function. Indeed, the BCG luminosity function, both before and after application of the α -correction, is consistent with a single Gaussian distribution. Half the BCGs in the sample show some evidence of small color gradients as a function of radius within their central $50 h^{-1}$ kpc regions but with almost equal numbers becoming redder as becoming bluer. However, with the central $10 h^{-1}$ kpc the colors are remarkably constant—the mean $B-R_c$ color is 1.51 with a dispersion of only 0.06 mag. The narrow photometric and color distributions of the BCGs, the lack of “second-parameter” effects, as well as the unique rich cluster environment of BCGs, argue that BCGs are the most homogeneous distance indicators presently available for large-scale structure research.

Subject headings: galaxies: clusters: general — galaxies: photometry

1. INTRODUCTION

We have recently completed a program to measure the motion of the Local Group with respect to a full-sky sample of Abell clusters (Lauer & Postman 1994, hereafter Paper I). Our goal was to explore galaxy bulk flows on much larger scales than had been done previously. Central to our investigation was the use of the luminosities of brightest cluster galaxies (BCGs) as distance indicators. BCGs had heretofore yielded conflicting results on large-scale peculiar motions (Sandage 1975; James, Joseph, & Collins 1987; Lucey & Carter 1988) because of heterogeneous data quality, small sample sizes, and the lack of full-sky coverage. However, we were impressed enough by the properties of BCGs to choose their use over the more conventional Tully-Fisher and D_n - σ methods. First, the BCGs as a class are probably more homogeneous than all other types of galaxies used as distance indicators. Second, the BCGs could be readily identified over the full sky from digitized sky surveys by looking in the vicinities of clusters in the Abell (1958) and Abell, Corwin, & Olowin (1989; hereafter ACO) catalogs out to distances well in excess of our $15,000 \text{ km s}^{-1}$ survey limit. Further, since the BCGs reside at the center of rich clusters, their individual peculiar velocities are minimized and their selection is largely invulnerable to inhomogeneous

Malmquist bias. Their photometric homogeneity also minimizes homogeneous Malmquist bias, as well as concerns over undiscovered biases due to possible “second parameters.” Last, use of BCGs is relatively easy; only photometric images and “redshift” quality spectra are required, in contrast to the Tully-Fisher and D_n - σ methods, which require higher signal-to-noise “dispersion” quality spectroscopic observations.

Humason, Mayall, & Sandage (1956) first identified the photometric homogeneity of BCGs, leading to the work of Sandage (1972a, b), who explored in detail the use of BCGs as cosmological probes. A key aspect of this work was a demonstration that the luminosities of the BCGs had an impressively small dispersion of 0.25 mag, after correction for a relationship between BCG luminosity and cluster morphology, and careful culling of the data. Gunn & Oke (1975) presented a somewhat different formulation of the BCG distance scale, focusing on the metric luminosity, L_m , of the BCGs within a physical aperture, r_m , of relatively small size compared to the extent of the BCGs. In this case, BCGs cannot be treated as point sources; errors in luminosity are related to errors in distance by α , the logarithmic slope of L_m as a function of r_m . Hoessel (1980) further showed the existence of a relationship between L_m and α that could be used to both reduce the luminosity scatter of BCGs as well as to counter some selection effects. Unfortunately, the L_m - α relationship measured from the more modern photometry of Hoessel & Schneider (1985) was poorly defined, showing little improvement over the scatter in L_m alone.

We have selected the 119 BCGs in the present sample to define an inertial frame suitable for measuring the motion of the Local Group (see Paper I for results). The sample covers all

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² Space Telescope Science Institute is operated by AURA, Inc., under contract to the National Aeronautics and Space Administration.

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the sky with $|b| \geq 15^\circ$ and is volume limited, in contrast to the previous BCG studies listed above, which were incomplete and largely based on galaxies observable from the north. Our program imposed stringent requirements on the photometric accuracy of our imaging data. Specifically, a 3σ detection of a 600 km s^{-1} bulk velocity on a scale of $15,000 \text{ km s}^{-1}$ demands that any total systematic error in galaxy magnitudes be limited to $\lesssim 0.02$ mag. A key component of the measurement of the Local Group motion presented in Paper I is, thus, the primary topic of this paper—a precise recalibration of the $z \approx 0$ photometric properties of the central regions of BCGs and a full assessment of their status as standard candles.

We find the metric luminosity, L_m , of the central $10 h^{-1} \text{ kpc}$ regions of BCGs within $z \leq 0.05$ to be independent of their $B-R_c$ color, their location within the host cluster, and the richness of the host cluster. We confirm the Hoessel (1980) L_m - α relationship, finding that the cosmic scatter in L_m is significantly reduced when its dependence on the logarithmic slope of the surface brightness profile, α , is accounted for. The residuals between the observed L_m and the metric luminosity predicted from the L_m - α relationship contain distance information and are shown to be independent of BGC luminosity, BCG color, and cluster parameters. These observations support the hypothesis that the central regions of BCGs are composed from quite similar stellar populations and, thus, validate their use as standard candles and as probes of the large-scale kinematics of the local universe.

The sample selection process is summarized in § 2. The data acquisition and reduction are described in § 3. Section 4 contains the full description of our BCG distance indicator and the various correlations between the photometric properties of the BCGs as well as between these properties and the properties of the host clusters. Our conclusions are in § 5.

2. SAMPLE SELECTION

2.1. The Abell Cluster Sample

We began by selecting for possible observation all Abell or ACO clusters (i.e., no richness class limits) with measured or photometrically estimated heliocentric redshifts $\leq 15,000 \text{ km s}^{-1}$ and galactic latitude $|b| \geq 15^\circ$. These criteria yielded a sample of 153 clusters. Of these 153 clusters, 22 ultimately turned out to have $z > 0.05$ —incorrectly measured redshifts were always cases where the cluster redshift was based on only one galaxy that we discovered to be in the foreground. Because the L_m - α distance indicator is strictly applicable to old stellar systems, we further restrict our observations to elliptical BCGs. Consequently, we exclude from the present analysis four of the clusters which have spiral BCGs (A2995, A3354, A3578, and A3816). In addition, A426 is excluded given the well-known A-type spectrum of its BCG, NGC 1275. There are six Abell and ACO “clusters” (A34, A256, A480, A762, A3388, and A3990) which were not observed. This is because there is no apparent overdensity we can identify on the POSS or SERC plates at the reported coordinates that is consistent with the indicated richness and estimated or previously measured redshift. A400 and A3555, while within the redshift cutoff, are not used in this analysis because they have anomalously faint ($\sim 4\sigma$ deviation) BCGs. A2295 and A3577 are also within the redshift cutoff, but the spectroscopic observations were unfortunately obtained after the conclusion of our imaging runs and hence no CCD image data are available.

We included three clusters (A3528, A3530, and A3532) from

the “apparent” background of the very rich Shapley concentration that are nominally outside the $15,000 \text{ km s}^{-1}$ redshift cutoff because they may indeed be within the volume, given plausible infall velocities around this supercluster. The Shapley supercluster is the only large enhancement in the cluster distribution which straddles our survey redshift limit (mean heliocentric redshift = $14,567 \text{ km s}^{-1}$). The next nearest large supercluster is a five-member system (mean heliocentric redshift = $16,670 \text{ km s}^{-1}$) located in the Leo supercluster complex (Tully 1987). This system is substantially less rich than the Shapley supercluster (which has 11 members at the same percolation length), and, consequently, the probable infall velocities are smaller and it is unlikely that any of the member clusters are really within the survey volume.

The final sample consists of 119 clusters and is given in Table 1. We emphasize that the exclusion of 32 of the 35 clusters discussed above is based solely on redshift, the lack of a significant overdensity, or nonelliptical BCG morphology. All cluster and BCG velocities given in Table 1 are heliocentric. The cluster space density is nearly constant out to $15,000 \text{ km s}^{-1}$, indicating that the sample is effectively volume limited (see Fig. 3 in Paper I). A Kolmogorov-Smirnov (K-S) test comparing the observed redshift distribution with that expected for a volume-limited survey yields a probability of 0.74 that the cluster sample is indeed volume limited. This is an especially attractive feature of the sample as distance-dependent biases are minimized or eliminated. The overdensity at $\sim 4500 \text{ km s}^{-1}$ is a $\lesssim 2\sigma$ fluctuation due to the Hydra-Centaurus supercluster. The prominence of this fluctuation is amplified because there are only 13 Abell and ACO clusters with $z \leq 0.02$.

2.2. The Brightest Cluster Galaxy Selection

We performed our own BCG selection for each cluster. This was necessary because accurate positions of BCGs in most of the southern ACO clusters were not available in the literature and because we wanted to eliminate scatter in the distance indicator due to galaxies incorrectly identified as BCGs by previous investigators. This latter consideration does in fact appear to be important, as in a number of cases we have selected different galaxies from those that Hoessel, Gunn, & Thuan (1980) did for the same clusters (see the notes to Table 1).

Our first pass at BCG selection was done from 14.5×14.5 images extracted from the STScI digitized Schmidt plate database, centered on the cluster positions given in the Abell and ACO catalogs. For clusters north of the celestial equator, the plate material came from the Palomar Quick V Survey (QV; epoch 1983). For clusters south of the celestial equator, the plate material came from the UK Schmidt SERC J Survey (SERC- J epoch 1975). Both surveys were digitized using the PDS microdensitometers at STScI during construction of the *Hubble Space Telescope* (HST) Guide Star Catalog (Lasker et al. 1990). The image pixel size in all cases is $1''.7$ ($25 \mu\text{m}$). The plate limits are $V \approx 19$ and $J \approx 21$ for the QV and SERC- J surveys, respectively, more than sufficient to assure complete BCG identification out to $z = 0.05$.

We selected BCG candidates by eye, but we also computed instrumental magnitudes of the galaxies from the plates to aid our decisions. We relied on redshift data (where available) as well to eliminate foreground galaxies. In cases where two BCG candidates in the same cluster could not be ranked unambiguously (primarily due to saturation effects), we obtained CCD data for both. There were no clusters in our sample with more than two BCG candidates.

TABLE 1
THE BCG SAMPLE

Abell	R.A. (J2000) Dec	BCG cz	Cluster cz	Runs	Notes	Abell	R.A. (J2000) Dec	BCG cz	Cluster cz	Runs	Notes		
76	00 39 26.32	+06 44 02.5	11162	11162	2	I1565	11 25 31.10	+35 30 15.3	10220	10297	4	Note 1	
119	00 56 16.13	-01 15 22.1	13320	13354	2, 3		11 28 36.40	+26 54 19.4	9798	9830	4		
147	01 08 37.90	+02 16 06.1	13138	13022	2	Note 1	11 33 05.09	-04 00 49.8	15508	15508	4		
160	01 12 59.76	+15 29 29.3	13095	13095	2		1314	11 34 49.42	+49 04 38.3	9977	9838	4	I 712
168	01 14 57.70	+00 25 52.1	13439	13446	2, 3	Note 1	1367	11 44 02.17	+19 56 58.7	6237	6469	4	N3842
189	01 23 26.33	+01 42 17.8	10236	9925	2		1631	12 53 18.20	-15 31 58.0	14275	14030	4	Note 4
193	01 25 07.54	+08 41 57.0	14825	14553	2	I1695	1644	12 57 11.76	-17 24 35.0	14237	14220	4, 5	
194	01 25 58.84	-01 20 21.8	5321	5357	2	N 545	1656	13 00 08.04	+27 58 36.3	6497	6961	4	N4889
195	01 26 54.75	+19 12 50.1	12796	12796	2, 3	I 115	1736	13 27 28.04	-27 19 27.6	13605	13583	4	I4252
260	01 50 42.96	+33 04 54.8	10699	10998	2		1836	14 01 41.84	-11 36 22.7	10999	11036	4, 5	
261	01 51 26.48	-02 15 11.5	13965	13965	3		1983	14 52 43.25	+16 54 12.9	13756	13617	4	Note 1
262	01 52 46.28	+36 09 05.3	4831	4913	2	N 705	2040	15 12 47.74	+07 26 02.0	13680	13698	4	
295	02 02 17.24	+01 07 40.6	12818	12890	2, 3		2052	15 16 44.55	+07 01 16.8	10331	10575	4	
347	02 25 26.47	+41 49 27.1	5257	5604	2	N 910	2063	15 23 05.46	+08 36 33.2	10626	10633	4, 5	
376	02 46 04.04	+36 54 16.7	14700	14700	2		2107	15 39 39.05	+21 46 55.7	12613	12496	4	U9958
397	02 56 28.91	+15 54 58.9	10286	9975	2		2147	16 02 17.04	+15 58 27.6	10384	10511	4, 5	
407	03 01 51.60	+35 50 27.4	14059	14130	2		2151	16 04 35.76	+17 43 18.0	11248	11115	5	N6041
419	03 08 28.73	-23 36 53.2	12240	12240	2	Note 2	2152	16 05 29.28	+16 26 11.0	13598	13598	4, 5	
496	04 33 37.68	-13 15 39.7	9893	9893	1, 2, 3, 4		2162	16 12 35.73	+29 29 03.5	9547	9629	5	N6086
533	05 01 08.25	-22 34 58.1	14365	14149	2, 3, 5		2197	16 29 45.04	+40 48 42.5	8800	9042	4	N6173, Note 1
539	05 16 55.12	+06 33 09.5	9682	8754	2, 3, 4		2199	16 28 38.45	+39 33 03.0	9348	9034	4	N6166
548	05 49 21.74	-25 20 47.0	11848	11949	1, 2, 5		2247	16 52 48.08	+81 37 57.9	12016	11641	4	
569	07 09 07.76	+48 36 55.3	5724	5749	2, 4	N2329	2572	23 17 13.54	+18 42 28.2	13009	12420	2	
576	07 22 06.83	+55 52 29.3	12072	11542	4		2589	23 23 57.55	+16 46 35.8	12144	12397	2	
634	08 15 44.90	+58 19 15.8	8135	8135	4		2593	23 24 20.15	+14 38 50.4	12514	12489	2	N7649
671	08 28 31.92	+30 25 48.0	14970	15219	4	I2378	2634	23 38 29.44	+27 01 49.7	9141	9153	2	N7720
779	09 19 46.80	+33 44 58.7	6867	6796	4	N2832	2657	23 44 30.49	+09 15 51.6	12454	12100	4	Note 1
912	10 01 09.41	-00 04 46.5	13572	13572	3, 4, 5		2666	23 50 58.66	+27 08 47.6	8123	8057	2	
957	10 13 38.27	-00 55 31.3	13438	13433	3, 4, 5		2717	00 03 12.77	-35 56 16.7	14936	14753	1	N 25
999	10 23 23.85	+12 50 05.8	9749	9603	4		2731	00 09 59.04	-57 01 17.7	9457	9457	1	N 215
1016	10 27 07.91	+11 00 36.9	9705	9669	4	I 613	2806	00 40 48.47	-56 12 53.4	8229	8252	1	N 215
1060	10 36 42.56	-27 31 42.9	3704	3719	5	N3311	2870	01 07 41.84	-46 54 31.2	6714	7233	1, 3	I1625
1100	10 48 45.73	+22 13 03.4	13990	13990	4		2877	01 09 54.82	-45 55 57.3	7261	7309	3	I1633
1139	10 58 11.02	+01 36 15.4	11460	11788	4	U6057, Note 3	2881	01 11 14.14	-17 04 14.5	13363	13324	2	
1142	11 00 45.39	+10 33 11.6	10118	10501	4	I 664	2896	01 18 18.60	-37 06 17.0	9555	9555	3	
1177	11 09 44.52	+21 45 31.5	9561	9561	4		2911	01 25 12.19	-38 05 38.0	5992	6056	3	N 544
1185	11 10 38.32	+28 46 03.0	10521	9917	4	N3550	3144	03 37 04.88	-55 01 20.8	13120	13546	3	
1213	11 16 22.87	+29 15 08.1	13535	14045	4	Note 1	3193	03 58 13.54	-52 19 44.4	10009	10350	1, 3	
1228	11 21 23.19	+34 21 24.4	10674	10973	4		3367	05 49 41.71	-24 32 43.7	13461	13461	2, 3	
							3374	05 56 42.97	-21 15 12.4	14201	14307	2, 3	

