

# The Spiral Structure of Our Galaxy Determined from H II Regions

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Received July 3, 1975, revised January 22, 1976

**Summary.** The overall, large-scale distribution of ionized hydrogen regions in our Galaxy defines a spiral structure comparable to the large-scale spiral patterns observed in external galaxies. In order to outline the spiral arms of our Galaxy we have combined our optical observations (distances of exciting stars and H $\alpha$  radial velocities) with the radio observations of H II regions [H 109 $\alpha$  radial velocities (Mezger, 1970)], using a method proposed by Bok (1971). Our new observations of the radial velocities of 268 H $\alpha$  regions and new distance determinations for 360 exciting stars, have allowed us to define a detailed rotation model exhibiting, in certain sectors, deviations from circular motion, thus permitting us to correctly fit together the stellar and kinematic distances; this rotation model has been supplemented, in the interior part of the Galaxy (galactocentric distance between 4 and 5 kpc, longitudes 330° to 340°) by using H 109 $\alpha$  velocity maxima. The optical detection of very distant H II regions—out to 9 kpc of the Sun—has permitted a good overlap be-

tween the optical data and the H 109 $\alpha$  radio data (obtained over the whole of our Galaxy with nevertheless a distance selection effect as shown in Fig. 11). In addition, we have analyzed the radial velocities and made precise identifications (using H $\alpha$  photographs and radio continuum maps) so as to group together all the H 109 $\alpha$  sources of a single complex. These identifications, along with the use of radio absorption components (molecules, 21-cm), have allowed us to resolve the distance ambiguity for all these H II complexes (with a high certainty for 66%). Eighty per cent of the high-excitation-parameter H II regions thus defined fall along two symmetrical pairs of arms (i.e. four altogether) of 12° inclination. The longitudes at which one sees these arms tangentially correspond exactly to the flux maxima in the radio continuum and in the total 21-cm profile integral.

**Key words:** H II regions — spiral structure

## I. Introduction

In the past it was impossible to reconcile the models of the spiral structure of our Galaxy established from optical results with those based on radio observations. There were a number of reasons for this.

The older optical observations were limited to the solar neighborhood, and the resulting interpretations of structure attributed undue importance to this part of the Galaxy.

Observations at radio wavelengths have the advantage of not being limited by interstellar absorption, thus permitting exploration of the entire Galaxy. The 21-cm line observations, in particular, give the column density of neutral hydrogen integrated along a line of sight as a function of the radial velocity. However, these neutral hydrogen observations have a number of shortcomings. Two of these are common to all kinematic methods of spiral structure determination. They are, first, the necessity of using a rotation model to transform the radial velocities to distances. The determination of such a model is plagued by non-circular motions (Shane and Bieger-Smith, 1966; Kerr, 1970; Burton, 1973):

small fluctuations ( $\sim 1$  kpc) and larger scale variations as a function of galactocentric longitude (i.e., the difference between the northern and southern hemispheres as has been shown by Kerr, 1970). Second, for the interior parts of the Galaxy the rotation model gives two possible distances for a given observed radial velocity, and it is sometimes difficult to choose between the “near” and the “far” distance. In addition, the 21-cm observations have the particular shortcoming of resulting from an integration along the line of sight. Burton (1973) has shown that “In fact it is not clear to what extent a particular peak in a line profile owes its characteristics to streaming motions, to a density concentration or to a variation in temperature”. This means that the spiral arms cannot be determined with any accuracy directly from 21-cm line data. This shortcoming does not exist in the case of observations of discrete H II regions.

In recent years much progress has been made in the spiral structure problem: (i) important new optical data has been obtained (OB stars and H II regions)

especially in the southern hemisphere, (ii) radio-recombination lines have been discovered and extensively observed (Mezger, 1970; Reifenstein *et al.*, 1970; Wilson *et al.*, 1970 etc.) as have the radio molecular lines (Manchester *et al.*, 1970; Goss *et al.*, 1969; Robinson *et al.*, 1971; Caswell and Robinson, 1974; Wilson, 1972 etc.), (iii) a model-fitting approach has been applied to the interpretation of the 21-cm profiles (Burton and Shane, 1970), and new analyses of H I data in terms of the gravitational theory of density-waves has been done (Burton, 1973).

The purpose of this article is to try to establish *from the study of a homogeneous population—H II regions and their exciting stars—a coherent schema of the spiral structure of our Galaxy* which is compatible with the optical and the radio results. This method, recommended by Bok (1971), has been successfully used in the Carina arm by Bok *et al.* (1970). We shall try here to generalize it to the entire Galaxy as it has been briefly shown by Georgelin (1975a). To do this we use our new determinations of distances of exciting stars and of H $\alpha$  radial velocities (Georgelin, 1975), as well as all available H 109 $\alpha$  radial velocities and particularly those of Reifenstein *et al.* (1970) and of Wilson *et al.* (1970).

Our new observational results have also permitted us to greatly extend the optical exploration zone, particularly in Cygnus and in Carina (up to 8 and 9 kpc). This considerably increases the region of overlap of the optical with the radio results, and therefore facilitates the harmonisation of the two.

For this synthesis, based primarily on H II regions and exciting stars, we have also used the results for other optical constituents (supergiants, young clusters) and radio constituents (H I, molecules) of Population I. To do this we had to examine all the optical and radio data in detail, nebula by nebula. This work was simplified in the Carina region and in the Crux-Centaurus-Norma region by the extensive contributions in the catalogues of Hine (1971) and Smith (1972).

## II. Method of Analysis

### 1. Observational Data: Accuracy and Homogeneity

To have any hope at all of defining the spiral structure of our Galaxy, one must use a distance scale (spectrophotometric and kinematic) based on data which are as homogeneous and accurate as possible. The radial velocities of H $\alpha$  regions have been measured by the Fabry-Perot interferometric method described by Courtès (1960); thus, at a mean reciprocal dispersion of 20 Å mm<sup>-1</sup>, we have measured the H $\alpha$  radial velocities of about 8000 elements of a few square seconds-of-arc in 268 H II regions distributed quasi-uniformly over the entire galactic plane. The spectrophotometric distances of the 360 exciting stars of these H II regions have been obtained from spectral classifications and *UBV* photo-

metry effectuated by about fifty different authors, principally Crampton (1971), Crampton and Fisher (1974), and Georgelin *et al.* (1973). The absolute magnitude-spectral type calibration which we use (Crampton and Georgelin, 1975) takes into account important advances made by Conti and Alschuler (1971), Walborn (1972), and Balona and Crampton (1974).

All of the observational results, the critical analysis of the spectral classifications and absolute magnitudes, as well as a detailed study of the exciting OB clusters have been assembled by Georgelin (1975).

In this work we have extensively used the H 109 $\alpha$  radial velocities observed by Reifenstein *et al.* (1970), Wilson *et al.* (1970), Caswell (1972), Dickel and Milne (1972).

### 2. Grouping of the Various Measurements

Because of the intrinsic scatter of radial velocity within each H II region—i.e.,  $\sigma = 4.2$  km s<sup>-1</sup> (Table 1)—it is not correct to derive a kinematic distance from each measurement, at the risk of finding a spiral structure more confused than the reality. Therefore the observations are judiciously grouped and one distance is derived for each H II region or each complex H II region from mean radial velocity. One must exercise great care to define real physical groupings; we have done this with the aid of H $\alpha$  photographs and radio continuum maps. Thus, for example, the 36 H II regions detected in radio at  $\ell = 332^\circ$  and  $\ell = 337^\circ$  have been grouped into nine H II complexes of various sizes, allowing—finally—the structural interpretation of this crucial zone of the Galaxy.

### 3. Rotation Model

Relatively recent studies (as Courtès *et al.*, 1967) of the detailed and accurate kinematics of the large nearby galaxies has shown that they present great irregularities, on both a large scale and on a small scale, of the rotational velocity  $\Theta$  as a function of  $R$  (distance from the galactic center) and of  $L$  (galactocentric longitude); the situation is the same in our Galaxy. Consequently we have tried to define separate rather detailed rotation models for two zones of galactocentric longitude, in order to best tie in the kinematic distances with the spectrophotometric distances. In the optical zone we have used the distances of the exciting stars and the radial velocities of the H $\alpha$  regions, and in the interior parts of the Galaxy the maximum velocities of the H 109 $\alpha$  regions.

### 4. Ambiguity of the Distances

For the internal parts of the Galaxy the choice between the two distances corresponding to the same observed radial velocity was made with the aid of the following

complementary information, listed in order of decreasing certitude:

- (1) The distance of the exciting star. If known, all ambiguity is removed.
- (2) The fact that an H II region is observed optically. This is an argument in favor of the "near" solution.
- (3) Absorption components of radio lines, at greater absolute velocities than that of the H II region are evidence in favor of the "far" solution.
- (4) The fact that an H 109 $\alpha$  region is not observed optically is an argument in favor of the "far" solution.
- (5) The latitude thickness of the arm. If the latitude range is large, the "far" distance is unacceptable.
- (6) Absorption components of radio lines at smaller absolute velocities than that of the H II region are evidence slightly in favor of the "near" distance.
- (7) The velocity continuity along the arms or the continuity of the arms. As observation of external galaxies demonstrates, arms have a spiral shape with pitch angles varying from 0° to 30°; the sense of this inclination is known for our Galaxy. In the absence of other indications of the distance (near or far), we choose the solution permitting the alignment of the H II regions along spirals of plausible inclination. There is ambiguity of distance for 67 H II regions among the 100 important H II regions. This ambiguity is resolved with a high certainty for 44 H II regions (criteria 1 to 4) and with a lower certainty for 23 H II regions (criteria 5 to 7). Details are given in Table 10.

### 5. Weighting of the H II Regions

The H II regions which trace the spiral arms in external galaxies are "important" (brilliant and extended). In order to trace the spiral arms in our Galaxy it is therefore best to place ourselves in the same conditions as the observer of external galaxies, as has been suggested by Bok (1969). The most representative criterion of the "importance" of the H II regions is the excitation parameter  $U$ ;  $U$  is proportional to  $(SD^2)^{1/3}$ , where  $S$  is the radio flux emitted by the H II region and  $D$  is its distance. In this study we have therefore selected and weighted the H II regions in such a way as to favor the largest excitation parameters, so as to simulate the observational conditions for external galaxies, and to reduce the selection effect due to distance from the Sun.

### 6. External Galaxies

In this paper we try to establish, from the study of H II regions, a coherent scheme of the spiral structure of our Galaxy. We need guidance in this analysis from the results of studies of the neighboring galaxies. For the optical astronomers it is easy to obtain the distribution of H II regions in spiral galaxies; these studies can be made with great precision in M 31 (Arp, 1964) and

Table 1. Dispersions of radial velocities

Scale	Standard deviation (km s <sup>-1</sup> )
Within an individual H II region	4.2
Between H II regions of an element of spiral arm (500 pc)	4.5
Between H II regions of an entire spiral arm (5000 pc)	7.4
Between H II regions within a radius of 5 kpc around the sun	6.7

in M 33 (Boulesteix *et al.*, 1974). From the observations of H II regions in external galaxies we can conclude:

— that the arms are observed approximatively between  $R=3$  and  $R=12$  kpc;

— that the pitch angle (Danver, 1942) is:

for Sb galaxies  $\mu=75^\circ 2$  with  $\sigma_{r.m.s.}=5^\circ 9$  and for Sc galaxies  $\mu=72^\circ 5$  with  $\sigma_{r.m.s.}=6^\circ 4$  and

— that there is a very precise symmetry in the spiral arm pattern, especially in the inner regions, and particularly for M 33, M 51 and M 31, but that these galaxies are not symmetrical in their kinematic properties but instead exhibit very significant non-circular motions.

