THE X-RAY LUMINOSITY FUNCTION AND GAS MASS FUNCTION FOR OPTICALLY SELECTED POOR AND RICH CLUSTERS OF GALAXIES

JACK O. BURNS, MICHAEL J. LEDLOW, CHRIS LOKEN, AND ANATOLY KLYPIN

Department of Astronomy, New Mexico State University, Las Cruces, NM 88003-0001; jburns@nmsu.edu, mledlow@nmsu.edu, cloken@nmsu.edu, aklypin@nmsu.edu

Wolfgang Voges

Max-Planck-Institut für Extraterrestrische Physik, Postfach 1603, D-85740, Garching bei München, Germany; whv@mpe.dnet.nasa.gov

GREG L. BRYAN AND MICHAEL L. NORMAN

Department of Astronomy and National Center for Supercomputing Applications, 5600 Beckman Institute, Drawer 25, University of Illinois at Urbana-Champaign, Urbana IL 61801; gbryan@ncsa.uiuc.edu, norman@ncsa.uiuc.edu

AND

RICHARD A. WHITE

Space Data and Computing Division, Code 932, NASA/Goddard Space Flight Center, Greenbelt, MD 20771; rwhite@jansky.gsfc.nasa.gov
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ABSTRACT

We present the first X-ray luminosity function (XLF) for an optically selected sample of 49 nearby poor clusters of galaxies and a sample of 67 Abell clusters with $z \le 0.15$. We have extended the measured cluster XLF by more than a factor of 10 in X-ray luminosity. Our poor cluster sample was drawn from an optical catalog of groups with $0.01 \le z \le 0.03$ composed of Zwicky galaxies. The X-ray emission was measured from the *ROSAT* all-sky survey. About 45% of the poor clusters were detected, with 0.5–2.0 keV luminosities from 1.7 to $65 \times 10^{41} \ h^{-2}$ ergs s⁻¹. These are among the X-ray brightest, optically selected poor clusters in the northern hemisphere. For this sample, the poor cluster XLF was found to be a smooth extrapolation of the rich cluster XLF. A new Hydro/N-body simulation of a hot + cold dark matter model with $\Omega_{\text{total}} = 1$, $\Omega_{\nu} = 0.2$, and a baryon fraction of 7.5% was used to model and understand our observational selection effects. We found that the observed cluster gas mass function was consistent with our model.

Subject headings: galaxies: clusters: general — intergalactic medium — X-rays: galaxies

1. INTRODUCTION

The presence of head-tail radio galaxies (e.g., Venkatesan et al. 1994) and extended X-ray emission (e.g., Ebeling, Voges, & Böhringer 1994; Doe et al. 1995; Pildis, Bregman, & Evrard 1995) demonstrates the existence of a relatively dense $(\approx 5 \times 10^{-4} \text{ cm}^{-3})$ and hot $(\approx 1 \text{ keV})$ intracluster medium within poor clusters. The X-ray emission, in particular, is a good tracer of the gravitational potential well and the dynamical state of the clusters. Einstein observations (Price et al. 1991) showed that the bulk X-ray and optical properties of these poor clusters appear to be scaled-down versions of the rich clusters, a result consistent with ROSAT observations (Doe et al. 1995). However, recent work (e.g., Mulchaey et al. 1996), including new Advanced Satellite for Astronomy and Astrophysics (ASCA) results, suggests that poor clusters may have a wide range of baryonic fractions (5%-30%) and low metallicity (<0.15 solar, Davis et al. 1996; Tawara et al. 1995), which might indicate that they are distinctly different from rich clusters. We are only beginning to understand the formation and evolution of poor clusters; models of their genesis remain controversial (e.g., Diaferio et al. 1993).

In an effort to investigate the properties of poor galaxy clusters and to provide a reference point for studies of cluster evolution, we have constructed an X-ray luminosity function (XLF) for an optically selected, nearby, volume-limited sample, using images from the *ROSAT* all-sky survey (RASS). Previous XLFs have concentrated mainly on rich clusters (e.g., Edge et al. 1990; Henry et al. 1992; Castander et al. 1994),

although Henry et al. (1995) have recently added one low $L_{\rm X}$ point to the XLF from a small sample of X-ray selected groups. In this Letter, we also calculate the gas mass function for poor and rich clusters, and compare this function with that computed from large-scale structure models. We use $H_0=100~h~{\rm km~s^{-1}~Mpc^{-1}}$ and $q_0=0.5$ throughout.

2. POOR CLUSTER STATISTICAL SAMPLE & RASS IMAGES

Using an algorithm similar to that of Turner & Gott (1976), White et al. (1996) compiled a catalog of ≈600 optically selected poor clusters composed of Zwicky et al. (1961–68) galaxies down to 15^m.7. These groups were identified by drawing the largest possible circle, centered on a Zwicky galaxy, within which the surface density of Zwicky galaxies was more than 24 times the average surface density in the north galactic cap. This catalog contains nearly all the Yerkes poor clusters (e.g., White 1978), a few Hickson groups (4 out of 49 clusters in Table 1), as well as many more loose condensations. A partial listing of the densest poor clusters in the catalog along with VLA observations were published by Burns et al. (1987). For this X-ray project, we selected a complete subsample of the 49 densest poor clusters, as defined in Ledlow et al. (1996), each with (1) ≥ 4 Zwicky galaxies, (2) $|b| > 30^{\circ}$, (3) a surface overdensity $\gtrsim 50$, and (4) $0.01 \le z \le 0.03$. As a result, the effective radii of the groups in Table 1 range from 60 to 180 h^{-1} kpc for z = 0.01 to z = 0.03, respectively. Eight Abell clusters also lie within this volume (see Ledlow et al. 1996).