TABLE 1—Continued

Abell	R.A. (J2000) Dec		BCG cz	Cluster cz	Runs	Notes
3376	06 00 41.01	-40 02 46.6	13814	13907	1, 3	
3381	06 09 53.76	-33 35 34.1	11488	11467	1, 3	
3389	06 22 21.30	-64 56 04.6	8294	8114	1, 3	
3395	06 27 36.42	-54 27 01.7	14520	14712	3	
3526	12 48 48.94	-41 18 42.0	3045	3454	5	N4696
3528	12 54 22.31	-29 00 45.9	16425	16317	5	Note 5
3530	12 55 36.10	-30 20 49.3	16162	16206	5	Note 5
3532	12 57 22.15	-30 21 47.5	16633	16646	5	Note 5
3537	13 01 00.71	-32 26 28.9	5104	5053	5	
3542	13 08 41.52	-34 34 31.3	10389	10387	5	
3553	13 19 15.08	-37 10 46.5	14446	14446	5	
3554	13 19 31.53	-33 29 19.7	14333	14333	5	
3556	13 24 06.76	-31 40 15.0	14459	14500	5	
3558	13 27 56.53	-31 29 46.8	14110	14312	5	
3559	13 29 51.02	-29 30 53.1	14105	14213	5	
3560	13 31 53.33	-33 14 04.4	3644	3734	5	N5193
3562	13 33 34.74	-31 40 20.3	14708	14708	5	
3564	13 34 55.37	-35 05 57.8	14496	14721	5	
3565	13 36 39.06	-33 57 56.7	3762	3834	5	I4296
3566	13 39 38.22	-35 36 32.6	14529	14529	5	
3570	13 46 24.00	-37 58 15.5	11377	11156	5	
3571	13 47 28.42	-32 51 51.8	11679	11913	5	
3572	13 48 14.26	-33 22 57.8	12134	12141	5	
3574	13 49 05.29	-30 17 44.9	4523	4657	5	I4329
3575	13 52 38.36	-32 53 17.3	11188	11188	5	
3581	14 07 29.51	-27 01 07.0	6618	6682	5	I4374
3656	20 00 49.97	-38 34 35.9	5965	5768	3, 5	I4931
3676	20 24 24.50	-40 21 59.5	12108	12108	3	
3677	20 26 23.56	-33 21 03.5	13787	13789	3	
3698	20 35 56.21	-25 16 45.3	5799	6040	3	N6936
3716	20 51 56.82	-52 37 47.9	13900	13426	3, 5	
3733	21 01 59.07	-28 03 34.5	10943	11039	1, 3	
3736	21 05 04.49	-43 25 09.2	14604	14604	3	
3742	21 07 52.28	-47 10 43.4	4764	4842	3	N7014
3744	21 07 16.19	-25 28 08.4	11041	11153	3	
3747	21 08 38.93	-43 29 11.1	9174	9170	1, 3	
3869	22 20 31.05	-55 07 29.9	12005	12005	3	N7249
4038	23 47 28.30	-28 06 34.6	8218	8501	3	I5353
4049	23 51 36.65	-28 21 55.7	8137	8512	1	I5362
4059	23 57 00.37	-34 45 32.4	14696	14730	3	

NOTE.—All velocities are heliocentric and are given in km s^{-1} . (1) The present BCG choice differs from the galaxy selected by Hoessel et al. 1980. (2) A0419 appears to consist of two overlapping clusters at $z = 0.041$ and $z = 0.068$. We have selected the lower redshift component. (3) The redshift for the BCG in A1139 (obtained from the CfA Redshift Survey) had a relatively large measurement error of 170 km s^{-1} (see Table 6). However, a remeasurement of its redshift by us in 1992 gives a heliocentric velocity of $11,495 (\pm 29) \text{ km s}^{-1}$, within 35 km s^{-1} of the less precise value. (4) A1631 appears to consist of two overlapping clusters at $z = 0.014$ and $z = 0.046$. We have selected the higher redshift component. (5) This cluster was included to test for inhomogeneous velocity bias at the sample edge (see Paper I).

To back up our first pass at BCG selection, we inspected the QV and SERC-J Schmidt plate prints directly by eye out to a projected radius of $5 h^{-1} \text{ Mpc}$ ($\sim 2^\circ$ at $z = 0.05$) to identify candidate BCGs lying just beyond the initial $14''.5$ field boundaries. To facilitate this step, we created transparent overlays of the positions of all CfA Catalog redshifts within the projected area (Huchra 1991; Huchra et al. 1992). There are 18 clusters (14.5% of the sample) with BCGs that lie beyond our initial $14''.5$ field. In these cases, a second digitized image was extracted and the instrumental magnitude was computed to verify the selection.

Our final BCG identifications are strictly based on the CCD photometry discussed in the next section and consistency between the BCG redshift and the mean cluster redshift. We

emphasize that the designation of “brightest galaxy” is thus based only on the metric magnitudes presented below, rather than on estimated total magnitudes. Indeed, in a few cases, the identity of the BCGs would change if a much larger aperture size were used. Our strict definition of the BCGs also means that we considered all galaxies regardless of their location within a $5 h^{-1} \text{ Mpc}$ projected radius from the cluster center. In the Virgo Cluster, for example (which is not in our sample), we would have selected NGC 4472 over NGC 4486, despite the latter’s more central location and extensive X-ray halo. In our sample, the median projected BCG separation from the published center of its host cluster is $70 h^{-1} \text{ kpc}$, 90% of the BCGs lie within a projected radius of $350 h^{-1} \text{ kpc}$, and the largest projected BCG separation as $1.2 h^{-1} \text{ Mpc}$ (A548).

There are no significant differences in the BCG properties between clusters selected from the original Abell catalog and clusters selected from the ACO southern hemisphere extension. The metric luminosities, colors, and profile shapes are consistent with being drawn from the same universal distributions.

The positions of all BCG candidates were computed using the astrometric solutions provided by version 1.0 of the *HST* Guide Star Catalog. A center-of-gravity algorithm was used to compute the centroid in a $15'' \times 15''$ box centered on the peak pixel. The J2000 BCG positions are given in Table 1. All positions have absolute errors less than $1''.5$.

3. OBSERVATIONS AND REDUCTION

3.1. Photometry

3.1.1. Imaging Observations and Basic Reduction

CCD Images of the BCGs were acquired under photometric conditions using the KPNO 4 m and 2.1 m telescopes, and the CTIO 1.5 m telescope between 1989 November and 1991 April (over a total of five observing runs). Images were obtained in the Kron-Cousins R_c band, which we use as the primary bandpass of the photometric distance indicator. We also imaged all but 13 of the galaxies in the Johnson B band to provide color information. Table 2 summarizes the runs and CCD detectors used.

The exposure times were typically 200–600 s and were set by the desire that the R_c -band CCD images have signal-to-noise ratios (S/Ns) of at least 100 in the central pixels of the BCGs. While this number is somewhat arbitrary, high-S/N images allow the best identification and removal of galaxies or other objects that contaminate the BCG envelopes. The BCGs are quite red, and we accept a lower S/N (typically 40–50) for our B -band images, as these are not directly used to estimate cluster distances.

Dome flats and blank night-sky frames were obtained to flatten the images. Our R_c -band frames typically could be flattened to better than 0.5% of the sky level. Background sky levels were measured from the pixel intensity modes in the image corners. For the most extended BCGs, where we were concerned that galaxy light might contribute to the sky, we obtained images offset from the nominal pointing to sample the sky at larger angular distances from the galaxies. The offset frames overlapped with the BCG images, thus allowing direct measurement of the amount of galaxy-light contamination, independent of any global fluctuations of the sky that might have occurred between the two images.

3.1.2. Measurement of the Aperture Photometry

The BCG distance indicator is based on the integrated luminosity of the galaxy through an aperture of constant physical

TABLE 2
 BCG IMAGING RUNS

Run	Date	Telescope	CCD	Pixel Scale	FOV
1.....	1989 Oct	CTIO 1.5 m	TI 800 × 800 (TI2)	0".273	3.6
2.....	1989 Nov	KPNO 4 m	TI 800 × 800 (TI2)	0.299	4.0
3.....	1990 Nov	CTIO 1.4 m	TI 800 × 800 (TI2)	0.273	3.6
4.....	1991 Mar	KPNO 2.1 m	Tek 1024 × 1024 (TE1K)	0.304	5.2
5.....	1991 Apr	CTIO 1.5 m	Tek 1024 × 1024 (TEK1K-1)	0.434	7.4

size, with a correction based on the logarithmic slope, α , of the curve of growth at the edge of the aperture. Because we wanted to investigate the best aperture size to use and, further, were likely to vary the *angular* size of the aperture as we iteratively derived the Local Group motion, we represented the photometry of each BCG as a set of apparent magnitudes measured through apertures of increasing geometric size, rather than just the flux through any fiducial aperture. We then used a spline both to interpolate between the apertures and to measure the local α at any desired radius.

Measurement of the aperture curve of growth proceeded in two steps. The first step was to extract the surface brightness profile of the BCG. Since other cluster galaxies are often interacting with or projected against the envelope of the BCG, it was necessary to decompose such composite systems into their individual galaxies. We did this using the multi-isophote decomposition algorithm of Lauer (1986), which can solve for the brightness distributions of several overlapping galaxies in a simultaneous least-squares fit. Compact contaminating objects, such as stars, can also be excluded from the fit. The decomposition algorithm recovers the galaxy brightness distributions only under the assumption that the isophotes of each system are concentric and elliptical; no model of the brightness profile is assumed and isophote position angle is allowed to vary.

Implicit in this approach is identification of the BCG as the dominant component in a composite system. Since 75% of multiple systems are only chance superpositions of cluster galaxies, or result from high-speed nonmerging encounters (Lauer 1988), this procedure is likely to be sensible. We thus contrast our method with that of Schneider, Gunn & Hoessel (1983a), who view the whole system as the BCG, adopting the luminosity centroid as the center of the apertures.

The multi-isophote decomposition is especially important for ~ 30 of the 119 BCGs because the flux contribution from non-BCG galaxies is ≥ 0.02 mag within the metric aperture. The observed L_m - α relationship and the scatter about it do not change significantly if we exclude these 30 galaxies from our analysis. We note in passing that exclusion of these 30 BCGs also does not significantly affect the derived dipole solution given in Paper I. The results of this exclusion test provide further assurance that our photometric measurement procedure is not introducing artificial signals into the data.

The aperture photometry itself is measured from the model BCG light distribution reconstructed from the measured surface photometry. The R_c -band photometry and $B-R_c$ color (prior to correction for extinction and K -dimming) for each BCG is given in Table 3 as a function of geometrically increasing aperture radius in arcseconds. The second aperture given is always the adopted metric aperture, calculated for the cluster redshifts given in Table 1, under the assumption that the Local Group is at rest with respect to the Abell cluster inertial frame. Where multiple observations of the given BCG exist, they have

been averaged. Last, since the B -band images are shallower than the R_c images, for a number of galaxies we provide colors for only the inner apertures.

3.1.3. Photometric Calibration

The photometric quality of the skies, extinction coefficients, and photometric zero points were monitored by periodic observations of Landolt (1983) standard stars (typically 10–15 observations per night). The majority of the standard stars were selected to have broadband $V-R_c$ and $B-V$ colors similar to those of BCGs. However, some bluer stars were also observed in order to evaluate the significance of second-order extinction terms—these turned out to be unimportant. The quality of the atmospheric extinction transformations in R_c was always excellent, with the rms deviations of the standard stars from the mean extinction line being always less than 0.01 mag on photometric nights.

As a variety of CCDs and filter sets were used during the course of this program, we observed a sample of 32 galaxies in common between the various runs to verify that all observations were reduced to consistent photometric zero points as well as to determine the basic accuracy of the galaxy photometry. We also used the overlap observations to verify that the relative pixel scales given in Table 2 were correct, an important consideration, given that the photometry depends on aperture size. Specifically, 25 BCGs were observed on two separate runs; five BCGs were observed on three separate runs; one BCG (A496) was observed on four separate runs; and one galaxy (the second brightest member in A548) was observed on all five runs. The apparent R_c -band metric magnitudes for galaxies with repeated observations are presented in Table 4. The rms scatter for the R_c -band galaxy photometry at the adopted metric aperture (see below) is 0.020 mag, which implies a random error of 0.014 mag for any given galaxy. The B -band photometry is somewhat poorer and has a random error of 0.025 mag. In principle, the repeat observations could be used to measure any zero-point offsets that might be required to correct for systematic differences between the runs. We found that such corrections were not necessary, however; the average photometry from all runs agreed to 0.01 mag without any adjustment.