Moreover, H I observations in external galaxies show that the H I spiral arms are strong in the outer parts of the galaxies where H II regions are inconspicuous, and vice versa in the inner parts.

### III. Rotation Model—Kinematic Distances

The determination of a rotation model from the radial velocities of H II regions is possible because of the sufficiently weak dispersion of these velocities. Table 1 shows the standard deviations found on different scales.

In various papers (Georgelin and Georgelin, 1970b; Georgelin, 1975; Crampton and Georgelin, 1975) we have studied the kinematics of optical H II regions ( $U_0$ ,  $V_0$ ,  $A$ ,  $\alpha$ ,  $R_0$ , etc.) with extensive comparison with the results obtained by several authors for stellar Population I objects. From our radial velocities of H II regions we have also made a tentative (Georgelin, 1975) approach to involve streaming motions along the spiral arms in terms of density-wave theory for the gas (Roberts, 1969). But in this paper we limit the kinematic study to the critical parameters used for kinematic distance determination.

#### 1. Mean Rotation Curve

For 151 optical H II regions we have obtained both the radial velocity of the H II region and the spectrophotometric distance of the exciting star ( $s$ ). From these data, with  $R_0=10$  kpc, we have calculated by means of a least-

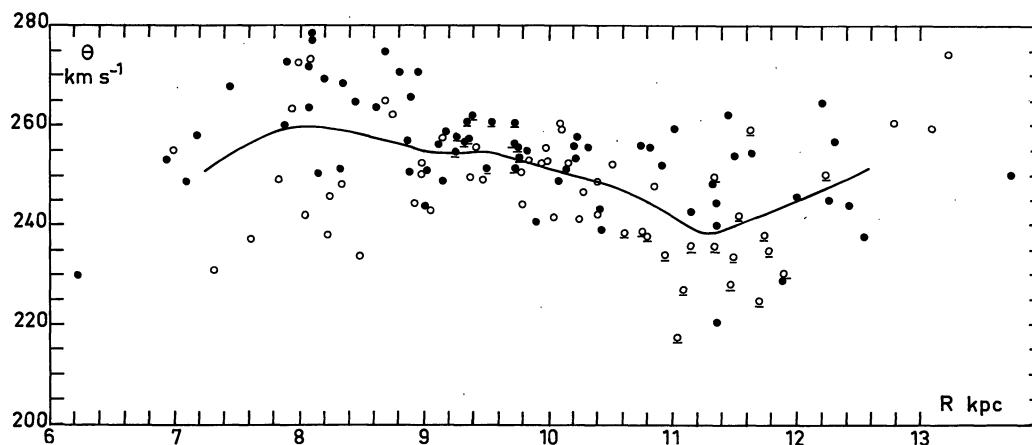


Fig. 1. Rotation curve of our Galaxy calculated from distances of exciting stars and radial velocities of H $\alpha$  regions. The abscissa represents the distance from the galactic center and the ordinate the rotation velocity. Underlined open circles corresponds to Perseus feature and underlined filled circles correspond to Carina feature.  $\bullet$ :  $\ell > 180^\circ$ ;  $\circ$ :  $\ell < 180^\circ$

squares method:  $U_0 = -7.2 \text{ km s}^{-1}$  and  $V_0 = +14.7 \text{ km s}^{-1}$ . In this paper the rotation model and all the distances have been computed with these values ( $V_{\text{HII}}$ ), even if we have given the  $V_{\text{LSR}}$  radial velocities for easier comparison. Figure 1 gives, for each H $\alpha$  region, the rotation velocity  $\Theta$  around the galactic center, calculated in the classical way by

$$\Theta = (R/R_0) (\Theta_0 + V_{\text{HII}}/\sin \ell)$$

and

$$R^2 = R_0^2 + r^2 - 2r R_0 \cos \ell,$$

where  $R$  and  $R_0 = 10 \text{ kpc}$  are the galactocentric distances of the H II region and of the Sun, respectively;  $\Theta$  and  $\Theta_0 = 250 \text{ km s}^{-1}$  are the rotational velocities around the galactic center;  $r$  is the distance to the exciting star; and  $V_{\text{HII}}$  is the radial velocity of the H II region corrected for the solar motion. The northern H II regions ( $0^\circ < \ell < 180^\circ$ ) are represented by open circles, and the southern H II regions ( $180^\circ < \ell < 360^\circ$ ) by filled circles. The mean curve was found by taking a running mean, weighted by  $\sin \ell$ , over groups of ten points. The mean rotation curve differs from the Schmidt model (Schmidt, 1965) by:

- (1) a maximum around 8 kpc;
- (2) a minimum around 11.3 kpc (Perseus arm);
- (3) an increase for  $R > 12 \text{ kpc}$ .

These last two points have already been noted by Humphreys (1970), Balona and Feast (1974), and Georgelin (1975) respectively for supergiants, OB stars, and H II regions.

The last point has also been noted by Burton and Bania (1974). They show the response of the correlation coefficient  $\rho$  (correlation between the observed deviations from circular motion and the model non-circular motions calculated at the locations of the observed objects) to variations in the parameter  $k$  specifying the model basic rotation curve. This correlation analysis (Burton and Bania, 1974) is made for associations, supergiants and H II regions, and the results are very

conclusive for H II regions where the correlation is certainly significant despite the limited number of regions used. The maximum correlation is obtained for a basic rotation curve characterized by  $k = +3$  instead of  $k = 0$  (Schmidt model). From H II regions we obtain (Fig. 2), for  $12 < R < 13 \text{ kpc}$ ,  $\Theta = 250 \text{ km s}^{-1}$ , corresponding to  $k = +3$  in agreement with Burton and Bania (1974).

## 2. North-south Discontinuity from Optical Results

A detailed, point-by-point analysis of Fig. 1 shows that the southern hemisphere points ( $180^\circ < \ell < 360^\circ$ ) tend to have a greater rotational velocity. In order to confirm this, we have separated the northern and southern points and smoothed in the same way as for the mean curve. Figure 2, on which we have indicated the root-mean-square dispersions of these two rotation curves, confirms their clear separation, particularly for 8.5 kpc and for 11.2 kpc. It is important to remember that this effect (southern rotation faster than northern rotation) is in the same sense as that observed in H I by Kerr (1970) between 8.2 and 10 kpc. We also note that from 7 kpc to 10 kpc the northern curve is very close to the Schmidt model, while the southern curve is located significantly above, the difference reaching  $14 \text{ km s}^{-1}$  at 8.3 kpc; beyond 12 kpc the curves are less well defined, but the northern and southern curves seem to be identical while remaining above that of the Schmidt model.

In an extensive work on supergiants, Humphreys (1970) has shown that their kinematics in the longitude sector  $30^\circ < \ell < 210^\circ$ , are different from those in the sector  $210^\circ < \ell < 30^\circ$ . A re-analysis of her data, in the same manner as for the H II regions, shows that the effect for the supergiants is indeed in the same sense as that for H II regions—i.e., the rotation curve for the southern hemisphere ( $180^\circ < \ell < 360^\circ$ ) is located above that for the northern hemisphere ( $0^\circ < \ell < 180^\circ$ ).



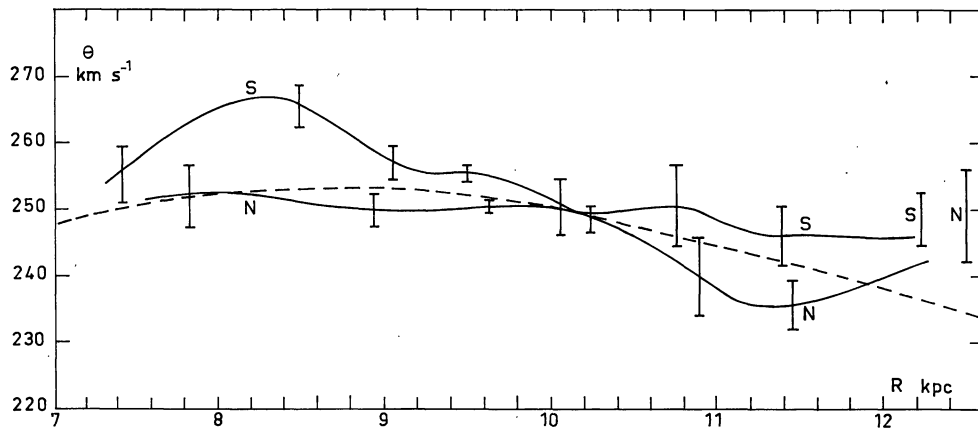


Fig. 2. Rotation curves of our Galaxy calculated from distances of exciting stars and radial velocities of H $\alpha$  regions for longitudes  $0^\circ < \ell < 180^\circ$  (north) and  $180^\circ < \ell < 360^\circ$  (south). The dotted line represents the Schmidt model

### 3. Deviations from Circular Motion. Carina Section and Perseus Section

Two undisputed features in the Population I structure of our Galaxy are the so-called “Carina section” and “Perseus section”. The “Carina section” (with  $2.5 < r < 8$  kpc,  $1282^\circ < \ell < 286^\circ$ ,  $9 < R < 10.5$  kpc, external longitude  $\ell = 282^\circ$ ) was extensively studied by Bok *et al.* (1970), and the “Perseus section” (with  $100^\circ < \ell < 150^\circ$ ,  $2 < r < 3$  kpc,  $11 < R < 12$  kpc) was discovered by Morgan *et al.* (1952) from OB stars and by Oort *et al.* (1958) from H I observations.

The analysis of the rotation curve (Fig. 2) shows, first, that the points of the “Carina section” between  $9 < R < 10$  kpc are located even higher than the southern rotation curve; and second, that the points of the Perseus arm ( $10.8$  kpc  $< R < 11.8$  kpc) are located lower than the northern rotation curve, which confirms the deviation from circular rotation already well known for these two regions. From our values we have determined, for the Carina arm ( $9.1$  kpc  $< R < 9.5$  kpc), a slope  $A = 19$ , which corresponds to a deviation from circular rotation of ( $\Delta V = -2A(R - R_0) \sin \ell \approx 6$  km s $^{-1}$ ), an effect already observed for H II regions (Bigay *et al.*, 1972) and for hot stars (Humphreys, 1972; Burton and Bania, 1974, Fig. 2).