We have also compared our photometry to the combined photoelectric and CCD R_c photometry of Colless et al. (1993) as an external check for the 21 galaxies that we have in common. The comparisons are presented in Table 5, where we have interpolated our own photometry to the apertures given by Colless et al. A straight comparison between the two data sets shows us to be brighter by -0.010 ± 0.012 mag with an rms difference of 0.057 mag. A few galaxies in Table 5, however, have much larger disagreements than any differences seen in Table 4. If we eliminate the two galaxies with deviations over 0.1 mag, located in A76 and A2589, then the mean offset between us and Colless et al. becomes -0.026 ± 0.007 mag,

TABLE 3

BRIGHTEST CLUSTER GALAXY APERTURE PHOTOMETRY

Abell	R_c	$B - R_c$	Abell	R_c	$B - R_c$	Abell	R_c	$B - R_c$	Abell	R_c	$B - R_c$	Abell	R_c	$B - R_c$						
76	15.53	13.491	1.848	27.16	13.618	1.627	496	17.79	13.160	1.714	45.82	12.077	1.642	50.81	13.926	...				
19.41	13.356	1.845	33.95	13.552	1.622	22.24	12.964	1.702	22.24	12.964	1.702	57.27	11.985	1.638	63.51	13.707	...			
24.26	13.223	1.843	42.44	13.493	1.613	27.80	12.780	1.689	27.80	12.780	1.689	71.59	11.900	1.633	957	13.45	13.544	1.742		
30.33	13.095	1.842	17.11	13.652	1.685	34.75	12.603	1.753	34.75	12.603	1.753	576	15.15	13.982	1.750	16.81	13.359	1.738		
37.91	12.974	1.844	21.39	13.523	1.681	43.44	12.441	1.757	43.44	12.441	1.757	18.94	13.903	1.739	21.02	13.178	1.735			
47.39	12.866	1.850	26.74	13.396	1.678	54.30	12.291	1.770	54.30	12.291	1.770	23.68	13.842	1.724	26.27	13.003	1.739			
119	13.21	13.711	1.729	15.25	13.860	1.765	33.42	13.268	1.676	67.87	12.132	1.806	29.60	13.799	1.704	32.84	12.841	1.745		
16.51	13.524	1.722	19.06	13.691	1.752	41.78	13.143	1.675	84.84	11.937	1.918	634	21.02	13.034	1.717	41.05	12.688	1.755		
20.64	13.342	1.717	29.78	13.367	1.719	52.22	12.999	1.700	533	12.76	14.042	1.660	21.02	13.034	1.717	51.31	12.543	1.791		
25.80	13.167	1.740	37.23	13.224	1.708	347	29.36	12.318	1.663	15.94	13.918	1.648	26.28	12.909	1.714	64.14	12.406	...		
32.25	13.006	1.737	194	31.11	12.045	1.609	36.70	12.173	1.668	19.93	13.802	1.635	32.85	12.797	1.712	80.18	12.278	...		
40.32	12.864	1.737	38.89	11.890	1.618	45.88	12.029	1.675	24.91	13.697	1.623	41.06	12.694	1.712	100.22	12.137	...			
50.40	12.772	1.713	48.61	11.752	1.631	57.35	11.893	1.689	31.14	13.603	1.609	51.32	12.594	1.712	125.27	11.977	...			
62.99	12.679	...	60.76	11.630	1.649	71.68	11.762	1.707	38.93	13.488	1.616	64.15	12.499	...	156.59	11.803	...			
147	13.51	13.767	1.668	75.96	11.519	1.677	48.66	13.383	1.757	48.66	13.383	1.757	80.19	12.404	...	999	18.42	13.344	1.690	
16.89	13.660	1.664	94.94	11.415	1.719	376	12.07	14.188	1.809	60.83	13.211	...	671	11.86	13.772	1.781	23.02	13.236	1.688	
21.11	13.562	1.661	15.08	14.017	1.805	15.08	14.017	1.805	76.03	13.078	...	76.03	13.078	...	28.78	13.131	1.684			
26.39	13.474	1.660	18.86	13.853	1.803	95.04	12.936	...	95.04	12.936	...	14.82	13.599	1.775	35.97	13.024	1.681			
32.99	13.395	1.661	23.57	13.698	1.802	539	19.91	13.295	1.754	19.91	13.295	1.754	23.16	13.261	1.770	44.97	12.921	1.682		
41.24	13.324	1.664	29.46	13.553	1.806	24.89	13.171	1.743	24.89	13.171	1.743	28.95	13.103	1.771	56.21	12.823	1.695			
160	13.40	14.250	1.765	33.41	13.492	1.705	31.11	13.051	1.737	31.11	13.051	1.737	36.18	12.956	1.778	70.26	12.728	...		
16.75	14.064	1.759	41.77	13.417	1.696	397	17.33	13.468	1.796	38.89	12.933	1.730	45.23	12.821	1.786	1016	18.31	13.531	1.665	
20.94	13.890	1.756	21.67	13.326	1.796	21.67	13.326	1.796	27.08	13.187	1.797	56.54	12.701	...	22.89	13.424	1.666			
26.18	13.734	1.753	27.08	13.187	1.797	33.85	13.049	1.803	60.77	12.699	1.736	70.67	12.607	...	28.61	13.323	1.667			
32.72	13.595	1.749	33.85	13.049	1.803	42.32	12.920	1.807	42.32	12.920	1.807	49.59	11.572	1.655	35.77	13.227	1.673			
40.90	13.470	1.751	52.90	12.804	1.810	52.90	12.804	1.810	52.90	12.804	1.810	61.99	11.438	1.659	44.71	13.136	1.688			
168	13.13	13.876	1.713	38.33	12.777	1.716	397	17.33	13.468	1.796	15.00	13.568	1.665	779	25.39	11.996	1.650			
16.42	13.733	1.711	47.91	12.651	1.723	407	12.52	14.495	1.890	18.75	13.446	1.657	31.74	11.850	1.649	39.67	11.709	1.650		
20.52	13.601	1.709	15.65	14.280	1.902	15.65	14.280	1.902	29.29	13.220	1.632	23.43	13.330	1.647	61.99	11.438	1.659			
25.65	13.478	1.708	19.62	13.181	1.717	33.85	13.049	1.803	45.77	13.036	1.655	45.77	13.036	1.655	77.48	11.311	1.657			
32.06	13.364	1.708	24.53	13.043	1.714	42.32	12.920	1.807	57.21	12.900	1.686	57.21	12.900	1.686	96.86	11.193	...			
40.08	13.259	1.712	30.66	12.907	1.714	30.66	12.907	1.714	71.51	12.778	...	71.51	12.778	...	912	13.32	14.409	1.702		
50.10	13.200	1.690	38.33	12.777	1.716	47.91	12.651	1.723	89.39	12.664	...	89.39	12.664	...	16.65	14.304	1.696			
62.63	13.139	1.686	47.91	12.651	1.723	15.65	14.280	1.902	111.74	12.543	...	111.74	12.543	...	20.81	14.209	1.696			
189	17.38	13.778	1.634	19.87	13.646	1.691	24.45	13.894	1.933	139.67	12.421	...	139.67	12.421	...	26.01	14.123	1.701		
21.73	13.694	1.631	31.04	13.405	1.701	30.56	13.702	1.944	419	14.48	14.320	1.615	419	14.48	14.320	1.615	32.52	14.046	1.752	
			38.81	13.295	1.718	38.81	13.295	1.718	18.11	14.238	1.614	22.63	14.162	1.615	26.01	14.123	1.701	40.65	13.986	...
			31.04	13.405	1.701	28.29	14.093	1.620	28.29	14.093	1.620	35.36	14.041	1.635	40.65	13.986	...			
			38.81	13.295	1.718	41.43	12.055	1.708	41.43	12.055	1.708									

TABLE 3—Continued

Abell	Rad	R_c	$B - R_c$	Abell	Rad	R_c	$B - R_c$	Abell	Rad	R_c	$B - R_c$	Abell	Rad	R_c	$B - R_c$
1139	16.07	13.895	1.703	1213	12.77	13.951	1.674	1367	26.84	12.149	1.624	2063	16.49	13.522	1.659
	20.09	13.766	1.700		15.97	13.821	1.666		33.56	12.011	1.621		20.61	13.321	1.650
	25.11	13.645	1.700		19.96	13.709	1.657		41.95	11.893	1.616		25.77	13.140	1.641
	31.39	13.530	1.700		24.95	13.618	1.651		52.43	11.784	1.612		32.21	12.977	1.631
	39.23	13.423	1.697		31.19	13.539	1.646		65.54	11.680	1.609		40.26	12.829	1.613
	49.04	13.322	1.695		38.99	13.458	1.636		81.92	11.586	1.605		50.32	12.691	...
					48.73	13.373	1.624	1367	102.41	11.513	...		62.91	12.558	...
	15.20	13.751	1.738									2063	16.49	13.522	1.659
	19.00	13.608	1.741	1228	16.05	13.717	1.636		33.56	12.011	1.621		20.61	13.321	1.650
	23.75	13.472	1.746		20.07	13.610	1.632		41.95	11.893	1.616		25.77	13.140	1.641
	29.69	13.339	1.755		25.08	13.509	1.625		52.43	11.784	1.612		32.21	12.977	1.631
	37.11	13.207	1.770		31.35	13.415	1.620		65.54	11.680	1.609		40.26	12.829	1.613
	46.39	13.083	1.798		39.19	13.326	1.615		81.92	11.586	1.605		50.32	12.691	...
	57.99	12.961	1.848		48.99	13.241	1.617		102.41	11.513	...		62.91	12.558	...
1142	16.91	13.463	1.663	1257	17.03	13.890	1.625	1631	12.91	13.841	1.724	2107	14.11	13.478	1.756
	21.14	13.332	1.657		21.29	13.826	1.612		16.14	13.677	1.717		17.64	13.286	1.750
	26.43	13.199	1.653		26.62	13.774	1.596		20.17	13.527	1.710		22.37	12.916	...
	33.03	13.062	1.647		33.27	13.727	1.575		25.22	13.390	1.696		27.97	12.794	...
	41.29	12.925	1.645		41.59	13.681	1.543		31.52	13.262	1.678		34.96	12.677	...
	51.62	12.801	1.643	1267	17.87	13.472	1.631		39.40	13.140	...		43.70	12.565	...
	64.52	12.699	1.645		22.34	13.386	1.627		49.25	13.025	...		54.62	12.464	...
					27.92	13.318	1.626		61.56	12.914	...		68.28	12.371	...
1177	18.42	13.132	1.647	1308	11.76	13.922	1.743	1644	12.75	13.887	1.792	2147	16.58	13.475	1.713
	23.03	12.954	1.643		14.70	13.788	1.734		15.94	13.642	1.785		20.72	13.306	1.711
	28.78	12.783	1.635		18.37	13.654	1.725		19.92	13.412	1.780		25.90	13.154	1.707
	35.98	12.623	1.628		22.97	13.522	1.716		24.91	13.203	1.771		32.38	13.014	1.703
	44.98	12.474	1.621		28.71	13.398	1.702		29.13	13.013	1.761		40.48	12.874	1.697
	56.22	12.331	1.614		35.89	13.288	1.683		31.13	13.013	1.761		50.60	12.731	1.692
	70.27	12.191	...		44.86	13.192	...		38.92	12.840	1.755		63.24	12.582	1.695
	87.84	12.057	...		56.08	13.107	...		48.64	12.682	1.751		79.06	12.361	...
	109.80	11.927	...		70.09	13.019	...		60.81	12.537	...		98.82	12.169	...
1185	17.71	13.210	1.594		84.86	12.940	...	1656	76.01	12.397	...		123.52	11.951	...
	22.14	13.049	1.598		95.01	12.261	...		95.01	12.261	...		154.41	11.704	...
	27.68	12.916	...		118.76	12.110	...		148.45	11.948	...		18.92	12.777	...
	34.60	12.795	...		148.45	11.948	...						23.65	12.582	...
	43.24	12.679	...										29.56	12.401	...
	54.06	12.566	...	1314	17.66	13.155	1.674		24.74	11.863	1.638	2151	15.72	13.591	...
	67.57	12.452	...		22.08	13.015	1.668		30.92	11.712	1.632		19.66	13.406	...
									38.66	11.572	1.618				
									48.32	11.441	...				

TABLE 3—Continued

Abell	R _c	B - R _c	Abell	R _c	B - R _c	Abell	R _c	B - R _c	Abell	R _c	B - R _c	Abell	R _c	B - R _c		
46.18	12.076	...	2657	14.38	14.219	...	73.28	11.750	...	16.75	14.242	1.658	47.27	13.244	...	
57.73	11.927	...	17.98	14.129	...	2877	23.76	11.771	1.640	20.94	14.154	1.668	3526	52.60	10.590	1.726
72.16	11.784	...	22.47	14.053	...	29.70	11.622	1.638	37.12	11.480	1.638	3374	12.66	14.409	1.611	
90.20	11.648	...	28.09	13.991	...	46.40	11.342	1.640	58.01	11.214	1.644	3376	13.04	13.846	1.736	
112.75	11.523	...	35.11	13.935	...	72.51	11.097	...	19.77	14.288	1.621	3381	15.64	13.987	1.693	
2247	14.88	13.860	...	2666	20.82	12.559	1.623	13.31	14.389	1.626	24.72	14.255	1.630	34.31	13.245	...
18.60	13.756	...	26.02	12.421	1.620	16.63	14.287	1.624	30.90	14.226	1.626	3376	13.04	13.846	1.736	
23.25	13.662	...	32.53	12.295	1.619	20.79	14.199	1.628	38.62	14.205	...	3376	13.04	13.846	1.736	
29.06	13.582	...	40.66	12.187	1.622	2881	13.31	14.389	1.626	24.72	14.255	1.630	34.31	13.245	...	
36.33	13.510	...	50.82	12.094	1.627	16.63	14.287	1.624	30.90	14.226	1.626	3376	13.04	13.846	1.736	
45.41	13.435	...	63.53	12.010	1.637	20.79	14.199	1.628	38.62	14.205	...	3376	13.04	13.846	1.736	
56.77	13.358	...	2717	12.16	14.283	1.680	25.99	14.124	1.629	16.30	13.695	1.727	27.45	13.530	...	
2572	14.00	13.662	1.868	15.20	14.065	1.658	18.31	13.368	1.623	16.30	13.695	1.727	27.45	13.530	...	
17.50	13.536	1.871	19.00	13.842	1.631	22.88	13.279	1.621	20.37	13.552	1.718	20.37	13.552	1.718		
21.88	13.403	1.862	23.75	13.620	1.600	28.60	13.199	1.624	25.46	13.416	1.710	25.46	13.416	1.710		
27.35	13.260	1.846	29.69	13.411	1.559	31.83	13.286	1.632	31.83	13.286	1.705	31.83	13.286	1.705		
34.19	13.107	1.821	37.11	13.221	...	35.75	13.126	1.632	39.79	13.170	1.736	39.79	13.170	1.736		
42.73	12.956	1.810	46.39	13.046	...	44.69	13.060	1.649	49.73	13.069	1.696	49.73	13.069	1.696		
2589	14.03	13.844	1.587	57.99	12.882	...	29.11	28.40	12.872	1.636	62.17	12.978	...	130.89	12.250	...
17.54	13.653	1.587	2731	18.60	13.049	1.660	35.50	12.788	1.636	16.30	13.695	1.727	27.45	13.530	...	
21.93	13.466	1.586	23.25	12.914	1.663	44.38	12.719	1.641	55.47	12.665	1.654	3381	15.64	13.987	1.693	
27.41	13.286	1.589	29.07	12.781	1.670	55.47	12.665	1.654	19.55	13.849	1.705	3381	15.64	13.987	1.693	
34.26	13.118	1.589	36.33	12.653	1.684	3144	13.32	14.200	1.666	24.44	13.710	1.726	24.44	13.710	1.726	
42.83	12.966	1.584	45.42	12.536	1.707	20.81	13.985	1.669	30.55	13.578	1.765	30.55	13.578	1.765		
2593	13.95	13.833	1.689	56.77	12.431	1.723	26.01	13.888	1.676	38.18	13.448	1.818	38.18	13.448	1.818	
17.43	13.626	1.687	2806	21.22	12.870	1.644	32.52	13.797	1.691	47.73	13.231	...	47.73	13.231	...	
21.79	13.445	1.683	26.52	12.766	1.657	40.65	13.721	1.720	59.66	13.095	...	59.66	13.095	...		
27.24	13.286	1.681	33.15	12.678	1.678	50.81	13.659	...	3389	21.96	12.620	1.657	67.45	12.740	...	
34.05	13.147	1.687	41.44	12.600	1.713	17.20	13.351	1.624	27.45	12.492	1.654	34.32	12.375	1.678		
2634	18.50	12.902	1.713	51.80	12.526	1.757	3193	17.20	13.351	1.624	42.89	12.273	1.704	42.89	12.273	1.704
23.12	12.741	1.710	64.75	12.458	...	21.50	13.228	1.612	21.50	13.228	1.612	3395	12.39	14.276	1.841	
28.91	12.586	1.710	2870	24.01	12.424	1.580	26.87	13.118	1.599	33.59	13.015	1.583	15.49	14.079	1.829	
36.13	12.433	1.712	30.02	12.269	1.575	41.98	12.917	1.640	19.36	13.896	1.823	15.49	14.079	1.829		
45.16	12.284	1.720	37.52	12.122	1.570	52.48	12.845	1.649	24.20	13.720	1.821	19.36	13.896	1.823		
56.46	12.139	1.731	46.90	11.988	1.566	58.62	11.870	1.578	30.26	13.552	1.826	24.20	13.720	1.821		
			58.62	11.870	1.578	3367	13.40	14.343	1.653	37.82	13.393	1.840	30.26	13.552	1.826	
									37.82	13.393	1.840	37.82	13.393	1.840		

TABLE 3—Continued

Abell	Rad	R_c	$B - R_c$	Abell	Rad	R_c	$B - R_c$	Abell	Rad	R_c	$B - R_c$	Abell	Rad	R_c	$B - R_c$	
19.59	13.341	1.737	47.18	13.069	1.660	23.56	12.823	1.722	125.73	11.548	...	29.68	12.971	1.658		
24.49	13.214	1.720	58.98	12.902	...	29.45	12.592	1.713	157.16	11.367	...	37.10	12.873	1.685		
30.61	13.095	1.701	73.72	12.736	...	36.82	12.392	1.700	3656	29.60	11.919	1.620	46.37	12.784	1.684	
38.27	12.990	1.681	92.16	12.570	...	46.02	12.218	1.680	37.00	11.771	1.606	79.44	12.739	...		
47.83	12.907	1.664	115.19	12.401	...	57.53	12.062	1.658	46.25	11.621	1.589	99.30	12.624	...		
59.79	12.838	...	143.99	12.233	...	71.91	11.913	1.634	89.88	11.471	1.572	3869	14.80	13.699	1.673	
3558	12.69	13.562	1.752	3564	12.36	14.231	1.717	112.35	11.605	...	3733	15.84	14.005	1.871		
15.86	13.339	1.739	15.46	14.118	1.709	89.88	11.762	...	57.81	11.471	1.572	18.51	13.604	1.671		
19.83	13.130	1.726	19.32	14.020	1.697	140.44	11.436	...	72.26	11.327	1.560	23.13	13.517	1.670		
24.79	12.937	1.714	24.15	13.936	1.682	175.55	11.252	...	90.32	11.179	...	28.91	13.438	1.677		
30.98	12.754	1.702	30.19	13.870	1.659	112.91	11.025	...	36.14	13.371	1.688		
38.73	12.585	1.688	37.73	13.826	1.627	14.81	13.795	1.698	141.13	10.865	...	45.18	13.317	1.708		
48.41	12.429	1.665	47.17	13.798	1.584	18.51	13.707	1.680	176.42	10.557	...	56.47	13.270	1.733		
60.51	12.283	1.639	58.96	13.780	1.532	23.14	13.627	1.658	3676	14.61	13.623	1.576	4038	20.27	12.980	1.593
75.04	12.144	1.606	3572	14.81	13.795	1.698	18.51	13.707	1.680	18.27	13.504	1.564	25.34	12.867	1.582	
94.55	12.014	...	3574	38.03	11.470	1.642	23.14	13.627	1.658	18.27	13.504	1.564	31.67	12.763	1.574	
118.19	11.883	...	3575	47.54	11.289	1.631	28.93	13.557	1.631	22.83	13.394	1.557	39.59	12.669	1.570	
3559	12.77	13.628	1.779	91.09	10.490	1.438	36.16	13.498	1.599	28.54	13.299	1.548	49.49	12.590	1.566	
15.96	13.477	1.782	113.86	10.370	1.435	45.20	13.444	1.550	35.68	13.221	1.538	30.01	12.992	1.699		
19.95	13.326	1.789	142.33	10.247	...	56.50	13.392	1.491	44.59	13.169	1.527	37.51	12.872	1.698		
24.94	13.182	1.797	3565	46.64	10.871	1.449	38.03	11.470	1.642	12.92	14.498	1.748	46.88	12.765	1.696	
38.97	12.920	1.816	3566	58.30	10.740	1.446	47.54	11.289	1.631	16.15	14.398	1.738	58.61	12.667	1.696	
48.71	12.811	1.836	3570	72.87	10.614	1.442	74.28	10.944	1.608	20.19	14.298	1.727	73.26	12.577	1.710	
60.89	12.716	1.876	3575	92.85	10.799	1.589	92.85	10.799	1.589	25.24	14.196	1.712	3742	35.28	12.046	1.597
76.11	12.617	1.949	3576	116.07	10.664	1.569	116.07	10.664	1.569	31.55	14.088	1.698	44.10	11.942	1.592	
95.13	12.511	...	3577	145.08	10.529	1.544	145.08	10.529	1.544	39.43	13.972	1.684	55.12	11.841	1.590	
3560	47.96	11.465	1.603	181.35	10.399	...	181.35	10.399	...	49.29	13.831	1.709	68.90	11.746	1.590	
59.94	11.384	1.602	3578	16.00	14.423	1.633	38.03	11.470	1.642	27.96	12.512	1.687	4059	12.17	13.716	1.728
74.93	11.307	1.601	3579	20.01	14.354	1.597	47.54	11.289	1.631	34.96	12.420	1.687	15.22	13.496	1.723	
93.66	11.227	...	3580	25.01	14.285	1.541	59.43	11.110	1.621	3744	15.67	13.478	1.798	19.02	13.287	1.721
3562	12.37	14.056	1.763	31.26	14.215	1.454	74.28	10.944	1.608	43.69	12.332	1.688	23.77	13.096	1.717	
15.46	13.883	1.752	3581	39.07	14.142	...	92.85	10.799	1.589	43.69	12.332	1.688	29.72	12.920	1.717	
19.33	13.716	1.739	3582	39.07	14.142	...	116.07	10.664	1.569	54.62	12.248	1.690	37.15	12.761	1.719	
24.16	13.550	1.724	3583	26.37	12.617	1.690	145.08	10.529	1.544	39.07	14.142	...	46.43	12.619	1.728	
30.20	13.387	1.705	3584	32.96	12.465	1.681	181.35	10.399	...	38.25	13.097	1.901	58.04	12.485	1.755	
37.75	13.228	1.683	3585	41.20	12.320	1.678	181.35	10.399	...	47.82	13.012	1.977	72.55	12.349	1.807	
			3586	51.50	12.183	1.678	100.58	11.784	...	32.54	13.185	1.669				
			3587	64.37	12.052	1.678	80.46	11.922	1.687	26.03	13.329	1.672				
			3588	80.46	11.922	1.687	100.58	11.784	...	32.54	13.185	1.669				
			3589	100.58	11.784	...										

TABLE 4
INTERRUN PHOTOMETRY COMPARISONS

Abell	Run 1	Run 2	Run 3	Run 4	Run 5
119-1.....	...	13.528	13.521
168-1.....	...	13.734	13.732
195-1.....	...	13.809	13.763
295-1.....	...	13.681	13.622
496-1.....	12.952	12.981	12.968	12.955	...
533-1.....	...	13.936	13.944	...	13.897
539-1.....	...	13.194	13.139	13.180	...
548-1.....	13.418	13.467	13.449
548-2.....	13.485	13.470	13.472	13.455	13.492
569-1.....	...	12.171	...	12.186	...
912-1.....	14.312	14.313	14.296
957-1.....	13.344	13.369	13.361
1631-2.....	11.691	11.695
1644-1.....	13.644	13.640
1836-1.....	13.141	13.146
2063-1.....	13.317	13.325
2147-1.....	13.327	13.285
2152-1.....	13.825	13.868
2657-2.....	...	14.217	...	14.186	...
2870-1.....	12.286	...	12.253
2881-2.....	14.314	14.295
2911-2.....	12.773	...	12.837
3193-1.....	13.248	...	13.209
3367-1.....	...	14.259	14.226
3374-1.....	...	14.327	14.341
3376-1.....	13.684	...	13.706
3381-1.....	13.817	...	13.881
3389-1.....	12.508	...	12.476
3656-1.....	11.753	11.780	...
3716-1.....	13.647	13.668	...
3733-1.....	13.843	...	13.857
3747-1.....	13.075	...	13.074
3747-2.....	13.472	...	13.459

NOTE.—Comparisons are shown for the R_c -band photometry at the adopted metric aperture in the Local Group frame. Details of the runs are given in Table 2. A portion of the comparisons were based on the second-ranked galaxies in some clusters.

TABLE 5
EXTERNAL PHOTOMETRY COMPARISONS

Abell	Radius	R_c	R_w	$R_w - R_c$
76.....	14.95	13.379	13.511	0.132
119.....	14.95	13.595	13.604	0.009
168.....	14.95	13.801	13.790	-0.011
260.....	14.95	13.340	13.354	0.014
262.....	14.95	12.981	12.926	-0.055
496.....	14.95	13.305	13.318	0.013
548.....	14.95	13.659	13.568	-0.091
1656.....	14.95	12.272	12.239	-0.033
1983.....	19.75	13.931	13.853	-0.078
2040.....	14.95	14.217	14.193	-0.024
2052.....	14.95	13.483	13.456	-0.027
2107.....	14.95	13.448	13.427	-0.021
2147.....	19.75	13.384	13.340	-0.044
2151.....	14.95	13.663	13.633	-0.030
2162.....	19.75	13.030	12.984	-0.046
2197.....	19.75	12.483	12.457	-0.026
2247.....	14.95	13.895	13.858	-0.037
2589.....	14.95	13.637	13.784	0.147
2593.....	14.95	13.740	13.760	0.020
2666.....	9.60	13.107	13.097	-0.010
2657.....	14.95	14.214	14.202	-0.012

NOTE.—Comparisons are shown for the R_c -band photoelectric photometry of Colless et al. 1993 (R_c) and the present work (R_w), for the selected Colless et al. aperture.

with an rms difference of 0.030 mag per galaxy, which nicely matches the quadrature sum of the Colless et al. stated internal error of 0.025 mag, and our own of 0.014 mag. In either case, however, the comparison with Colless et al. suggests that there may be a slight zero-point mismatch between the two sets. We have emphasized internal consistency in our photometry, but have invested much less effort in tying our observations to the true photoelectric R_c system. We thus present our photometry with the caveat that it may contain a slight offset from other external sets of R_c galaxy observations.