### 4. Rotation Curve of Ionized Gas from Radio Results for $3.9 < R < 5$ kpc

Between  $\ell = 330^\circ$  and  $\ell = 339^\circ$  the H 109 $\alpha$  results of Wilson *et al.* (1970, Fig. 1, p. 380) show two alignments of numerous H II regions at  $\ell = 332^\circ$  and  $\ell = 337^\circ$ . These two typical alignments show a “hole” of 4 kpc around the tangential point which could mean either that there are no H II regions in this zone—in which case we would have a rather curious 4 kpc interarm gap—or that the observed velocities are not as great as the maximum velocity predicted by the Schmidt model.

Table 2. Maximum radial velocities observed near  $335^\circ$

Component	$V_{\text{LSR}}$	Reference
H 109 $\alpha$	-93, -91, -89, -89, -88	Wilson <i>et al.</i> (1970)
OH absorption	-89, -88	Goss <i>et al.</i> (1970)
CH $_2$ O	-90, -85	Whiteoak and Gardner (1970)
Interstellar Ca absorption	-94, -82, -82, -82	Chu-Kit (1973)
	-82, -81	

Therefore, we have searched in the sector  $330^\circ < \ell < 340^\circ$  for maximum velocities measured for each separate constituent. The results are given in Table 2.

We have deduced from these results a mean observed maximum velocity of  $-90$  km s $^{-1}$  at  $\ell = 332^\circ$ , and  $-93$  km s $^{-1}$  at  $\ell = 337^\circ$ . No significant error is introduced by considering this observed maximum velocity as the velocity at the tangential point, because two weak effects acting in opposing directions tend to compensate each other:

- (1) Because of the statistical scatter of the radial velocities, one risks overestimating the velocity of the tangential point by a few km s $^{-1}$  when one selects the maximum velocity observed.
- (2) There is not necessarily an H II region at the tangential point; thus the maximum measured velocity may correspond to a point located in front of or behind the tangential point, whose velocity, therefore, is underestimated.

The two maximum velocities obtained for the longitudes  $332^\circ$  and  $337^\circ$  allow modification of the rotation curve of our Galaxy for the southern hemisphere. With  $V_{\text{LSR,max}} = -90$  km s $^{-1}$  at  $\ell = 332^\circ$  and  $V_{\text{LSR,max}} = -93$  km s $^{-1}$  at  $\ell = 337^\circ$  one obtains, for the rotation velocities,  $\Theta = 210$  km s $^{-1}$  and  $193$  km s $^{-1}$ , respectively. These values are respectively  $19$  km s $^{-1}$  and  $13$  km s $^{-1}$  less than the rotation velocities predicted by the Schmidt model.

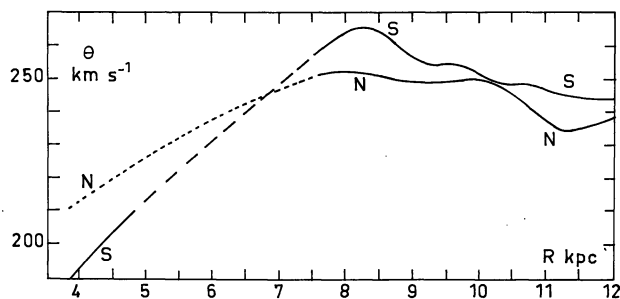


Fig. 3. Rotation curves of our Galaxy calculated from distances of exciting stars and radial velocities of H $\alpha$  regions for  $7 < R < 12$  kpc and from H 109 $\alpha$  velocity maxima for  $3.9 < R < 5$  kpc in the southern hemisphere. Schmidt model for  $R < 7$  kpc in the northern hemisphere. -----: Schmidt model; —: H II regions

This is to be compared to the somewhat weaker asymmetry, in the same sense, of about  $10 \text{ km s}^{-1}$  observed at 21-cm by Kerr (1970) between the northern and southern hemispheres. The observed rotational dissymmetry ( $\Delta\theta = 20 \text{ km s}^{-1}$ ) at  $R = 4.5$  kpc, for which  $\theta = 200 \text{ km s}^{-1}$ , for points  $65^\circ$  apart with respect to the galactic center, is not particularly important if one compares this to the results obtained in external galaxies. For example, in M 33, for points diametrically opposed,  $\Delta\theta = 40 \text{ km s}^{-1}$  at  $R = 4$  kpc, for which  $\theta = 70 \text{ km s}^{-1}$ ; and in the galaxy NGC 6946 (Sc multi-arm) it is even greater:  $\Delta\theta = 60 \text{ km s}^{-1}$  at  $R = 4$  kpc, at which  $\theta = 100 \text{ km s}^{-1}$  (Monnet, 1972).

#### 5. "Optical-radio" Rotation Model (Fig. 3)

For the northern hemisphere we have used the H $\alpha$  results for  $7 < R < 13$  kpc, and the Schmidt model for  $R < 7$  kpc since the observed maximum H 109 $\alpha$  velocity agrees with the maximum velocity predicted by the Schmidt model.

For the southern hemisphere we have used the H $\alpha$  results for  $7 < R < 13$  kpc, and the maximum observed H 109 $\alpha$  velocity for  $3.9 < R < 5$  kpc. For  $5 < R < 7$  kpc there are very few H II regions and we cannot consider the observed maximum H 109 $\alpha$  velocity as the velocity at the tangential point; thus for  $5 < R < 7$  kpc we have interpolated the results of the two other sections of the rotation curve. It is clear that this section of the rotation curve is poorly defined but it has been used for the computation of only 11 kinematical distances of H II regions.

The use of these separate northern and southern models has the two following advantages: it eliminates the 4 kpc "hole" near the tangential point at longitudes  $\ell = 332^\circ$  and  $337^\circ$ , and it greatly improves the correlation between the kinematic and spectrophotometric distances, particularly for the H II regions located between  $\ell = 180^\circ$  and  $\ell = 270^\circ$ .

#### 6. Calculation of the Kinematic Distances

In the rest of this article, for tracing the spiral structure we use the stellar distances whenever they are available. For optical and radio H II regions for which spectrophotometric distances of the exciting star are not available, we have calculated the "kinematic" distances using the northern and southern rotation models. These kinematic distances and a comparison with the stellar distances have been discussed elsewhere (Georgelin, 1975).

### IV. Identification and Correlation of Optical with Radio Data

#### 1. Optical Results

Young clusters, exciting stars, and H $\alpha$  regions allow one to trace the spiral structure of our Galaxy in the solar neighborhood. We have published elsewhere (Crampton and Georgelin, 1975, Fig. 1) a study of the distribution of H $\alpha$  regions obtained from exciting star distances (based upon recent absolute magnitude calibrations) and from kinematic distances (based upon the rotation model of the H II regions of § III). However, a discussion and a synthesis of these results in conjunction with those concerning the distribution of OB clusters—Becker and Fenkart (1971), Moffat and Vogt (1973)—and associations is necessary in order to disentangle the most important and best established spiral tracers. We shall show here, for a particular example, the type of study of detail that we have carried out for the entire Milky Way.

#### Associations, Young Clusters, Exciting Stars, and H II Regions in the Perseus Arm

Figure 4a is a wide-field H $\alpha$  photograph of Dubout (1975). We have indicated in Fig. 4b, on the same scale, the positions of associations, clusters, and H II regions. Examination of the positions, of the morphology, of the distances, and of the radial velocities shows that the Perseus arm is distinguishable from the local feature not only by distance (or radial velocity) but also by its position in the sky, as shown by the dashed line.

In Fig. 4 the local feature is at positive latitudes above the dashed line. The local feature is at a mean distance of 0.8 kpc; it includes three associations (Cep OB 2, Cep OB 3 and Cep OB 4), two clusters, and nineteen H II regions detailed in Table 3.

In Fig. 4 the Perseus arm is at low negative latitudes below the dashed line. The Perseus arm is at a distance of 2.5 to 5 kpc; it includes ten associations (Cep OB 1, Cep OB 5, Cas OB 2, Cas OB 5, Cas OB 4, Cas OB 7, Cas OB 1, Cas OB 8, Per OB 1 and Cas OB 6), twenty-one clusters, and twenty-nine H II regions detailed in Table 3. The boundary between these two groups permits, moreover, a more precise delimitation of the associations.

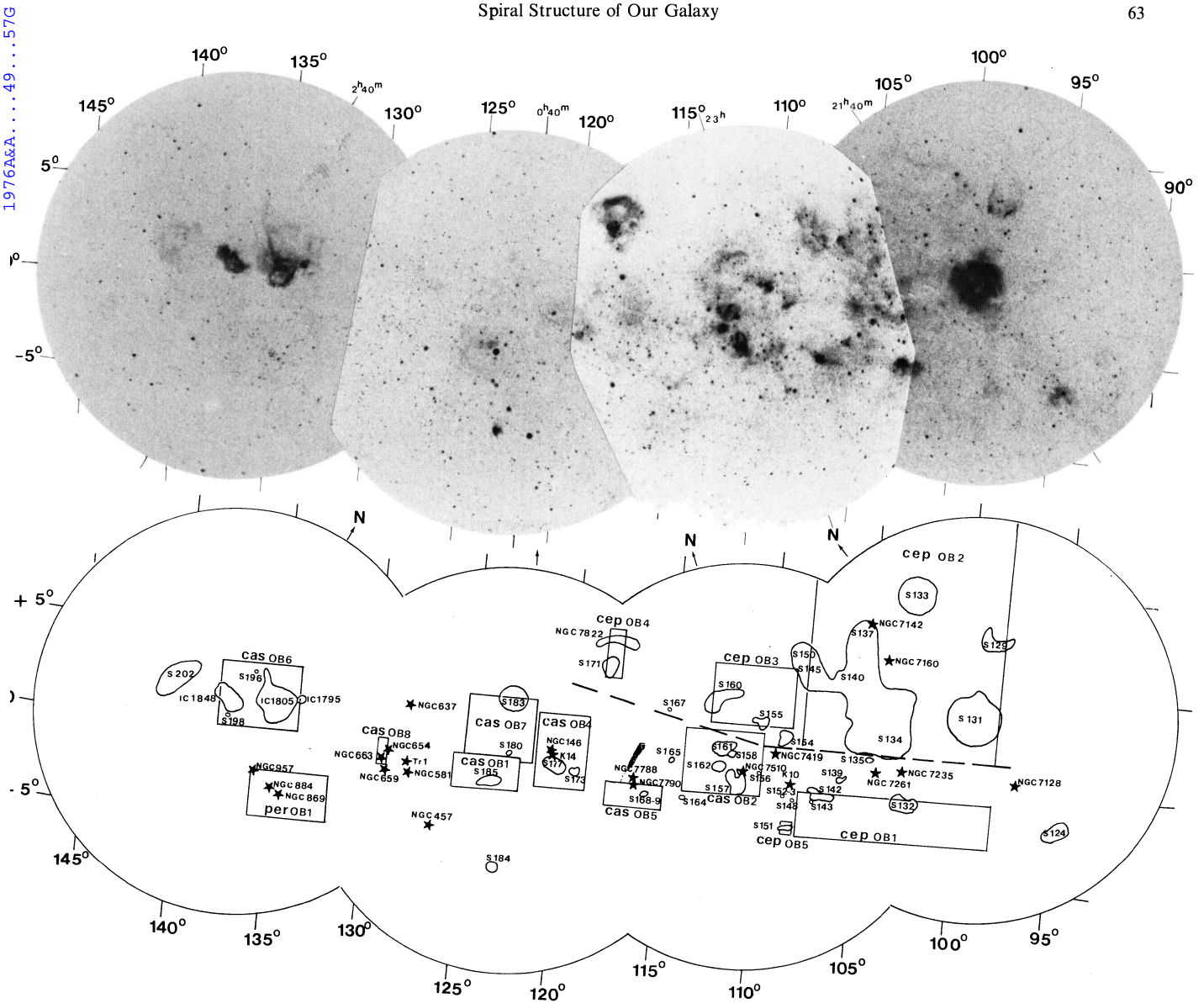


Fig. 4. (a) Monochromatic  $H\alpha$  photographs of the Perseus arm between  $\ell = 95^\circ$  and  $\ell = 145^\circ$  (Dubout, 1975). (b) On the same scale, an identification chart of  $H\text{ II}$  regions, stellar associations and OB clusters. Above the dashed line the local feature at positive latitudes and mean distance of 0.8 kpc. Below the dashed line the Perseus arm at low latitude and distance between 2.5 and 5 kpc