3.2. Redshifts

The majority of the redshift data used here are obtained from the literature. However, new redshifts were obtained for 33 ACO clusters (and their BCGs) by us at CTIO and for 17 northern BCGs by J. Huchra (1992) and A. Zabludoff (1992) at the Multiple Mirror Telescope (MMT). For the southern clusters, we used the CTIO 1.5 m telescope and the GEC-CCD spectrograph with grating No. 09 (8.4 Å resolution) over the spectral range 3900–6300 Å. Redshifts were determined primarily from the Mg I and Na D absorption lines using the Tonry-Davis (Tonry & Davis 1979) cross-correlation technique. The galaxies NGC 1316, NGC 1427, NGC 3115, NGC 6868, and NGC 7507 were used as templates. The typical error for the CTIO redshifts is 60 km s⁻¹. A similar reduction procedure was used on the MMT spectrograph data. The MMT redshifts have a typical uncertainty of 30 km s⁻¹. We also observed 10 galaxies (from CTIO) for which independent redshift data already exist to estimate any systematic errors. We find a typical scatter of ±60 km s⁻¹ between our CTIO redshifts and the published values, consistent with expectations based on the internal errors. There is no significant velocity offset. Table 6 presents all the new heliocentric redshifts obtained during the course of our survey and, therefore, includes galaxies which have $z > 0.05$ as well. We note that the new BCG redshift data are also included in Table 6; thus there is some redundancy with Table 1. The source codes in Table 6 for data obtained by us at CTIO are “2” and “3”; source code “1” refers to data obtained by Huchra and Zabludoff at the MMT.

At the minimum, we have velocities for at least four galaxies per cluster including the BCG, and 51% of the clusters have measurements for 10 or more member galaxies. Fortunately, BCGs are generally close to the kinematic centers of the clusters (Quintana & Lawrie 1982; Zabludoff, Huchra, & Geller 1990; Zabludoff et al. 1993), so even where we have few velocities, because we include the BCGs, the error in the mean cluster redshift is small. BCGs, in general, are not at rest with respect to their cluster, but deviate from the mean velocity with a typical dispersion of σ_1 , which can be compared to the typically larger one-dimensional velocity dispersion for all cluster members, σ_c . For the 42 clusters with more than 20 members, we find $\sigma_c = 666$ km s⁻¹, in excellent agreement with Zabludoff et al. (1990). For the same clusters, we find $\sigma_1 = 264$ km s⁻¹ after correcting for the typical error in the mean cluster redshift of 102 km s⁻¹. The distribution of BCG-cluster velocity differences appears to be Gaussian, so we argue that clusters with “significant” BCG peculiar velocities simply represent the more extreme examples of a general phenomenon. From these results we conclude that we should give the BCGs a weight of $(\sigma_c/\sigma_1)^2 \approx 6.0$ compared to other cluster members when calculating mean cluster redshifts.

TABLE 6
GALAXY REDSHIFT OBSERVATIONS

Abell	R.A. (J2000) Decl.		cz^a	σ	Source	Notes
147-G1	01 ^h 08 ^m 37.90	+02°16'06".1	13138	33	1	
260-GA	01 50 42.96	+33 04 54.8	10699	28	1	
GB	01 51 23.57	+33 01 51.9	10411	26	1	
295-G1	02 02 17.24	-01 07 40.6	12818	29	1	
G2	02 02 20.17	-01 06 37.2	12996	31	1	
419-G1	03 08 15.84	-23 41 29.0	20391	34	1	1
484-G1	04 16 23.19	-07 40 59.0	20914	42	1	2
G2	04 16 47.68	-07 40 32.3	20714	31	1	2
533-G2	05 01 35.99	-22 36 03.1	14642	34	1	
912-G6	10 01 10.83	-00 03 46.8	27828	84	2	1
1100-G1	10 48 45.73	+22 13 03.4	13990	45	1	
1139-G1	10 58 11.02	+01 36 15.4	11460	170	1	
1644-G1	12 57 11.76	-17 24 35.0	14237	42	1	
G2	12 57 49.30	-17 32 44.5	14013	50	1	
1837-G1	14 01 36.37	-11 07 44.0	20722	3	1	2
2040-G1	15 12 47.74	+07 26 02.0	13680	30	1	
2107-G1	15 39 39.05	+21 46 55.7	12613	31	1	
2148-G1	16 03 19.79	+25 27 14.3	26606	82	1	2
2806-G1	00 40 12.92	-56 09 15.0	8244	49	3	
G3	00 40 03.53	-56 10 53.3	8478	48	3	
2860-G1	01 04 03.85	-39 46 59.2	32224	63	3	2
2870-G2	01 07 44.43	-46 51 48.6	7571	39	3	
G3	01 07 14.69	-46 50 20.6	9317	46	3	
2882-G1	01 11 22.08	-17 04 16.7	13661	32	1	Spiral
G2	01 10 54.46	-17 11 51.6	12665	38	1	
G3	01 11 14.14	-17 04 14.5	13363	43	1	
2933-G1	01 40 59.38	-54 37 27.8	27894	57	3	2
G2	01 40 35.13	-54 30 52.8	27355	65	3	2
2995-G3	02 14 55.44	-24 51 18.3	11050	56	3	
3202-G1	04 01 00.55	-53 41 08.6	21133	52	3	2
G2	03 59 29.13	-53 38 22.6	21063	75	3	2
3367-G1	05 49 41.71	-24 32 43.7	13461	34	1	
G2	05 48 50.27	-24 21 06.3	19984	31	1	1
3374-G1	05 56 42.97	-21 15 12.4	14201	71	1	
3381-G1	06 09 53.76	-33 35 34.1	11488	57	3	
G2	06 09 48.68	-33 35 51.4	11310	56	3	
3389-G1	06 22 21.30	-64 56 04.6	8294	48	3	
G2	06 21 26.41	-64 59 37.1	8027	44	3	
G3	06 21 24.42	-64 57 45.4	8391	45	3	
3528-G1	12 54 22.31	-29 00 45.9	16425	78	2	
G4	12 54 24.91	-28 58 24.6	14754	102	2	
3537-G3	13 01 28.49	-32 20 03.2	4949	39	2	
G4	13 01 25.84	-32 31 08.5	25647	79	2	1
3545-G1	13 11 21.65	-34 04 48.2	29172	73	2	2
G3	13 11 14.52	-34 05 36.1	15029	93	2	2
3554-G1	13 19 31.53	-33 29 19.7	14333	78	2	
G4	13 19 19.76	-33 28 06.0	14605	82	2	
3557-G1	13 24 55.16	-28 53 15.8	23405	61	2	2
3570-G1	13 46 47.33	-37 54 28.4	11237	74	2	
G2	13 46 24.00	-37 58 15.5	11377	61	2	
G4	13 46 47.15	-37 54 41.9	11310	86	2	
3577-G1	13 54 14.83	-27 50 51.8	14848	83	2	
3581-G1	14 07 29.51	-27 01 07.0	6618	39	2	I4374
G2	14 07 44.13	-27 04 58.8	6595	41	2	
3653-G1	19 53 02.85	-52 02 14.8	32717	97	3	2
3656-G1	20 00 49.97	-38 34 35.9	4965	34	3	I4374
G3	20 00 09.79	-38 30 59.5	5487	39	3	
G4	20 00 25.38	-38 27 55.6	5888	34	3	
G5	20 00 00.33	-38 30 17.9	5902	32	3	
3664-G1	20 13 59.00	-80 39 08.7	41166	117	3	2
3676-GB	20 24 19.70	-40 21 39.2	12372	56	3	
3677-G1	20 26 23.56	-33 21 03.5	13787	60	3	
G2	20 26 13.18	-33 21 38.5	11930	65	3	
3695-G1	20 34 45.40	-35 49 25.9	26851	86	3	2
3698-G1	20 35 56.21	-25 16 45.3	5799	35	3	NGC 6936
G2	20 36 22.07	-25 20 30.1	5732	57	3	
G3	20 36 07.13	-25 11 54.1	6456	64	3	
3706-G1	20 42 14.56	-38 19 57.2	30763	66	3	2
3733-G1	21 01 59.07	-28 03 34.5	10943	70	3	
G3	21 02 23.82	-28 10 18.2	10816	50	3	Emission lines
G4	21 01 36.54	-28 03 33.3	12193	64	3	
3742-G1	21 06 47.32	-47 11 15.4	5163	30	3	

TABLE 6—Continued

Abell	R.A. (J2000) Decl.		cz^a	σ	Source	Notes
3744-G1	21 07 16.19	-25 28 08.4	11041	56	3	
G2	21 07 20.93	-25 29 15.3	12780	135	3	
G3	21 07 19.80	-25 29 17.1	10364	69	3	
G4	21 07 24.96	-25 25 46.8	11628	63	3	
G5	21 07 25.95	-25 25 48.5	11456	60	3	
3747-G2	21 08 28.99	-43 29 24.5	5086	30	3	
G4	21 08 41.23	-43 30 41.1	8967	49	3	
3816-G1	21 50 16.38	-55 16 35.5	11681	50	3	Emission lines
G2	21 50 44.25	-55 14 55.1	11906	69	3	
G3	21 49 58.70	-55 15 56.0	20234	54	3	1
3869-G1	22 20 45.73	-55 06 08.9	11920	61	3	
GW	22 20 45.14	-55 06 08.9	12407	75	3	
3879-G1	22 27 48.85	-69 01 24.1	19956	52	3	2
3911-G2	22 46 15.76	-52 43 42.0	29198	64	3	2

NOTES.—(1) This galaxy is a background object. (2) This galaxy is in a cluster beyond the 15,000 km s⁻¹ heliocentric velocity limit. The redshift measurement, however, was obtained by us to better determine the membership of the $z \leq 0.05$ cluster sample.

^a All velocities are heliocentric and are given in km s⁻¹.

SOURCES.—(1) MMT, J. Huchra & A. Zabludoff, private communication. (2) CTIO 1.5 m, 1991 Spring (3) CTIO 1.5 m, 1990 Fall.

A separate problem in calculating cluster redshifts is to identify the galaxies properly belonging to the cluster. Our approach is to first select candidate cluster members as those galaxies projected to be within the cylinder 1.5 h^{-1} Mpc in radius and ± 3000 km s⁻¹ in depth centered on the BCG location, and then to reject outliers by demanding that the velocity distribution be Gaussian. The rejection process is allowed to proceed until the probability that the resultant distribution is Gaussian exceeds 15%. The redshift error for any given cluster will always be 264 km s⁻¹ (σ_1) or less depending how many cluster members contribute to the average. The mean error in the mean cluster redshift for the entire sample is 184 km s⁻¹.

4. THE BCG DISTANCE INDICATOR

4.1. The Luminosity- α Relationship

As we have discussed in the introduction, the BCG distance indicator works by using

$$\alpha \equiv d \log L_m / d \log r |_{r_m} .$$

to predict L_m , where L_m is the total BCG luminosity within the circular aperture of physical radius r_m centered on the BCG nucleus. Effective use of the L_m - α relationship depends on two issues that we consider in this section, namely, (1) determination of the best relationship between L_m and α , and (2) selection of the best value of r_m . For convenience we calculate L_m using $H_0 = 80$ km s⁻¹ Mpc⁻¹. This choice is motivated by analysis of secondary distance indicators applied to the Virgo Cluster BCG, NGC 4472 (Lauer & Postman 1992), although the present analysis and our results are independent of the value adopted for H_0 . We have adopted the appropriate frame as that having the Abell clusters at rest on average (the F solution of Paper I). We have applied K -corrections which were computed by convolving the CCD + filter response with the spectral energy distribution of a typical BCG (Whitford 1971) converted to photons. For the B and R_c filters, the following analytic expressions provide excellent estimates for $z \leq 0.50$:

$$K_R = 2.5 \log_{10} (1 + 0.96z) ,$$

$$K_B = 2.5 \log_{10} (1 + 4.00z + 22.4z^2) .$$

The nominal corrections for Galactic extinction for the B and R_c filters are

$$A_R = 2.35E(B-V) ,$$

$$A_B = 4.05E(B-V) ,$$

where $E(B-V)$ values are obtained from Burstein & Heiles (1982) and are tabulated in Paper I. Last, we calculate r_m from angle θ_m (in arcseconds) as

$$r_m = 20.07 h^{-1} \theta_m [1 - (1+z)^{-1/2}] / (1+z) \text{ kpc} , \quad (1)$$

which assumes $q_0 = 0.5$ (although this choice matters little at the low redshifts of the present sample).

The relationship between L_m and α for our BCG sample is shown in Figure 1, using the optimal aperture of 10 h^{-1} kpc adopted below (the data plotted are published in Paper I). The relationship looks qualitatively similar to that of Hoessel

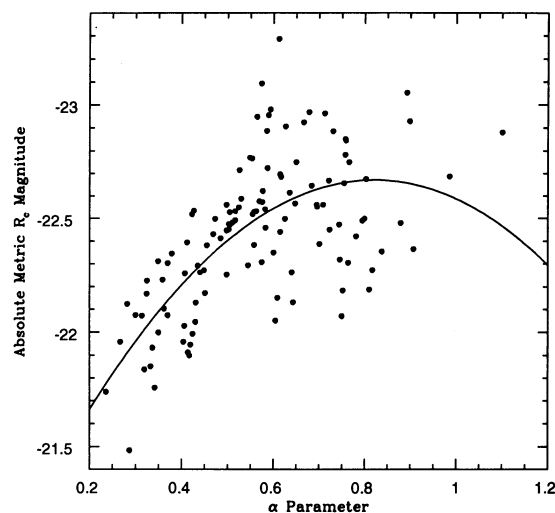


FIG. 1.—BCG R_c absolute metric luminosities are shown as a function of α , the logarithmic slope of the surface brightness profile evaluated at the metric radius. $H_0 = 80$ km s⁻¹ Mpc⁻¹ has been assumed. The solid line shows the mean L_m - α relationship. The relationship is shown after correction for the motion of the Local Group with respect to the Abell cluster inertial frame.

(1980). For small α the relationship appears to have a constant linear slope, but it flattens out toward larger α ; a parabolic fit appears adequate to match this nonlinearity and produces significantly smaller residuals than does a purely linear fit. The scatter about the best-fit parabola is $\sigma_m = 0.244$ mag, as compared to the scatter in L_m alone of $\sigma_m = 0.327$ mag. The form of the L_m - α relationship for $H_0 = 80 \text{ km s}^{-1} \text{ Mpc}^{-1}$ is

$$L_m = -20.896 - 4.397\alpha + 2.738\alpha^2 \quad (R_c \text{ band}). \quad (2)$$

with a total uncertainty of about 0.02 mag. Since the L_m - α relationship describes the flux of a BCG within a constant physical aperture, we can recast the relationship equally well as a predictor of surface brightness within the metric aperture. For $\alpha = 0.57$ (the mean value for this sample), the average surface brightness within r_m is $\mu_m = 20.787R_c \text{ mag arcsec}^{-2}$.

The residuals about the mean metric luminosity for the entire sample ($n = 0$ order polynomial fit) and the residuals about the quadratic L_m - α relationship ($n = 2$) are Gaussian. Fits to a Gaussian distribution yield probabilities of 0.248 ($\chi^2_v = 1.35$) and 0.916 ($\chi^2_v = 0.24$), respectively, that the residuals are normally distributed. Further, a K-S test appropriate for comparing an observed distribution with the distribution expected for a Gaussian with the same mean and standard deviation (Press et al. 1986, pp. 472–475) shows that the residuals of the L_m - α relationship are consistent with having been drawn from a Gaussian distribution at a 26.3% confidence level (rejection of the Gaussian hypothesis would require confidence levels of 5% or less). The histograms of the residuals about the mean metric luminosity for the entire sample ($n = 0$ order polynomial fit) and the residuals about the quadratic L_m - α relationship ($n = 2$) are shown in Figure 2a. The cumulative distribution functions (CDFs) for the residuals

about the $n = 0$ and $m = 2$ order polynomial fits to the L_m - α relationship are shown in Figure 2b, along with the corresponding CDFs for Gaussians with $\sigma_m = 0.244$ mag and $\sigma_m = 0.327$ mag. The aperture magnitudes alone, prior to the α -correction, are also Gaussian, but of course with larger scatter. The excellent agreement with a Gaussian distribution both before and after accounting for the L_m - α relationship argues strongly against the bimodal BCG luminosity function proposed by Bhavsar (1989). We suspect that the origin of the BCG luminosity distribution seen by Bhavsar is a combination of BCG misclassification by Hoessel and, in some cases, poor sky subtraction due to the small format of the CCD used by Hoessel. These effects can introduce a low-luminosity tail to the BCG luminosity distribution, perhaps causing it to appear bimodal.

Figure 3 shows the relative reduced χ^2 (normalized by the reduced χ^2 for the quadratic fit) and the scatter about the best-fit polynomial relation to the L_m - α relationship as functions of polynomial order, n . Goodness of fit is not significantly improved for $n > 2$. The minor reduction in χ^2_v for $n > 2$, which does occur, appears to be due to better fitting of the few BCGs with extreme α , where the relationship is poorly defined. We are currently extending the BCG sample to higher redshifts and thus should fill in the endpoints of the relationship much more densely.

The choice of metric aperture size is driven by the following considerations: (1) the aperture should be large enough to avoid seeing effects that could affect the photometry at the 0.5% level, (2) the aperture should be small enough to minimize sky-subtraction errors and to avoid encompassing the outer regions of the BCGs which can vary dramatically in surface brightness depending on whether or not the BCG is also a cD

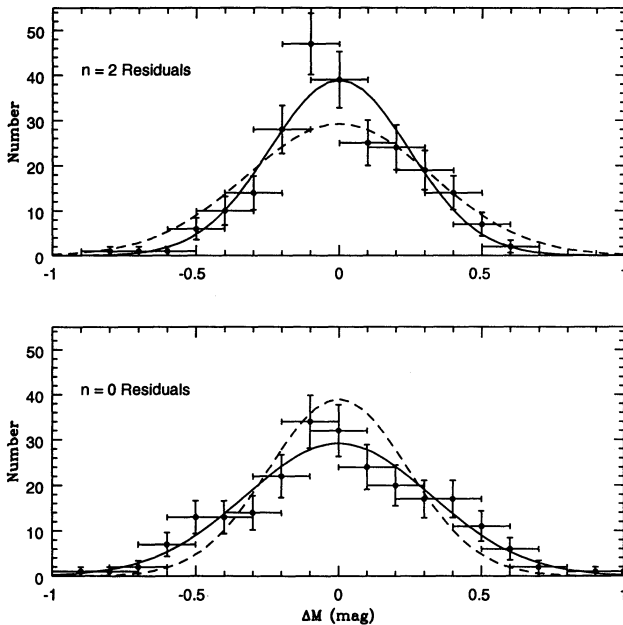


FIG. 2a

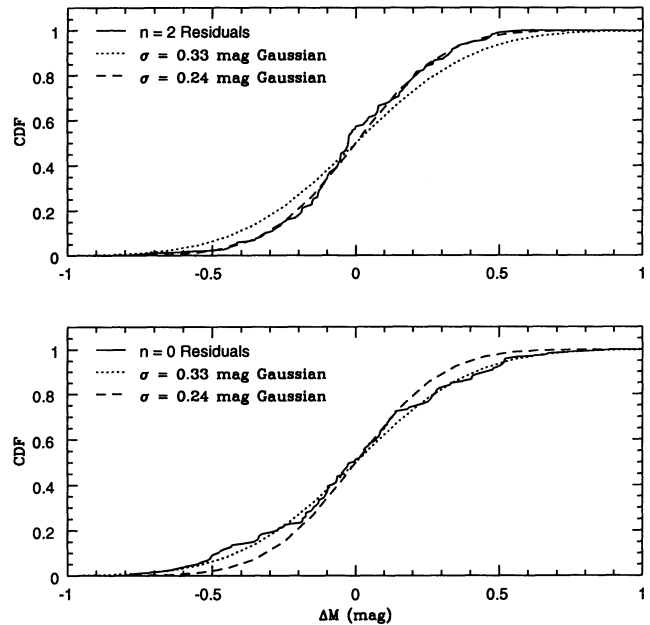


FIG. 2b

FIG. 2.—(a) Histograms of the residuals about the mean metric luminosity for the entire sample ($n = 0$) and the residuals about the quadratic L_m - α relationship ($n = 2$). The best-fit Gaussian distributions for each are superposed. (b) CDFs for the residuals about the mean metric luminosity for the entire sample ($n = 0$) and the residuals about the quadratic L_m - α relationship ($n = 2$). The CDFs for Gaussians with $\sigma_m = 0.244$ mag and $\sigma_m = 0.327$ mag are also shown. A K-S test between each residual CDF and the CDF for the best-fit Gaussian distribution yields a rejection of the hypothesis that the residual distributions are different from Gaussian at the 26% confidence level.

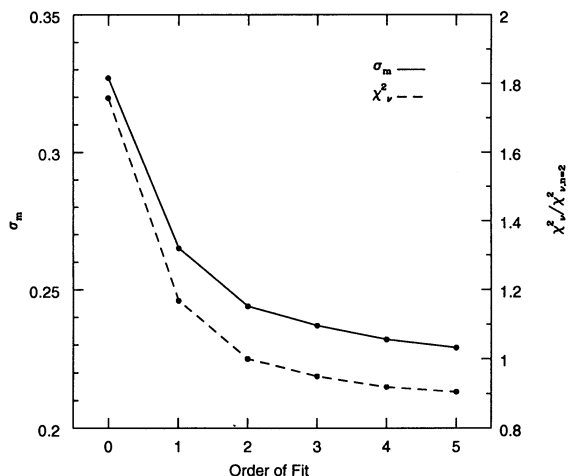


FIG. 3.—Scatter, σ_m , about the L_m - α relationship (solid curve) and the normalized χ^2_v value (divided by the χ^2_v for the quadratic fit) as functions of the polynomial order used to fit the relationship.

galaxy, and (3) within these two limits, the aperture should be selected to minimize the scatter in the BCG metric magnitudes. Figure 4 shows the scatter in the BCG metric luminosity as a function of aperture radius. The data for this plot are generated by placing apertures ranging in radius from $5 h^{-1}$ kpc to $20 h^{-1}$ kpc in intervals of 0.5 kpc on each BCG in our sample. At $z = 0.05$, this translates to an angular aperture range running from $7''.50$ to $30''.0$. The curve shows the scatter when cluster redshifts (and thus apertures) are corrected for the best-fit Local Group motion given in Paper I. The scatter is relatively constant for aperture radii less than $10 h^{-1}$ kpc but increases for larger radii. We thus chose $10 h^{-1}$ kpc to minimize any seeing-dependent effects and at the same time maximize distance accuracy. For comparison, Hoessel (1980) used a $9.6 h^{-1}$ kpc aperture. The metric magnitudes at the survey redshift limit change by only ~ 0.002 mag, independent of α , when the images are convolved with a typical stellar point-spread function. The effects of seeing on the measured photometry at $r_m = 10 h^{-1}$ kpc are, thus, negligible, which is not surprising, since the metric aperture radius is $14''.99$ at the survey limit of $z = 0.05$.

4.2. BCG Properties and Parameter Correlations

The central regions of BCGs interior to the metric aperture appear to contain similar stellar populations. Not only are the absolute metric magnitudes constant to within $\sim 24\%$ (after correcting for the L_m - α relationship), but the $B-R_c$ colors are impressively constant—the average $B-R_c$ color corrected for

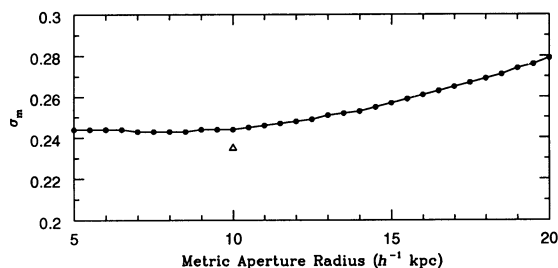


FIG. 4.—Scatter about the quadratic L_m - α relationship as a function of metric aperture radius.

extinction and K -dimming is 1.507 ± 0.006 mag, with a dispersion of only 0.055 mag, after the random errors in the photometry are accounted for. This is in excellent agreement with Schneider et al. (1983a), who found BCGs to have a dispersion of only 0.055 mag in $g-r$. The K -dimming and extinction corrections are indeed important, however; trend lines in the Figures 5a and 5b show how the color would redden if extinction and K -dimming corrections had not been made. Figures 5a-5d show the $B-R_c$ color as functions of $E(B-V)$, redshift, L_m , and the residual from the quadratic L_m - α relation, respectively. Once corrected, the colors show no dependence on extinction or redshift, nor on BCG metric luminosity.

The median color gradient, $d(B-R_c)/d \log r$, over the range $8-50 h^{-1}$ kpc is -0.031 with a scatter of 0.192 . There is no systematic trend for the BCGs to become bluer or redder with radius. The most extreme BCG color gradient seen in our sample is for A3554, which has $\Delta(B-R_c) = 0.36$ between radii of $40 h^{-1}$ kpc and $10 h^{-1}$ kpc.

BCG metric luminosities and residuals from the best-fit quadratic L_m - α relationship are also independent of the BCG location within its host cluster. Figures 6a and 6b show, respectively, the residuals as functions of BCG radial velocity offset and projected separation from the cluster center.

If BCGs just represent the brightest members drawn from a normal cluster luminosity function, then one might expect the BCG luminosity to depend on cluster richness (Scott 1957). However, the constancy of the BCG metric luminosity over an order of magnitude in cluster galaxy surface density argues against this hypothesis, as does the analysis of Tremaine & Richstone (1977). Figures 7a-7d show cluster richness (as represented by the Abell/ACO galaxy count) as functions of the BCG $B-R_c$ color, α -parameter, absolute R_m magnitude, and the residual from the quadratic L_m - α relation, respectively. Again, there are no significant trends between richness and any of these parameters. Note that the metric magnitudes in Figure 7c are *not* corrected for the L_m - α relationship, demonstrating that BCG metric luminosity is simply not dependent on cluster richness (alternatively, the Abell galaxy count may not be a good indicator of richness, but for $z \leq 0.05$ this is not likely). Of course, our sample contains only 12 clusters with Abell richness class greater than 1, and, hence, our constraints on the dependence of BCG parameters on richness are really limited to clusters with $R \leq 1$.

Finally, we look for correlations with BCG ellipticity, which is measured from the luminosity-weighted average of isophotes falling within the $10 h^{-1}$ kpc metric radius (cf. Ryden, Lauer, & Postman 1992). Figures 8a-8d show BCG ellipticity as functions of BCG $B-R_c$ color, α , L_m , and the residual from the quadratic L_m - α relation, respectively. The ellipticity is weakly correlated with both the metric luminosity and α (but since there exists an L_m - α relationship, this dual correlation is expected). The linear correlation coefficients are 0.306 and -0.241 , respectively, for the e - α relation and the e - L_m relation. The probabilities that a random sample of 119 points would produce correlation coefficients as large as this are 7.2×10^{-4} and 8.4×10^{-3} , respectively. Adding ellipticity as an independent variable does not significantly improve the dispersion in metric luminosity, however (with ellipticity, $\sigma_m = 0.242$ mag). Consequently, we expect and find that the residuals from the quadratic L_m - α relation are not significantly correlated with ellipticity.