Note also that the dispersion in distance is greater for the second group (the Perseus arm), even allowing for the greater imprecision due to its distance. It appears that there are in fact many sub-groups at different distances in the interior of the Perseus arm. These sub-groups are:

a) From  $\ell = 100^\circ$  to  $\ell = 106^\circ$  the exciting stars and clusters have stellar distances between 3.1 and 4.1 kpc and latitudes between  $-0^\circ6$  and  $+0^\circ8$ ; this is therefore one sole group.

b) From  $\ell = 106^\circ$  to  $\ell = 123^\circ$  one can distinguish two subgroups by their distances:

— A sub-group at 2.5 kpc is comprised of the  $H\text{ II}$  regions and clusters S 142 (NGC 7380), S 143, NGC 7510, S 157, S 161, S 162, S 163, S 164, S 165, S 168,

S 169, NGC 7788, NGC 7790, S 170, S 173, S 175, S 177, King 14, NGC 146, S 184.

— Another, at about 5 kpc, is more localised ( $106^\circ < \ell < 112^\circ$  and  $-3^\circ < b < +1^\circ$ ) and contains the two clusters NGC 7419 and King 10, and some small  $H\text{ II}$  regions (S 151, 152, 153, 148, 149, 156, 158—for 148 and 149 we have only preliminary data for the exciting stars).

c) From  $\ell = 124^\circ$  to  $\ell = 133^\circ$  there are no  $H\text{ II}$  regions, but seven clusters at distances between 2.1 and 2.8 kpc (NGC 457, NGC 581, Tr 1, NGC 637, NGC 654, NGC 659, NGC 663) which are in the extension of the first of the preceding sub-groups.

d) From  $\ell = 133^\circ$  to  $\ell = 138^\circ$  at 2.5 kpc, the associations Per OB 1 ( $h$  and  $\chi$  Per) and Cas OB 6 (IC 1795, IC 1805,

Table 3

HII Region	Cluster	d (kpc)	$V_{LSR}$ (km s <sup>-1</sup> )
a) Perseus arm			
	NGC 7128	3.1	
	NGC 7235	4.1	
S 132		3.6	- 50.9
	NGC 7261	3.2	
S 139		3.3	- 41.0
S 142	NGC 7380	2.4	- 35.8
S 143		3.7	- 28.8
	King 10	4.8	
S 151		5.1 k	- 54.9
S 152		3.4 k	- 44.3
S 153		4.0	- 52.8
	NGC 7419	6.0	
S 156		5.7	- 50.0
	NGC 7510	2.9	
S 157	An. Base 1 3	2.5	- 37.7
S 158		(4.3)	- 59.6
S 161		2.2	
S 162		3.5	- 45.2
S 163		2.3	- 42.6
S 164		3.1 k	- 44.0
S 165		1.6	- 27.5
S 168		3.0 k	- 43.5
S 169		1.5	
	NGC 7788	2.4	
	NGC 7790	3.2	
S 170		2.2	- 53.3
S 173		2.7	- 37.0
S 175		(1.7)	- 50.7
S 177		2.5	
	King 14	2.4	
	NGC 146	2.4	
S 184		2.2	- 27.8
	NGC 457	2.8	
	NGC 581	2.4	
	Tr 1	2.4	
	NGC 637	2.1	
	NGC 654	2.5	
	NGC 659	2.1	
	NGC 663	2.1	
IC 1795		2.3	- 40.7
S 190	IC 1805	2.3	- 46.0
	h Per	2.1	
	X Per	2.5	
	NGC 957	2.1	
S 196		3.6 k	- 43.0
S 198		2.8	- 33.5
S 199	IC 1848	2.3	- 36.6
b) Local feature			
S 125	IC 5146	1.0	+ 4.8
S 129		0.4	- 9.6
S 130 *			
S 131	IC 1396 (Tr 37)	0.9	- 2.2
S 133 *			
S 134		0.9	
	NGC 7160	0.7	
S 135		1.4	- 17.6
S 136 *			
	NGC 7142	0.6	
S 137		0.6	
S 140		0.9	- 10.0
S 144		0.1 k	- 1.7
S 145 *			
S 150 *			
S 154		1.0 k	- 11.4
S 155	Cep OB 3	0.7	- 15.1
S 160		0.8	
S 171	Cep OB 4	0.8	- 10.7

\* S 130, S 133, S 136 have  $b > 9^\circ$

S 145 and S 150 are in the same complex as S 140

IC 1848, S 196, S 198 and S 201—this last region having been studied by Dickinson *et al.*, 1974) continue the same alignment as the preceding associations.

## 2. H II Regions Observed Simultaneously in H $\alpha$ and H 109 $\alpha$

Since the optical data and the radio data are complementary it is important to make sure that there is a fair amount of overlap between them. This requires that a sufficiently large number of H II regions be observed simultaneously by optical and radio means, and that the optical regions be observed out to fairly large distances.

### Proportion of H 109 $\alpha$ Sources Detected Optically

Four hundred twenty H II regions are detected optically and 1300 H II regions are observed in the radio continuum. The radial velocities have been measured for 258 optical regions (H $\alpha$ ) and for 212 radio regions (H 109 $\alpha$  or H 158 $\alpha$ ). The optical and radio velocities are known simultaneously for only 58 of them, i.e., 22% of the measured optical regions and 27% of the radio regions. The radio observations are limited to the most intrinsically intense objects whereas the optical observations are mainly limited by absorption.

Optically, one can detect H II regions whose emission measures are 100 times weaker than the radio detection limit (Monnet, 1974). Hence even a weak radio region will be detected optically unless the absorption is greater than about five magnitudes.

### Detection of Very Distant H II Regions

The optical detection of H II regions and of their exciting stars was formerly limited to a fairly small distance: about 3 kpc. In the past few years great optical progress has been made, on the one hand; and on the other hand the systematic search for H $\alpha$  regions and of hot stars in the vicinity of distant radio sources has led to the discovery of very distant H $\alpha$  regions.

Here are some examples of very distant H II regions detected optically:

a) The measurements of H $\alpha$  radial velocities and the distances of exciting stars have permitted us to detect, optically, very distant H II regions at  $l = 291^\circ$ . As Fig. 5 shows, one observed some nearby optical nebulae such as Gum 38 a (NGC 3576) and Gum 37 in this region, and the somewhat farther region RCW 55 with kinematic distance 5.7 kpc and stellar spectrophotometric distance 5.1 kpc (CPD -62° 1824). Figure 5 shows, especially, three very distant optical and radio regions: Gum 38b (NGC 3603) with an H $\alpha$  kinematic distance of 8.0 kpc, Gum 35 excited by Wolf-Rayet star LSS 2063 of stellar spectrophotometric distance 7.9 kpc, and the source 291.9-0.7 with an H $\alpha$  kinematic distance of 9.7 kpc. Thus there is a very good overlap between the



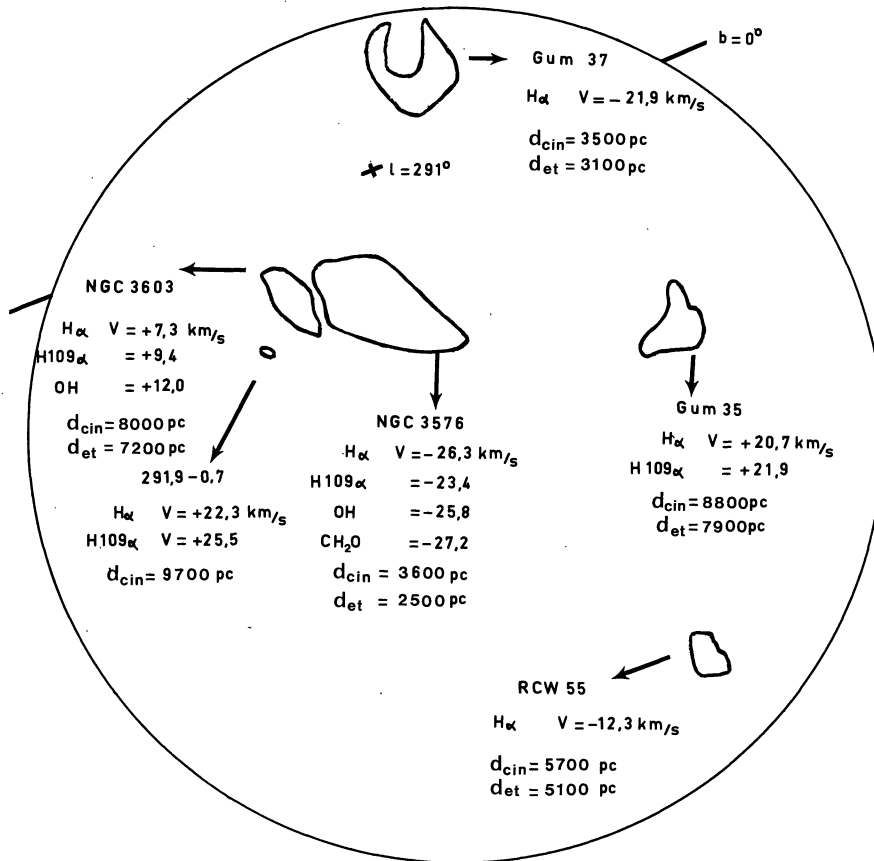


Fig. 5. Identification chart of H $\alpha$  regions near  $\ell = 291^\circ$ , showing radial velocities of H $\alpha$ , H 109 $\alpha$  and OH and CH $_2$ O absorption components. Stellar distances and kinematic distances are indicated too

optical and radio data. Only the radio sources 298.8–0.3 and 298.9–0.4 at 11.7 kpc, and 289.1–0.4 at 8.9 kpc, are not observed optically.