In summary, residuals from the quadratic L_m - α relationship are independent of BCG color, BCG location, BCG ellipticity,

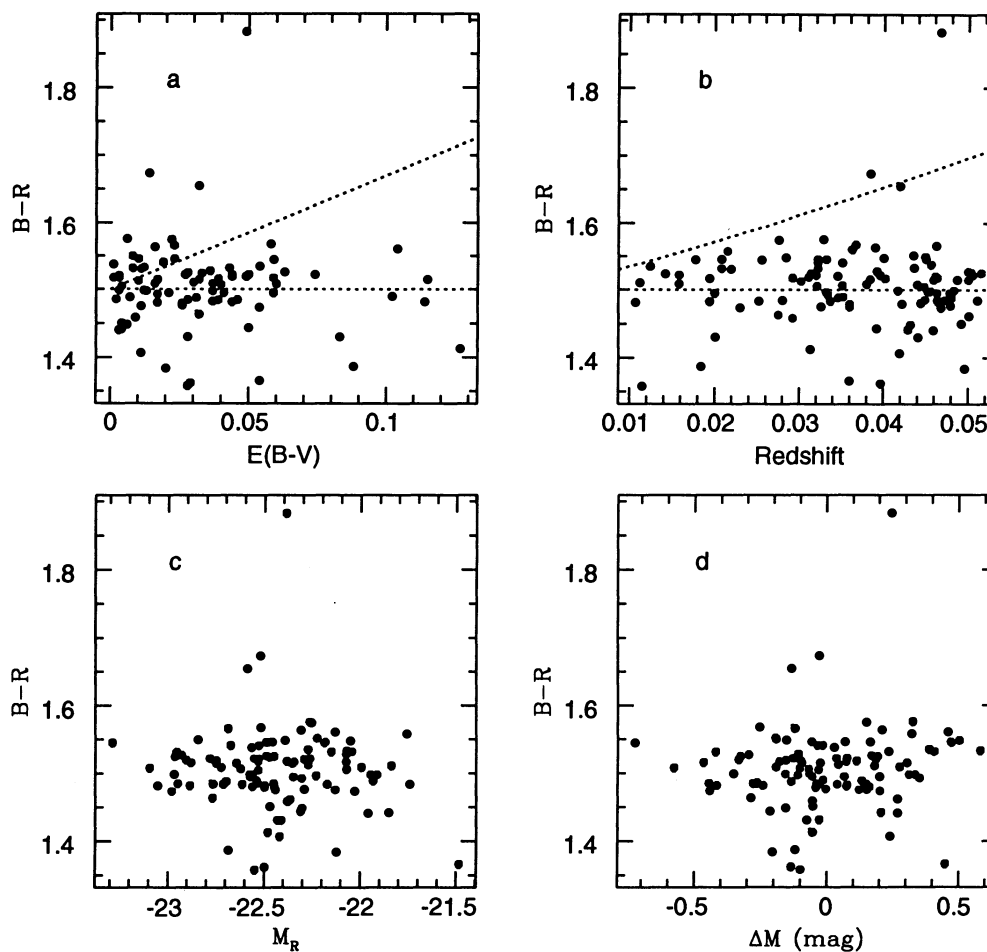


FIG. 5.—(a) BCG $B-R_c$ color as a function of $E(B-V)$. The rising dashed line shows the expected reddening when no extinction correction is applied. (b) BCG $B-R_c$ color as a function of redshift. The rising dashed line shows the expected reddening when no K -dimming correction is applied. (c) BCG $B-R_c$ color as a function of absolute R_m magnitude. (d) BCG $B-R_c$ color as a function of the residual from the quadratic $L_m-\alpha$ relationship.

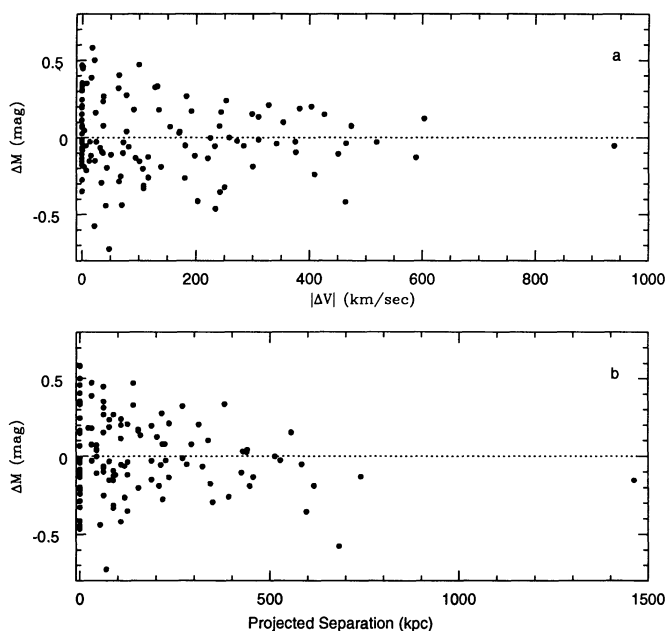


FIG. 6.—(a) Residual from the quadratic $L_m-\alpha$ relationship as a function of the BCG radial velocity offset from the mean cluster redshift. (b) Residual from the quadratic $L_m-\alpha$ relationship as a function of the BCG projected separation from the published Abell/ACO cluster center.

and cluster richness. We note in passing, however, that there is evidence to suggest that the internal BCG stellar velocity dispersion is correlated with the residuals. A full investigation of this correlation and its effect on distance determination will be presented in a subsequent paper.

4.3. Computing Estimated Redshifts for BCGs

For all of our discussion here, and in Paper I, we have assumed that the cluster redshifts give the best initial distance estimate to the BCGs with the $L_m-\alpha$ relationship providing independent information on BCG distances that we use to statistically correct for departures of the clusters from a perfect Hubble flow. One can use the $L_m-\alpha$ relationship directly as well, however, to estimate the redshift of any elliptical BCG. As we discuss below, this technique may provide the best distances to clusters lacking proper redshifts.

We estimate redshifts by measuring the brightness of the BCG through a set of apertures, finding the best aperture to place the BCG on the nominal $L_m-\alpha$ relationship. We begin by selecting one of the apertures (the initial choice is not important), assuming that it corresponds to the $10 h^{-1}$ kpc metric radius. The aperture choice thus implies the redshift (by eq. [1] in § 4.1), the corresponding K -correction, and hence the absolute metric magnitude; α is also measured at the aperture selected. We then compare the derived absolute metric magnitude with the expected absolute metric magnitude for the

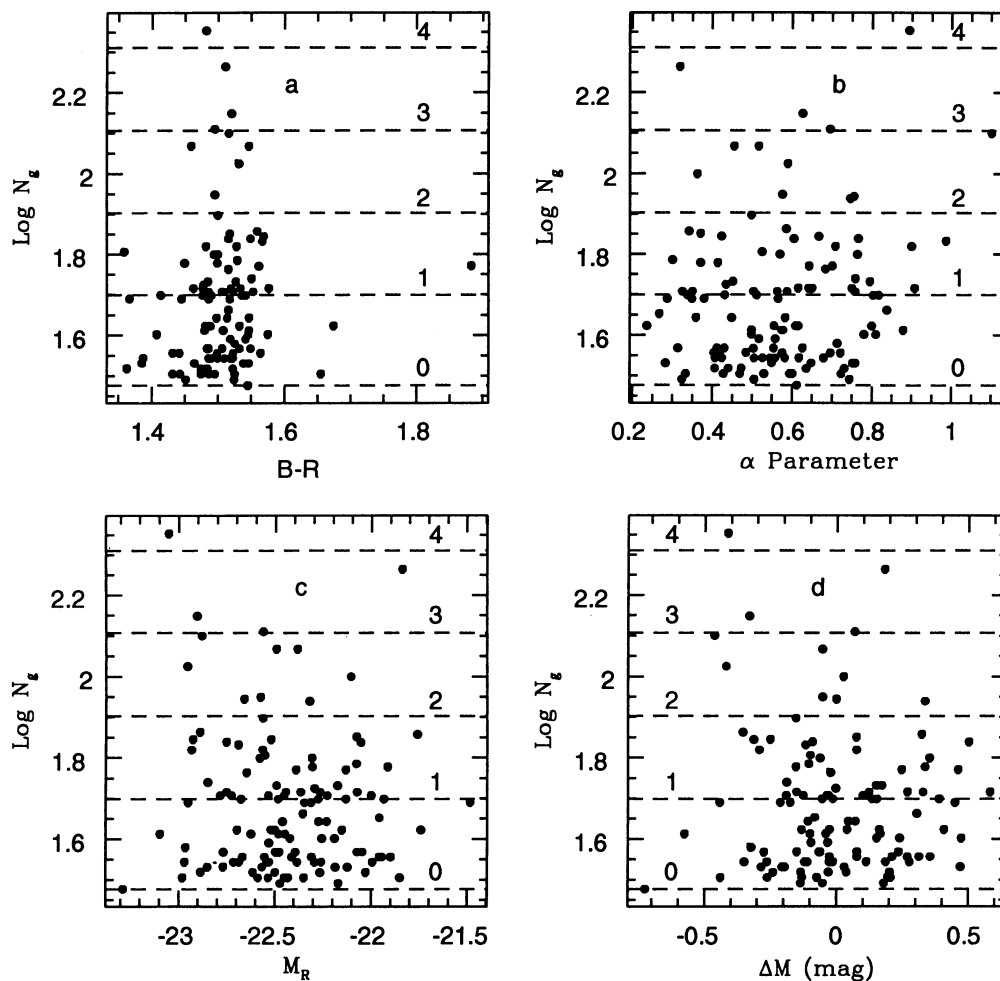


FIG. 7.—(a) Cluster richness as a function of BCG $B-R_c$ color. The horizontal dashed lines represent Abell richness class boundaries. (b) Cluster richness as a function of α . (c) Cluster richness as a function of L_m . (d) Cluster richness as a function of the residual from the quadratic L_m - α relationship.

observed α based on the L_m - α relationship (eq. [2] in § 4.1; valid for the Cousins R_c band, $H_0 = 80$, $q_0 = 0.5$). If the metric aperture is too small, then the BCG redshift will be initially overestimated, and the derived metric luminosity will be brighter than that expected for a BCG really at that redshift. Conversely, if the BCG redshift is underestimated (aperture too big), the derived metric luminosity will be fainter than expected. The size of the error leads to a revised aperture choice. In practice, the proper aperture, and thus redshift estimate, can be found precisely by performing an interpolation between a series of aperture radii (providing that the apertures selected bracket the true metric radius). The final aperture is the one that gives an absolute metric magnitude that exactly matches the prediction from the L_m - α relationship for the α measured. We note that the quadratic fit to the L_m - α relationship in the Johnson B band differs from eq. (2) (§ 4.1) only in the zero point—the first- and second-order coefficients are the same to within the errors. Consequently, this redshift estimation technique is applicable over the full optical bandpass.

The errors in the estimated redshifts depend on α , as has already been discussed by Gunn & Oke (1975). This can be seen as follows. Near the metric radius, the estimated metric luminosity, L_e , varies with estimated metric radius, r_e , away

from the nominal values as

$$L_e = L_m \left(\frac{r_e}{r_m} \right)^\alpha.$$

At any α , this can be cast in terms of surface brightness as

$$\mu_e = \mu_m \left(\frac{r_e}{r_m} \right)^{\alpha-2},$$

where μ_m is the average metric surface brightness for the given α . Now the dispersion in L_m , σ_m , implies an identical dispersion in μ_m . From the equation above, we can see that as μ_e varies from the nominal μ_m the ratio of the estimated distance, D_e to the nominal value will be

$$\frac{D_e}{D_m} = \left(\frac{\mu_e}{\mu_m} \right)^{1/(\alpha-2)}.$$

For small σ_m , this implies that the relative redshift error is

$$\Delta z/z \approx \sigma_m/(2-\alpha).$$

The limits of this expression are simple to understand. For $\alpha = 0$ (a point source) the flux is simply proportional to D^{-2} as

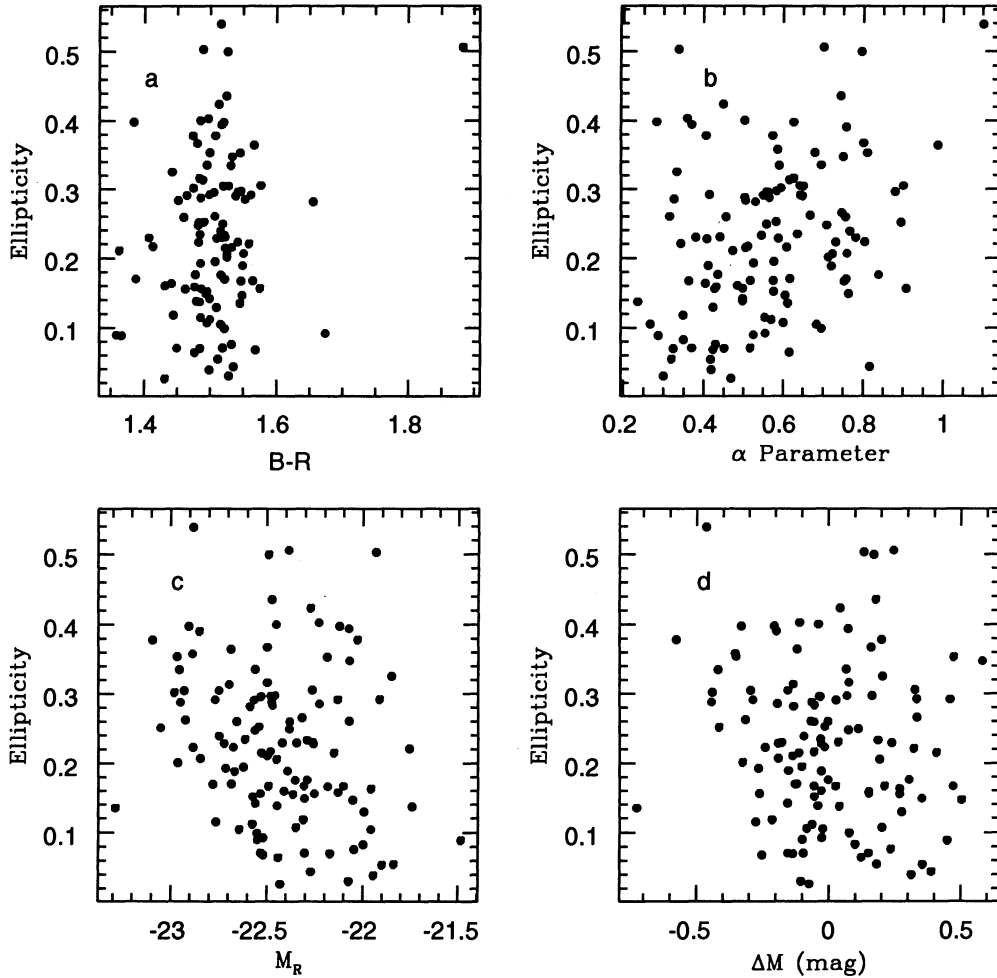


FIG. 8.—(a) BCG ellipticity as a function of BCG $B-R_c$ color. (b) BCG ellipticity as a function of α . (c) BCG ellipticity as a function of L_m . (d) BCG ellipticity as a function of the residual from the quadratic L_m - α relationship.

expected. For $\alpha = 2$ (a constant surface brightness sheet), the flux is simply proportional to the area of the aperture and is insensitive to distance. For the current sample, $\bar{\alpha} = 0.57$ and $\sigma_m = 0.244$ mag, implying a mean distance uncertainty of $\pm 17\%$ per BCG.