b) H $\alpha$  photography (Georgelin and Georgelin, 1970a, Fig. 24, p. 29) shows a weak H II region coinciding with the very distant radio source 316.8–0.1. This identification has been confirmed by Smith (1972). The observed radial velocities are given in Table 4. The mean velocity of  $-37 \text{ km s}^{-1}$  corresponds to distances of 2.1 and 12.5 kpc. The absorption profiles of Whiteoak and Gardner (1970) and Goss *et al.* (1972) show that this source is beyond the tangential point, i.e., at 12.5 kpc. The nebula observed optically at this location seems therefore the farthest known to date, and it would be interesting to be able to detect its exciting star.

c) Recently we have optically detected three H II regions in W 51 with H $\alpha$  monochromatic photographs, the first H $\alpha$  region corresponding to 49.0–0.3 and the two others to 49.5–0.4. From Fabry-Perot rings we have measured their radial velocities which agree with the H 109 $\alpha$  radial velocities. These results will be published later with more details (Georgelin, 1976).

d) We note also, among the very distant optical groups, the complex of H II regions at 8.8 kpc and at  $\ell = 70^\circ$ . This radio complex W 58 contains three H $\alpha$  regions cited by Minkowski (1946): Mi I.18 (NGC 6857=S 100),

Table 4. Radial velocities of 316.8–0.1

Component	$V_{LSR}$ (km s $^{-1}$ )	Reference
H 109 $\alpha$	-36.1	Wilson <i>et al.</i> (1970)
OH absorption	-38.0	Robinson <i>et al.</i> (1971)
CH $_2$ O	-37.7	Whiteoak and Gardner (1970)
	-41.8	
	-45.6	
	-45.6	
H I absorption	-38.1	Goss <i>et al.</i> (1972)
	-45.4	
	-52.9	

Mi I.16 (S 99) and Mi I.17 (70.4+1.6); these three regions are within a more extended H $\alpha$  region which has since been observed in the radio continuum by Israel (1974). We have found (Georgelin *et al.*, 1973) that the radial velocities of the three H $\alpha$  regions and of the surrounding region are the same, which shows that this is indeed a complex of regions at the kinematic distance of 8.7 kpc. The distance of the exciting star of NGC 6857 confirms this distance.

To the south of W 58 Dickel and Milne (1972) have detected, in H 109 $\alpha$ , a new H II region—69.9+1.5— with a radial velocity of  $-61 \text{ km s}^{-1}$ . Recently, in H $\alpha$ , we have measured the radial velocity of 69.9+1.5 as  $V_{LSR} = -35 \text{ km s}^{-1}$ . This result shows that 69.9+1.5

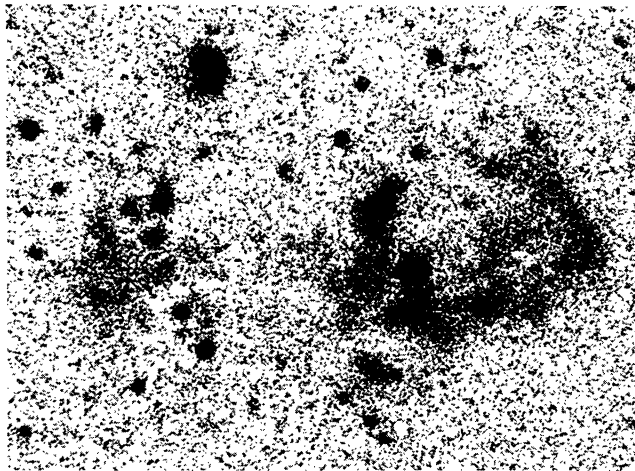
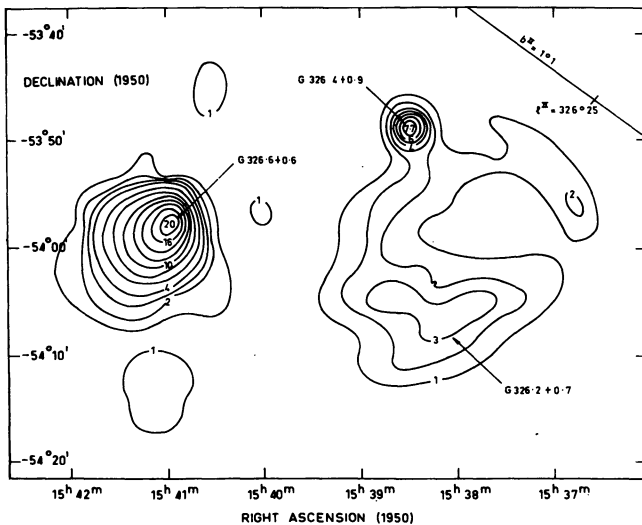


Fig. 6. (a) Radio continuum map of Shaver and Goss (1970). (b) On the same scale, H $\alpha$  photograph of the optical sources RCW 94, RCW 95 and 326.6+0.6

lies either at 8.8 kpc (the distance of W 58), or at 9.5 kpc ( $V_{\text{LSR}} = -35 \text{ km s}^{-1}$ ) which put it in the same spiral arm that W 58.

At the neighboring longitude  $\ell = 79^\circ$  is the H 109 $\alpha$  source 79.3+1.3 (DR 7). Reifenstein *et al.* (1970) have calculated the kinematic distance of this source (8.0 kpc) which certainly is a part of the same spiral arm as W 58 but is not detected optically.

e) Around the same longitudes, several optical and radio regions are located at distances of about 5 kpc, e.g., S 104 at  $\ell = 74^\circ 8'$  and distance 4.8 kpc, with matching H $\alpha$  and H 109 $\alpha$  velocities (Dickel and Milne, 1972); and the peculiar nebula S 106 at about 5.8 kpc, which shows strong expansion velocities (Maucherat, 1975).

Thus there now is a good overlap between the optical and radio data, which facilitates the interpretation of the radio results and permits the determination of a unique rotation model from H $\alpha$  and H 109 $\alpha$  velocities and stellar distances.

### 3. Optical Identification of Radio Sources

To trace the spiral structure of our Galaxy, it is indispensable to harmonize and to clarify all of the observational results. For a single well-defined H II region we can have the spectrophotometric distance of the exciting star, the H $\alpha$  kinematical distance, and the H 109 $\alpha$  kinematical distance (s) (near and far): of course we have compared these results in order to obtain only one distance for this region (examples § IV 2, 3). For a complex of H II regions, when there was convincing evidence (H $\alpha$  photograph, radio continuum map, morphology, matching radial velocities or distances, etc.) that the various features do belong together we have assigned a unique distance to this complex (examples § IV 3, V). For the whole Milky Way we have thus made optical identifications of these H 109 $\alpha$  sources, especially by using H $\alpha$  photographs of H II regions (Georgelin and Georgelin, 1970a), and radio continuum maps (Shaver and Goss, 1970).

We present below three examples of this analysis, for  $\ell = 326^\circ$ ,  $\ell = 328^\circ$  and  $\ell = 291^\circ$ .

#### a) H II Regions at $\ell = 326^\circ$

An H $\alpha$  photograph (Georgelin and Georgelin, 1970a) shows two previously known H II regions, RCW 94 and RCW 95, and two new H II regions that we have detected. These four nebulae are connected by very weak emission. The optical radial velocities  $C = -43.8, -38.6, -41.3$  and  $-39.5 \text{ km s}^{-1}$  confirms the physical connection.

The continuum radio map of Shaver and Goss (1970), reproduced in Fig. 6 along with an enlargement, at the same scale, of part of the H $\alpha$  photograph, shows that the radio source 326.5+0.9, also detected in H 109 $\alpha$ , is associated with the optical nebula RCW 94. The radio contours correspond rather well with the optical form, particularly in the south and in the west. The radial velocities are  $-41.3$  and  $-39.0 \text{ km s}^{-1}$  in H $\alpha$  and in H 109 $\alpha$ , respectively.

The same continuum map shows also that the radio source 326.6+0.6 coincides with one of the two new nebulae discovered in this photograph and not with RCW 95. The observed radial velocities are given in Table 5.

We therefore conclude that there is a single complex of H II regions comprising four optical nebulae, of which two were detected by radio (H 109 $\alpha$ , OH and CH $_2$ O). The mean of the H $\alpha$  and H 109 $\alpha$  observed velocities,  $-40.6 \text{ km s}^{-1}$ , gives the kinematic distance—2.5 kpc—of this unique complex.

#### b) H II Regions at $\ell = 328^\circ$

The H $\alpha$  photograph of this field (Georgelin and Georgelin, 1970a, Fig. 32, p. 37) shows an annular region of diameter  $3^\circ$  with a number of condensations including RCW 96, RCW 97, RCW 98, RCW 99 and 328.1-0.5.

The association of these various nebulae with the extended region is confirmed by their radial velocities which are, respectively,  $-51.5$ ,  $-49.2$ ,  $-40.1$ ,  $-47.8$ ,  $-42.0$  and  $-47.7$   $\text{km s}^{-1}$ .

The radio continuum map of Shaver and Goss (1970) shows that the nebula RCW 97 coincides with an extension towards the southwest of the radio source 327.3–0.5. The various radial velocities are in excellent accord, as can be seen from Table 6.

The radio source 328.0–0.6 is close to the optical region 328.1–0.5. The radial velocities agree:  $-42.0$   $\text{km s}^{-1}$  for  $\text{H}\alpha$  (Georgelin, 1975) and  $-42.9$   $\text{km s}^{-1}$  for  $\text{H 109}\alpha$  (Caswell, 1972).

The radio source 328.0–0.1, observed in  $\text{H 109}\alpha$ , has no optical counterpart but its velocity of  $-44.7$   $\text{km s}^{-1}$  agrees well with that of this complex with which it certainly is associated.

We therefore conclude that, morphologically and kinematically, all these regions form a single complex of mean radial velocity  $-46.2$   $\text{km s}^{-1}$ , corresponding to a kinematic distance of 2.9 kpc. This is in agreement with the distance of the exciting stars—2.6 kpc—and with the concentrations of OB stars at 2.5 kpc observed by Muzzio and McCarthy (1973).

### c) H II Regions at $\ell = 291^\circ$

The  $\text{H}\alpha$  and  $\text{H 109}\alpha$  velocities, as well as the velocities for diverse molecules observed in absorption, are indicated in Fig. 5 for each H II region. For each region the agreement between the various velocities is excellent, and therefore we are sure that these H II regions are at different distances. In addition, these kinematic distances, ranging from 2.5 kpc to 9.5 kpc, agree with the independently determined stellar distances (Georgelin and Georgelin, 1970). This is therefore a good example of convergence between all of the optical and radio data; it shows the spread in distance of the Carina arm which is aligned along the line of sight at this longitude.

We have made these identifications and correlations between all of the optical and radio results for the galactic plane in order to have homogeneous observational material available for the tracing of the spiral structure.

## V. Regions at Longitudes $332^\circ$ and $337^\circ$

This zone of longitude is crucial<sup>1)</sup> for tracing the spiral structure. For one thing, the line of sight intersects the spiral arms out to great distances and over the entire range of galactocentric distance from 4 kpc to 12 kpc; also the differential rotation is large and there is there-

<sup>1)</sup> The symmetrical region at  $\ell = 23^\circ$ – $28^\circ$  is less crucial, because for trailing spirals, we intersect—at near distance,  $r < 10$  kpc—more spiral arms and in a more perpendicular way at  $\ell = 330^\circ$  than at  $\ell = 30^\circ$ .