Figure 9 shows the estimated redshift (as derived above) as a function of the observed redshift for the 119 BCG in our sample. The solid line is $z_{\text{est}} = z_{\text{obs}}$; the dashed lines denote the $\pm 17\%$ limits. The data are plotted in log-log coordinates to demonstrate that the fractional distance error is independent of redshift. The mean and standard deviations for $(z_{\text{obs}} - z_{\text{est}})/z_{\text{obs}}$ are -0.012 and 0.173 , respectively. Note the excellent agreement between the predicted distance uncertainty and the observed dispersion in $(z_{\text{obs}} - z_{\text{est}})/z_{\text{obs}}$. The scatter about the line $z_{\text{est}} = z_{\text{obs}}$ is relatively insensitive to α , although the error in any given BCG redshift estimate is, of course, dependent on its α -value. The best-fit line is

$$z_{\text{est}} = (0.965 \pm 0.055)z_{\text{obs}} + (0.002 \pm 0.002),$$

with a linear correlation coefficient of 0.849.

The L_m - α redshift estimation scheme is superior to other photometrically based redshift estimators for clusters (at least

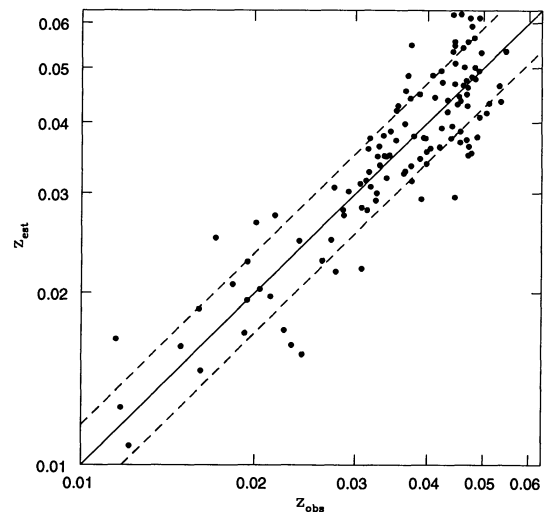


FIG. 9.—Estimated cluster redshift as a function of the observed cluster redshift. The redshift estimates are derived solely from the BCG photometry. The dashed lines show the $\pm 17\%$ limits corresponding to the average expected BCG distance error.

to $z = 0.05$ and most likely to much higher redshifts as long as evolution effects within the central $10 h^{-1}$ kpc remain small compared to the 0.244 mag cosmic scatter). For comparison, the mean and standard deviation for $(z_{\text{obs}} - z_{\text{est}})/z_{\text{obs}}$ using Leir & van den Bergh (1977) redshift estimates for Abell clusters with $z_{\text{obs}} < 0.055$ are -0.125 and 0.299 , respectively. The larger mean indicates a significant bias in the redshift estimate (the error in the mean is ± 0.038), and the scatter is about twice as big as that for the BCG method. For $z > 0.1$, photometrically derived redshift estimates from the tenth-ranked cluster galaxy become seriously biased due to interloper contamination (Postman et al. 1985). While BCG selection can also be problematic when redshift data are not available, the narrow color range and elliptical morphology can make their selection much more robust against foreground and background confusion.

4.4. The Second-ranked Galaxy

As part of our program, we considered the question of whether the *second-ranked galaxy* (SRG) in a cluster might also be a good standard candle. For example, in A1656 (Coma) the R_c -band metric magnitude of the BCG (NGC 4889) is 0.41 mag brighter than the SRG (NGC 4874), but NGC 4874 is itself brighter than many of the BCGs from other clusters. If the L_m - α relationship were just a property of the most luminous elliptical galaxies, then perhaps it might also be valid for the brighter SRGs. This possibility is especially intriguing as one might expect that many of the present SRGs were once BCGs in their own right prior to their original clusters merging with their neighbors.

During the course of our survey, we obtained CCD data for 30 SRGs (see Table 7) as part of the BCG selection process, but also with the above issues in mind. Figure 10 shows where

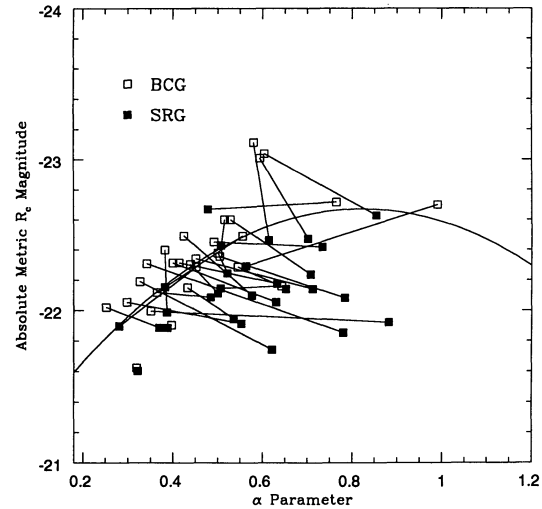


FIG. 10.—Luminosity of SRGs as a function of α . Each SRG (solid square) is linked by a line to its corresponding BCG (open square). The BCG L_m - α relationship from Fig. 1 is also plotted.

these SRGs lie in the L_m - α plane. Each SRG (solid square) is linked by a line to its corresponding BCG (open square). The BCG L_m - α relationship from Figure 1 is also plotted for comparison. Unfortunately, it appears that SRGs are distinct from BCGs in their structural properties. If SRGs were simply fainter versions of BCGs, then one might expect that the SRGs would be displaced along the mean L_m - α relation toward lower α -values (fainter absolute magnitudes). While this is true for a few SRGs, in most cases exactly the opposite shift is seen—the SRGs typically have *larger* α -values than the BCGs, despite their lower metric luminosities.

TABLE 7
SECOND-RANKED GALAXY OBSERVATIONS

Abell	R.A. (J2000)	Decl.	$M_2 - M_1$	α	cz	σ	Notes
076-2	00 ^h 39 ^m 33 ^s .51	+06°48'52".5	0.334	0.383	11916	31	I1566
119-2	00 56 25.53	-01 15 44.9	0.046	0.478	11575	250	
147-2	01 08 11.85	+02 11 35.6	0.393	0.577	12587	75	
533-2	05 01 35.99	-22 36 03.1	0.132	0.522	14642	34	
539-2	05 16 37.34	+06 26 28.0	0.274	0.785	8318	47	
548-A	05 48 38.48	-25 28 38.4	0.033	0.734	11942	57	
576-1	07 21 32.51	+55 45 25.9	0.141	0.553	12177	100	
1142-1	11 00 57.45	+10 30 20.9	0.144	0.653	11150	43	NGC 3492A
1228-2	11 21 42.70	+34 21 47.2	0.207	0.536	10578	39	I2744
1644-2	12 57 49.30	-17 32 44.5	0.405	0.564	14013	50	
1656-2	12 59 35.65	+27 57 33.0	0.412	0.855	7176	15	NGC 4874
1736-2	13 26 48.69	-27 08 36.9	0.650	0.614	13719	49	
1983-2	14 52 55.24	+16 42 09.1	0.457	0.781	13581	52	
2152-2	16 05 26.57	+16 26 32.8	0.411	0.387	
2197-2	16 27 41.19	+40 55 36.8	0.537	0.702	9408	77	NGC 6160
2247-2	16 50 59.06	+81 34 28.5	0.139	0.633	11409	98	
2572-1	23 18 30.39	+18 41 21.1	0.363	0.708	11610	150	NGC 7597
2657-1	23 44 57.44	+09 11 30.7	0.074	0.883	12361	46	
2806-1	00 40 12.92	-56 09 15.0	0.260	0.631	8244	49	
2881-2	01 10 54.46	-17 11 51.6	0.017	0.388	12665	38	
2911-A	01 24 44.69	-38 07 42.3	0.017	0.322	5826	23	NGC 534
3374-2	05 56 54.98	-21 15 06.0	0.133	0.370	14502	36	
3389-2	06 21 26.41	-64 59 37.1	0.170	0.507	8027	44	
3564-2	13 34 11.98	-35 20 14.9	0.193	0.501	
3566-1	13 40 13.34	-35 40 30.2	0.033	0.485	
3570-1	13 46 47.33	-37 54 28.4	0.448	0.622	11237	74	
3733-2	21 01 37.58	-28 01 56.2	0.013	0.507	11505	107	NGC 6998
3747-2	21 08 28.99	-43 29 24.5	0.390	0.281	5086	30	
4038-1	23 47 44.96	-28 08 31.4	0.200	0.713	7940	150	

We emphasize that selection of our SRG sample was highly biased, but in a way that we had presumed would strongly favor SRGs that resembled BCGs; SRGs that were dramatically fainter than the BCGs were systematically excluded. The average $M_1 - M_2$ for these 30 clusters is -0.242 mag, and the scatter about the mean absolute SRG metric magnitude is 0.255 mag. The small offset and artificially small dispersion (Schneider, Gunn, & Hoessel 1983b get a dispersion of 0.55 mag for a complete SRG sample) show that many of the SRGs selected will be brighter than the fainter BCGs. While a complete sample might show the faintest SRGs to fall on the low end of the L_m - α relationship, clearly the brightest SRGs do not fit it at all. The possibility that bright SRGs are fundamentally different from BCGs was also raised by Ryden et al. (1993), who find that these same SRGs are both rounder and have a narrower axis-ratio dispersion than BCGs and ordinary elliptical galaxies. These observations suggest different formation scenarios and/or merging histories for BCGs and these SRGs. In this context, it is interesting to note that Hoessel (1980) argued that the L_m - α relationship was just the track followed by bright ellipticals as they grew by cannibalism (Hausman & Ostriker 1978). This scenario, however, will not work for SRGs, since α is expected to increase rather than decrease as the galaxies accrete their fellow cluster members. The results are intriguing enough to make acquisition of photometric data for a complete SRG sample worthwhile and, indeed, a necessity to fully determine the status of SRGs as standard candles.

5. SUMMARY

We have completed a reinvestigation of BCGs as distance indicators with the specific goal of establishing a frame suitable for measuring galaxy bulk flows at large distances. We emphasize the following aspects of our sample selection and observations:

1. Our sample covers the full sky, except for regions with $|b| < 15^\circ$. Further, the sample is volume limited and as such includes all Abell clusters known with $z \leq 0.05$. In contrast, previous BCG samples have been limited to clusters observable from the north and were not complete or volume limited.
2. We have performed our own BCG selection. The final BCG selected for a given cluster is the galaxy that has the greatest flux within the metric aperture, regardless of position within the cluster. We have also recalculated the redshifts of all clusters, giving the BCGs high weight due to their preferential location at the kinematic centers of the clusters.
3. Large-format CCD cameras were used to ensure the best sky subtraction. Multi-isophote fitting software was used to remove other cluster galaxies and stars seen in projection against the BCGs.
4. A network of overlap galaxy observations ensures that the photometric observations are homogeneous over the sample. We find that all runs have zero points consistent to 0.01 mag and that the random photometric error in the R_c -band metric luminosities is 0.014 mag. We thus expect that any systematic errors in the velocity solutions presented in Paper I, due to photometric calibration errors, will be limited to 75 km s^{-1} .

We have confirmed the existence of the Hoessel (1980) L_m - α relationship for BCGs. The L_m - α relationship appears to be linear for small α , flattening off, however, as α increases. The

L_m - α relationship is described by a second-order curve. We have adopted $r_m = 10 h^{-1}$ kpc, after exploring the dependence of the BCG cosmic luminosity scatter, σ_m as a function of r_m . The basic scatter in L_m is 0.327 mag, which decreases to 0.244 mag, once the relationship between L_m and α is accounted for. The scatter about the L_m - α relationship for $\alpha < 0.6$ decreases to 0.181 mag. The distributions of L_m and L_m - α residuals are both consistent with single Gaussian distributions. The L_m - α relationship can be used to estimate BCG redshifts with 17% mean accuracy. This is as expected, given $\sigma_m = 0.244$ and the mean BCG $\bar{\alpha} = 0.57$. We have conducted an extensive search for additional "second parameters" that might be used to minimize σ_m further as follows:

1. There are no correlations between BCG $B - R_c$ color and L_m , L_m - α residuals, redshift, extinction, BCG ellipticity, or cluster richness, after extinction and K -corrections have been applied.
2. There are no correlations between L_m - α residuals and BCG offsets from the mean cluster velocity or projected cluster center. Both L_m and L_m - α residuals are uncorrelated with cluster richness.
3. BCG ellipticity may weakly correlate with L_m and L_m - α residuals, but is not correlated with α .
4. We have obtained central stellar velocity dispersion measurements for all 119 BCGs. The BCG internal velocity dispersions may correlate with L_m - α residuals. We are currently investigating this and will report our results in a subsequent paper.

We conclude that BCGs are a highly homogeneous population, making them suitable for statistical studies of galaxy peculiar velocities on large scales. We are currently extending the cluster frame defined in Paper I to $z = 0.08$. This program includes the observation of approximately 500 additional BCGs. The deeper sample is well suited to further investigation of the BCG properties considered in this paper.

We are greatly indebted to several of our colleagues who made key contributions to this project. We thank Ofer Lahav and Jim Gunn for useful conversations that helped define the observational program. John Huchra provided invaluable support by supplying essential galaxy redshift data from the CfA Redshift Catalog and, along with Ann Zabludoff, obtained several new galaxy redshifts at our request. Andy Fruchter provided photometric calibrations for a few of the images, and John Tonry obtained images of a few galaxies missed during our own runs. Dave Burstein, Tina Bird, Sandy Faber, Jim Gunn, Bohdan Paczyński, and Michael Strauss provided many critical and constructive discussions for which we are most grateful. This program was based on imaging observations obtained exclusively at the facilities of the National Optical Astronomy Observatories—we thank the Kitt Peak National Observatory and Cerro Tololo Inter-American Observatory TACs for their generous grants of observing time. We also thank the KPNO and CTIO staff for the tremendous support we received throughout this program. M. P. was supported by NASA grant NAGW-2166 and by the STScI Director's Discretionary Research Fund. BCG identifications and astrometry were obtained using the Guide Stars Selection System Astrometric Support Program developed at the STScI.

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