Table 5. Radial velocities of 326.6+0.6

Component	$V_{\text{LSR}}$ ( $\text{km s}^{-1}$ )	Reference
$\text{H}\alpha$	$-43.8$	Georgelin (1975)
$\text{H 109}\alpha$	$-44.5$	Wilson <i>et al.</i> (1970)
OH absorption	$-44.1$	Goss <i>et al.</i> (1970)
	$-21.2$	
$\text{CH}_2\text{O}$	$-45$	Robinson <i>et al.</i> (1971)
	$-21$	
	$-43.6$	
$\text{CH}_2\text{O}$	$-22.2$	Whiteoak and Gardner (1970)

Table 6. Radial velocities of 327.3–0.5

Component	$V_{\text{LSR}}$ ( $\text{km s}^{-1}$ )	Reference
$\text{H}\alpha$	$-49.2$	Georgelin and Georgelin (1970)
$\text{H 109}\alpha$	$-48.8$	Wilson <i>et al.</i> (1970)
H I absorption	$-49.5$	Goss <i>et al.</i> (1972)
OH absorption	$-49.0$	Robinson <i>et al.</i> (1971)
$\text{CH}_2\text{O}$	$-50.8$	Whiteoak and Gardner (1970)

fore a wide spread of radial velocities, allowing one to discriminate between various velocity groups—i.e., spiral arms. Observation of H II regions at  $\text{H 109}\alpha$  is very fruitful since fifteen and seventeen regions are detected at  $\ell = 332^\circ$  and  $337^\circ$  respectively. Clearly, a general interpretation of the spiral structure of our Galaxy first requires a satisfactory interpretation of these thirty-two regions (which have 64 possible kinematic distances).

### 1. H 109 $\alpha$ Regions at $\ell = 332^\circ$

Between  $\ell = 330^\circ.9$  and  $\ell = 333^\circ.7$  are fifteen H 109 $\alpha$  regions, whose velocities range from  $-47.0$   $\text{km s}^{-1}$  to  $-90.8$   $\text{km s}^{-1}$ . When we plot the kinematic distances this gives thirty possible points along the line of sight. We have already seen that a number of H 109 $\alpha$  sources thus detected may be the intensity maxima of a single region; this is particularly true here, in this field which is very rich in H II regions (as well as hot stars), as is shown in the  $\text{H}\alpha$  photograph on which are also indicated the H 109 $\alpha$  sources (Georgelin, 1970, Fig. 2, p. 325).

Precise optical identification shows that

- the two H 109 $\alpha$  sources—332.7–0.6 and 332.8–0.6, with velocities of  $-47.0$   $\text{km s}^{-1}$  and  $-57.2$   $\text{km s}^{-1}$  respectively—correspond exactly to two optical zones of similar velocity in the same H II region, RCW 106 ( $-47.0$  and  $-58.3$   $\text{km s}^{-1}$  respectively);
- the two H 109 $\alpha$  sources 333.0–0.4 and 333.1–0.4, with velocities of  $-53.8$  and  $-55.8$   $\text{km s}^{-1}$ , situated on the edge of the optical H II region 333.1–0.6 and having the same radial velocity, correspond to this same single H II region which is, moreover, at the same distance as RCW 106; and
- similarly, the sources 332.2–0.5 and 333.3–0.4, of velocities  $-55.0$  and  $-50.1$   $\text{km s}^{-1}$ , have weaker optical

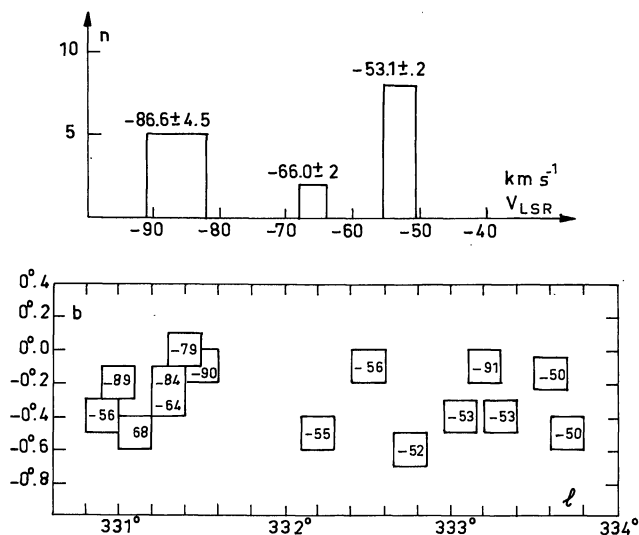


Fig. 7. (a) H  $109\alpha$  radial velocities of fifteen sources near  $\ell=332^\circ$ . Histogram of the number of sources as a function of radial velocity. (b) Grouping of these fifteen H  $109\alpha$  sources as a function of longitude and latitude

and  $-86 \pm 4.5 \text{ km s}^{-1}$  which correspond to three different distances. Figure 7b shows the positions of these regions in the sky ( $\ell, b$ ), along with their velocities. We note that:

- The five sources (group 1) with the greatest velocities ( $-86.6 \text{ km s}^{-1}$ ) are concentrated at very low latitudes ( $-0^\circ 2 < b < 0^\circ$ ). None of them is detected optically. Their distance is 7.8 or 9.7 kpc.
- The eight H  $109\alpha$  sources (group 3) of velocities  $-50$  to  $-56 \text{ km s}^{-1}$  constitute a single “near” complex (4.2 kpc) of which a large part is detected optically (six H II regions and numerous OB stars).
- The two H  $109\alpha$  sources (group 2) of velocities  $-64$  and  $-68 \text{ km s}^{-1}$  are closely neighboring H II regions and constitute an optically non-detected group which cannot be integrated with the two preceding groups.

## 2. H $109\alpha$ Regions at $\ell=337^\circ$

Between  $\ell=335^\circ 8$  and  $\ell=338^\circ 4$  are seventeen H  $109\alpha$  regions of velocities between  $-4$  and  $-93 \text{ km s}^{-1}$ .

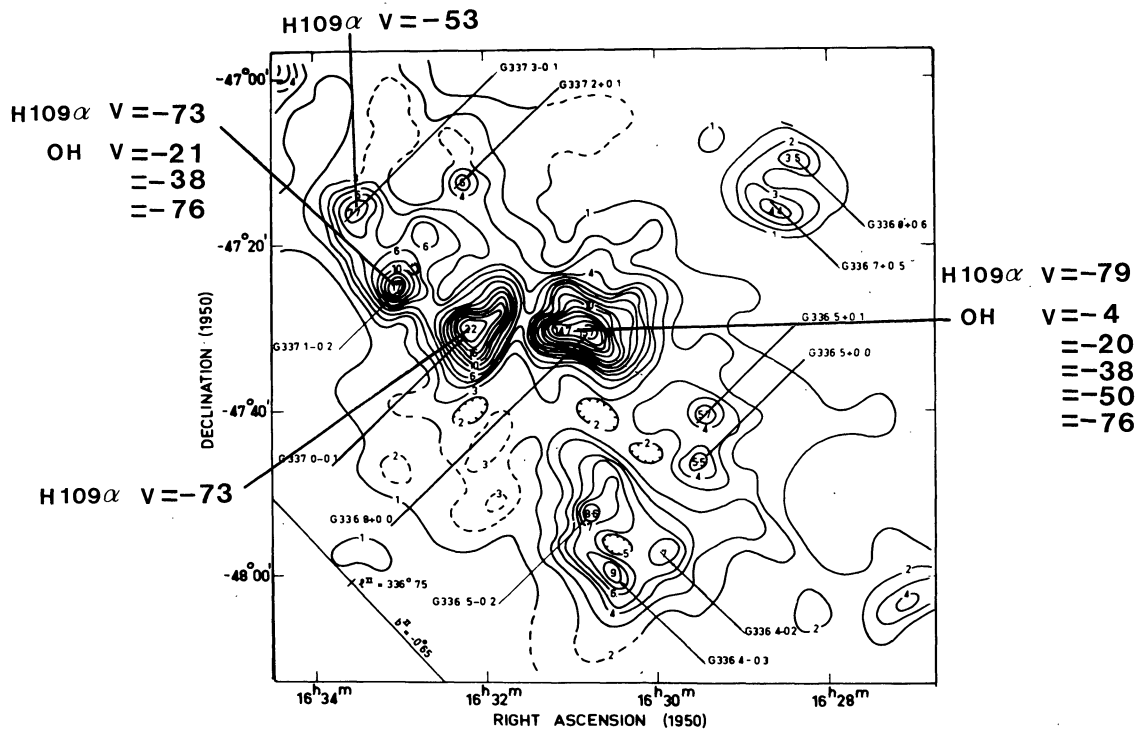


Fig. 8. 408 MHz continuum maps of Shaver and Goss (1970) at longitude  $336^\circ$ . Radial velocities (in  $\text{km s}^{-1}$ ) of group 4 (No. 89 in Table 9)

counterparts. We may conclude that these six H  $109\alpha$  sources are certainly part of the same complex as the six optical H II regions of this field.

We have grouped  $332.7-0.6$  with  $332.8-0.6$ ,  $333.0-0.4$  with  $333.1-0.4$ , and  $333.6-0.1$  with  $333.6-0.2$ . We have taken the mean of the radial velocities from H $\alpha$  and H  $109\alpha$ , in agreement with other constituents (three CH<sub>2</sub>O sources, two H<sub>2</sub>O sources and one OH absorption region). In Fig. 7a the mean velocities are assembled in three groups centered on  $-53 \pm 2$ ,  $-66 \pm 2$

The optical H II regions identified in this sector of longitude are:

- RCW 107, comprised of the two small symmetric nebulae NGC 6164 and NGC 6165 which are in expansion around the Of star HD 148937. This nebula has been studied in detail by Fabry-Perot interferometry by Pismis (1974). The identification of NGC 6164 with the H  $109\alpha$  source 336.4-0.2 is sure; the identification of NGC 6165 with 336.4-0.3 is a little less sure but the distance of 7.3 s of arc between the two radio sources



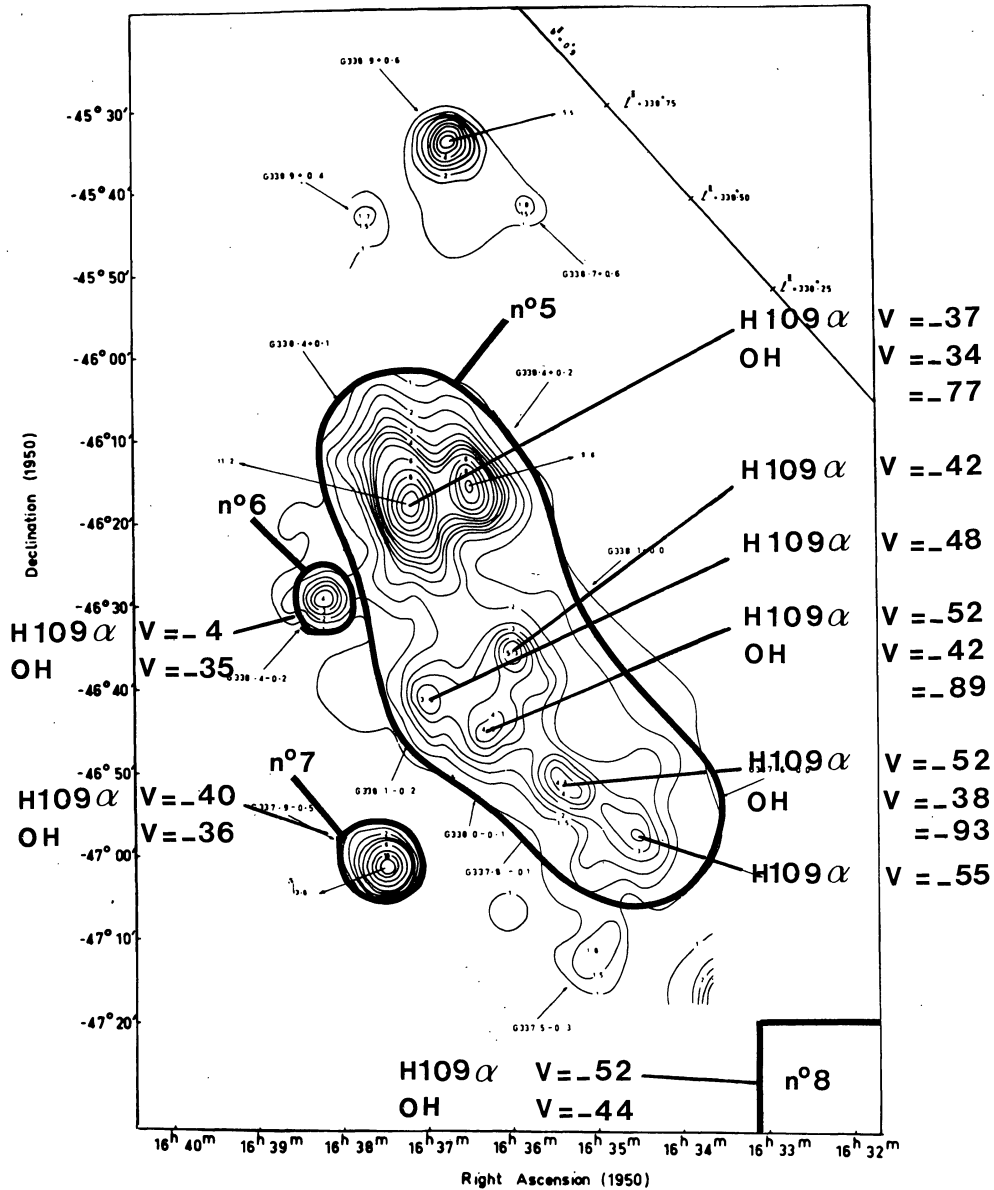


Fig. 9. 5000 MHz continuum maps of Shaver and Goss (1970) at longitude 338°. Identifications and radial velocities (in  $\text{km s}^{-1}$ ) of groups 5, 6, 7 and 8 (Nos. 91, 92, 90 and 87 in Table 9)

coincides well with the distance between the two principal optical condensations, although the orientation is slightly different. These two radio sources have radial velocities of  $-68$  and  $-93 \text{ km s}^{-1}$  respectively, with a greater than normal "turbulence". Optically too, one observes (Pismis, 1974) very large internal motions—from  $+20 \text{ km s}^{-1}$  to  $-120 \text{ km s}^{-1}$ —with mean velocities of  $-20$  and  $-80 \text{ km s}^{-1}$  for the two condensations. Analysis of the radio continuum map of Shaver and Goss (1970) shown in Fig. 8 suggests that the two other closely neighboring sources,  $336.5+0.0$  and  $336.5-0.2$ , with similar radial velocities ( $-63$  and  $-88 \text{ km s}^{-1}$  respectively), should be grouped with this very peculiar complex which we shall not use for tracing the spiral structure.

— RCW 108, a very large (diameter  $3^\circ$ ) and brilliant optical H II region of mean velocity  $-19.9 \text{ km s}^{-1}$  which corresponds to a single H 109 $\alpha$  source,  $336.5-1.5$ , of velocity  $-24.9 \text{ km s}^{-1}$  (group 9).

Analysis of the radio continuum maps of Shaver and Goss (1970) reproduced in Figs. 8 and 9 shows that:

— The three sources  $336.8+0.0$ ,  $336.9-0.1$  and  $337.1-0.2$  constitute a single complex (group 4, Fig. 8), which is confirmed by the H 109 $\alpha$  velocities, respectively  $-78.9$ ,  $-73.1$  and  $-72.7 \pm 1.8 \text{ km s}^{-1}$ .

The clearly weaker source  $337.3-0.1$  seems linked to this complex in spite of a different, but more uncertain, radial velocity:  $-53.5 \pm 6.6 \text{ km s}^{-1}$ . From this complex's mean velocity we obtain a distance of 11.7 or 6.7 kpc.

— The neighboring sources  $337.6-0.0$ ,  $337.8-0.1$ ,

Table 7. H II Region Complexes at  $\ell = 332^\circ$  and  $\ell = 337^\circ$ 

Group	$\ell$	Mean $V_{\text{LSR}}$ ( $\text{km s}^{-1}$ )	Number of regions	Distances (kpc)	Numero (Table 9)
1	331 $^\circ$ 7	-86.6	5	7.8 or 9.8	85
2	331 $^\circ$ 2	-66.0	2	5.4 or 12.1	84
3	332 $^\circ$	-53.1	8	4.2 or 13.4	86
4	336 $^\circ$ 9	-74.9	3	6.7 or 11.7	89
5	338 $^\circ$	-46.9	5	4.4 or 14.1	91
6	338 $^\circ$ 4	-4.3	1	0.6 or 18.0	92
7	337 $^\circ$ 9	-40.4	1	3.7 or 14.8	90
8	335 $^\circ$ 8	-52.1	1	4.6 or 13.7	87
9	336 $^\circ$ 5	-19.9	1	1.3	88

Table 8. Observed and calculated radial velocities

Velocities observed at 332 $^\circ$ (groups 1 and 2)	-88.6 $\text{km s}^{-1}$ (group 1)	-66.0 $\text{km s}^{-1}$ (group 2)
Predicted velocities at 337 $^\circ$ assuming circular arms	-71.7 (group 1)	-51.9 (group 2)
Predicted velocities at 337 $^\circ$ assuming spiral arms inclined at 12 $^\circ$ , for the case of near solution	-77.7 (group 1)	-53.9 (group 2)
Predicted velocities at 337 $^\circ$ assuming spiral arms inclined at 12 $^\circ$ , for the case of the distant solution	-63.6 (group 1)	-46.3 (group 2)
Velocities observed at 337 $^\circ$ (groups 4 and 5)	-74.9 (group 4)	-46.9 (group 5)

338.1-0.2 and 338.1+0.0 are connected together and to the two intense sources 338.4+0.2 and 338.4+0.1; they constitute a complex (group 5, Fig. 9) spread out along the galactic plane. The H 109 $\alpha$  velocities (-54.8, -52, -52.5, -47.7, -41.7 and -38  $\text{km s}^{-1}$ ) are in agreement, and the mean velocity of -46.9  $\text{km s}^{-1}$  gives a distance of 4.4 or 14.1 kpc for this complex.

— The source 338.4-0.2 (group 6, Fig. 9) is not really part of the complex; this is confirmed by a very different velocity—-4.3  $\text{km s}^{-1}$ —corresponding to a distance of 0.6 or 18.0 kpc.

— The two sources 335.8-0.2 (group 8, Fig. 9) and 337.9-0.5 (group 7, Fig. 9) are not connected with the preceding complexes.

### 3. Interpretation

After examining the various optical and radio data, the thirty-two H 109 $\alpha$  regions were separated into the nine groups of Table 7.

We have seen in Fig. 7 that at longitude 332 $^\circ$ , groups 1, 2 and 3 have velocities of -86.6, -66.0 and -53.1  $\text{km s}^{-1}$ . By simple continuity of the spiral arms one should find these same velocities slightly shifted at a nearby longitudes. The predicted velocities of these regions at longitude 337 $^\circ$  are given in Table 8.

This table shows:

- that there is velocity continuity between groups 1 and 4,
- that there is velocity continuity between groups 2 and 5,
- that the continuity between 1 and 4 is much better in the case of the “near” distance,
- that the continuity between 2 and 5 is much better in the case of the “far” distance.

#### Choice of Distances (Fig. 10)

*Group 3 (No. 86).* Six of the eight sources of group 3 are detected optically and are situated at 4.2 kpc. The sources detected in absorption have no higher velocity component. There is no doubt concerning the “near” distance for group 3.

*Group 9 (No. 88).* Group 9 is identified with RCW 108, which is part of the Sagittarius-Carina arm.

*Group 5 (No. 91).* For group 5 there is no question that the distance solution at 14.1 kpc is correct, because there are three OH sources and one H I component in absorption at greater absolute velocities (-77, -89, -93 and -60  $\text{km s}^{-1}$ ). The very low latitude (0 $^\circ$ 0, 0 $^\circ$ 0, -0 $^\circ$ 1 and -0 $^\circ$ 2), and the shape of the complex, spread out along the galactic plane, does not contradict this hypothesis.

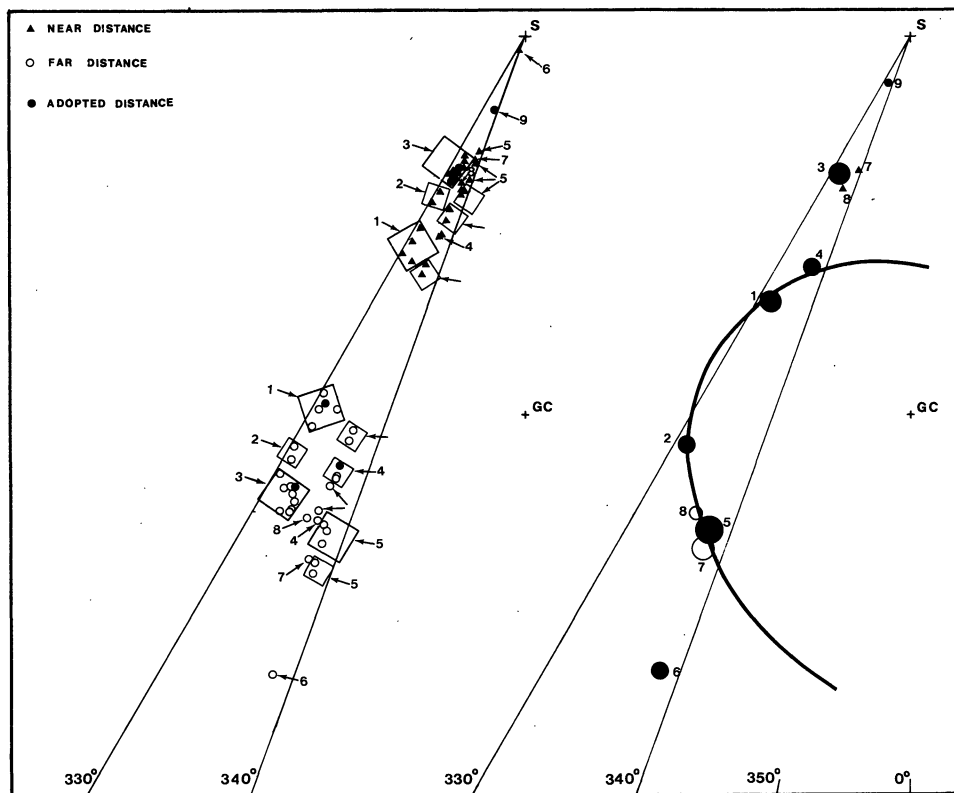


Fig. 10. Grouping and interpretation of the thirty-two H 109 $\alpha$  regions between  $\ell = 330^\circ$  and  $\ell = 340^\circ$

*Group 6 (No. 92).* If group 6 were at 0.6 kpc it would be detected optically. In addition, Caswell and Robinson (1974) have found OH in absorption at  $-35 \text{ km s}^{-1}$ . The “far” distance at 18 kpc is therefore adopted.

*Group 4 (No. 89).* In front of two of the sources of group 4, Caswell and Robinson (1974) have detected OH in absorption, but none of the radial velocity components is more negative than that of H 109 $\alpha$ , even though OH absorption components with radial velocities up to the subcentral velocity have been observed at the same longitude (e.g., in front of group 5). Thus group 4 is at the “near” distance.

*Group 2 (No. 84).* We have already seen that group 2 is associated (Table 8) with group 5, which is distant. Therefore we chose the distant solution for group 2.

*Group 1 (No. 85).* There is a continuity of velocity between group 1 and the “near” group 4. Furthermore, the hypothesis of the “near” solution places these two groups on the same spiral as groups 2 and 5.

*Groups 7 and 8 (Nos. 90 and 87).* For groups 7 and 8 the choice is difficult. We adopt the “near” distance without strong conviction; OH absorption components have been detected (Caswell and Robinson, 1974) but the H 109 $\alpha$  velocity is slightly offset from the OH velocities.

Figure 10 groups the results of all the § V.

We have made the same discussion for each complex in the whole galactic plane.

In Table 9 we give the data: running number, mean longitude, adopted  $V_{\text{LSR}}$ , spectrophotometric or kinematic distance, excitation parameter (see § VI, 1) and exciting stars for each complex with  $U > 70 \text{ pc cm}^{-2}$ ; H 109 $\alpha$  and H $\alpha$  identification and their  $V_{\text{LSR}}$  for each component.

In Table 10 we give the criteria used (§ II, 4) to resolve distance ambiguity for each complex.

## VI. Outline of the Large-Scale Spiral Structure

### 1. Description of the Ionized Hydrogen Spirals Arms (Fig. 11)

Figure 11 has been made from the pertinent data (longitude, radial velocity, spectrophotometric or kinematic distance, choice of “near” or “far” distance) obtained for each complex of H II region. Details are in Tables 9 and 10.

In Fig. 11 we show only the “important” H II regions<sup>2)</sup> ( $U > 70 \text{ pc cm}^{-2}$ ), and we have “weighted” each of them according to its excitation parameter  $U$ : bright (b) for  $U > 200 \text{ pc cm}^{-2}$ , medium (m) for  $110 < U < 200 \text{ pc cm}^{-2}$ , faint (f) for  $70 < U < 110 \text{ pc cm}^{-2}$ . The main effect of this weighting is to show the relative insignificance of our local region of the Galaxy, which does not

<sup>2)</sup> The “giant” H II regions of Mezger (1970) are defined by  $\delta D^2 > 400 \text{ f.u. kpc}^2$ , corresponding to  $U > 100 \text{ pc cm}^{-2}$ .

Table 9

No	$l$	adopted $V_{LSR}$ (km s <sup>-1</sup> )	$d_{star}$ (kpc)	$d_{kin}$ (kpc)	$U$ (pc cm <sup>-1</sup> )	H109 $\alpha$ source	H $\alpha$ source	$V_{LSR}$ H109 $\alpha$ (km s <sup>-1</sup> )	$V_{LSR}$ H $\alpha$ (km s <sup>-1</sup> )
1	6°0	+ 4.4	1.5		f	6.0 - 1.2	M 8	+ 3.0	+ 5.7
cluster NGC 6530									
2	7°0	+13.8	2.7		f	7.0 - 0.2	M 20	+16.4	+11.3
HD 164492									
3	8°5	+35	5.3		f	8.5 - 0.3		+35	
4	10°3	+ 8.0	18.7	b		10.2 - 0.3 10.3 - 0.1 10.6 - 0.4		+13.9 + 9.7 + 0.3	
5	12°8	+36.3	4.2	m		12.8 - 0.2		+36.3	
6	13°2	+57.0	5.8	f		13.2 + 0.0		+57.0	
7	14°6	+37.2	4.0	f		14.6 + 0.1		+37.2	
8	15°1	+18.1	2.5	m		15.1 - 0.7	M 17	+17.2	+19.0
cluster NGC 6618									
9	17°0	+26.4	2.0	f		17.0 + 0.8	M 16	+24.5	+28.4
cluster NGC 6611									
10	18°2	+47.4	4.3	f			S 53		+47.4
11	18°6	+31.1	2.2	m		18.5 + 1.9	S 54	+32.9	+29.4
BD -11°4581 BD -12°4964 BD -12°4970 BD -12°4973 BD -12°4975 BD -12°4979 HD 167971 BD -12°4982 BD -12°4984 HD 168112 BD -12°4994 HD 168206									
12	19°1	+67.3	13.4	m		19.1 - 0.3		+67.3	
13	19°7	+43.4	3.8	f		19.7 - 0.2		+43.4	
14	20°7	+57.4	14.0	m		20.7 - 0.1		+57.4	
15	22°8	+82.5	12.4	b		22.8 - 0.3		+82.5	
16	23°4	+101.5	7.3	m		23.4 - 0.2		+101.5	
17	24°6	+113.1	9.1	m		24.6 - 0.2 24.5 + 0.2 24.8 + 0.1		+109.2 +116.0 +114.1	
18	25°4	+ 60.4	13.6	b		25.4 - 0.2		+ 60.4	
19	25°8	+110.0	9.0	m		25.8 + 0.2		+110.0	
20	28°0	+ 96.9	6.9	m		*27.3 - 0.2 28.6 + 0.0		+ 97.6 + 96.2	
21	29°9	+ 96.4	7.0	m		29.9 - 0.0 (30.2 - 0.2)		+ 96.4 (+101.0)	
22	30°8	+ 92.5	6.6	b		30.8 - 0.0 (31.1 + 0.0)		+ 92.5 (+ 99.2)	
23	34°3	+ 53.9	3.6	f		34.3 + 0.1		+ 53.9	
24	35°2	+ 46.5	3.0	f		35.2 - 1.7		+ 46.5	
25	36°4	+ 57.4	3.9	f			S 72	+ 57.4	
26	37°6	+ 51.8	12.4	b		37.4 - 0.2 37.6 - 0.1 37.9 - 0.4		+ 39.0 + 55.8 { + 60.2 + 61.0	
27	43°2	+ 8.6	14.1	b		43.2 - 0.0		+ 8.6	
28	45°5	+ 55.5	3.9	f				45.5 + 0.0 45.5 + 0.1	+ 57.4 { + 53.0 + 54.0
29	48°6	+ 17.0	11.9	m		48.6 + 0.0			+ 17.0
30	49°2	+63.5	5.8	b		48.9 - 0.3 49.0 - 0.3 49.1 - 0.4 49.2 - 0.3 49.4 - 0.3 49.5 - 0.4		[49.0 - 0.3] [49.5 - 0.4]	+66.4 { +63.2 +65.8 +72.4 { +66.1 +67.2 +52.8 { +59.0 +58.2
31	-61°3	+52.6	4.4	m		51.2 - 0.1 51.4 - 0.0			+55.3 +49.8
32	70°3	-23.4	8.8	b		70.3 + 1.6 70.3 + 1.6	NGC6857-8100 K 3.50 S 99	-24.4 -24.2	-21.7 -22.0 -24.4 (-35)
An. $\alpha_{1950} = 19^h39^m9.6$ , $\delta_{1950} = 33^{\circ}23'$ (Liller and Chang Yuan Shao, 1968 Crampton and Georgelin, 1975)									
33	75°8	- 4.8	5.7	f		75.8 + 0.4			- 4.8
34	76°4	- 7.6	5.8	f		76.4 - 0.6	S 106		+ 5.6 - 7.6
35	79°3	-41.4	7.6	m		79.3 + 1.3			-41.4
36	79°9	-18.7	5.7	m		79.8 + 1.2 80.0 + 1.5			-22 -15.4
37	80°0	(x)	1.4	m		Cygnus X complex	S 109		
38	84°9	+ 2.1	2.0	f			S 115		+ 2.1
BD 46°2972 BD 46°2978 LSIII 46°11 LSIII 46°12									
39	85°5	+ 3.0	0.8	f		84.9 - 0.7 (158 $\alpha$ ) (166 $\alpha$ )	NGC7000-8117	- 0.4 + 4.2	+ 3.0
HD 199579									
40	102°8	-50.9	3.6	m			S 132		-50.9
LSIII 55°32 HD 211853 BD 55°2722									
(x) The kinematic distances have little significance for Cygnus X sources.									
41	111°3	-37.7	2.5	f			S 157		-37.7
HD 219286 BD 59°2677 LSIII 59°57 HD 219460 LSIII 59°66 Cluster An. Rosal 3									
42	111°5	-60.1	2.8 (5.8)	m			S 158	-60.6	-59.6
Astrographic Cat. Vatican n° 37622 An. at 1' in SW									
43	133°7	-42.7	2.3	m		133.7 + 1.2 134.8 + 1.0	IC 1795 IC 1805	-42.3 -41.7	-40.7 -46.0
BD 61°411 Cluster IC 1805									
44	137°5	-36.6	2.3	f		137.3 + 1.0	IC 1848	-37.0	-36.6
Cluster IC 1848									
45	150°6	-23.0	3.6	f		150.6 - 0.9	S 206	-22.0	-24.0
BD 50°886									
46	151°6	-46.6	> 4.8	f		151.6 - 0.2	S 209	-46.1	-47.1
47	173°5	- 0.9	3.4	m		173.4 - 1.8	IC 410	- 1.8	+ 0.0
Cluster NGC 1893									
48	206°3	+19.6	1.5	f		206.4 - 1.9	S 275	+17.8	+21.5
Cluster NGC 2244									
49	209°0	+ 3.2	0.4	f		206.5 - 16.4 208.9 - 19.3 209.0 - 19.4	NGC 2024 NGC 1976	{ + 7.0 + 6.0 + 7.7 - 2.8	{ - 3.5 + 8.7 + 1.2
Association Ori OB1									
50	227°7	+56.3	5.2	f			S 298		+56.3
HD 56925									
51	243°3	+46.6	4.1	f		243.3 + 0.6	S 311		+46.6
Cluster NGC 2467									

