TABLE 1 POOR CLUSTER SAMPLE

$\begin{array}{cccccccccccccccccccccccccccccccccccc$	Cluster	z	$L_{\mathrm{x}}^{}}$	M _{gas} ^b
$\begin{array}{cccccccccccccccccccccccccccccccccccc$	S34-115	0.0225	13.0 ± 2.4	3.1 ± 0.5
S49-140 0.0179 9.0 ± 1.6 2.6 ± 0.4 S49-141 0.0264 <3.4	S49-147	0.0191	12.4 ± 2.1	3.1 ± 0.5
S49-141 0.0264 <3.4	S34-111	0.0173	16.9 ± 2.3	3.6 ± 0.5
S49-141 0.0264 <3.4	S49-140	0.0179	9.0 ± 1.6	2.6 ± 0.4
$\begin{array}{cccccccccccccccccccccccccccccccccccc$		0.0264	<3.4	< 1.6
$\begin{array}{cccccccccccccccccccccccccccccccccccc$	S49-145	0.0227	<3.7	<1.7
N45-342. 0.0165 <2.2	S49-142	0.0211	4.5 ± 1.5	1.8 ± 0.4
$\begin{array}{cccccccccccccccccccccccccccccccccccc$		0.0165	<2.2	<1.3
$\begin{array}{cccccccccccccccccccccccccccccccccccc$	N34-169	0.0292	18.9 ± 3.4	3.8 ± 0.6
$\begin{array}{cccccccccccccccccccccccccccccccccccc$	N45-388	0.0291	<4.1	<1.8
$\begin{array}{cccccccccccccccccccccccccccccccccccc$	N45-384	0.0266	7.0 ± 1.9	2.3 ± 0.4
$\begin{array}{cccccccccccccccccccccccccccccccccccc$	N34-170	0.0292	<2.3	<1.3
$\begin{array}{cccccccccccccccccccccccccccccccccccc$	N34-172	0.0207	<3.2	< 1.6
$\begin{array}{cccccccccccccccccccccccccccccccccccc$	N56-393	0.0221	6.4 ± 1.5	2.2 ± 0.4
$\begin{array}{cccccccccccccccccccccccccccccccccccc$		0.0264	15.0 ± 3.2	3.4 ± 0.6
$\begin{array}{cccccccccccccccccccccccccccccccccccc$	N67-322	0.0210	<2.2	<1.4
$\begin{array}{cccccccccccccccccccccccccccccccccccc$		0.0206	9.5 ± 1.9	2.7 ± 0.4
$\begin{array}{cccccccccccccccccccccccccccccccccccc$	N56-371	0.0276	<4.3	<1.8
$\begin{array}{cccccccccccccccccccccccccccccccccccc$	N67-300	0.0129	< 2.4	<1.3
$\begin{array}{cccccccccccccccccccccccccccccccccccc$	N56-392	0.0272	< 6.5	<2.2
$\begin{array}{cccccccccccccccccccccccccccccccccccc$		0.0153	< 2.2	<1.3
$\begin{array}{cccccccccccccccccccccccccccccccccccc$	N79-299B	0.0235	4.1 ± 1.4	1.8 ± 0.4
$\begin{array}{cccccccccccccccccccccccccccccccccccc$	N67-335	0.0204	49.5 ± 5.8	6.1 ± 0.9
$\begin{array}{cccccccccccccccccccccccccccccccccccc$	N79-299A	0.0235	14.4 ± 2.4	3.3 ± 0.5
$\begin{array}{cccccccccccccccccccccccccccccccccccc$	N79-282	0.0132	<1.7	<1.1
$\begin{array}{cccccccccccccccccccccccccccccccccccc$	N79-268	0.0253	<3.6	<1.6
$\begin{array}{cccccccccccccccccccccccccccccccccccc$	N79-283	0.0259	9.8 ± 2.1	2.7 ± 0.5
$\begin{array}{cccccccccccccccccccccccccccccccccccc$	N79-292	0.0235	<4.4	
$\begin{array}{cccccccccccccccccccccccccccccccccccc$	N79-284	0.0246	4.0 ± 1.5	1.7 ± 0.4
$\begin{array}{cccccccccccccccccccccccccccccccccccc$	N67-323	0.0289	<4.6	
$\begin{array}{cccccccccccccccccccccccccccccccccccc$		0.0217	< 6.2	<2.2
$\begin{array}{cccccccccccccccccccccccccccccccccccc$		0.0226		
$\begin{array}{cccccccccccccccccccccccccccccccccccc$	N79-296	0.0232	19.2 ± 3.8	3.8 ± 0.6
$\begin{array}{cccccccccccccccccccccccccccccccccccc$	N67-329	0.0231		
$\begin{array}{cccccccccccccccccccccccccccccccccccc$	N67-318	0.0235		
$\begin{array}{cccccccccccccccccccccccccccccccccccc$				
$\begin{array}{cccccccccccccccccccccccccccccccccccc$	N67-336	0.0196		
$\begin{array}{cccccccccccccccccccccccccccccccccccc$				
$\begin{array}{cccccccccccccccccccccccccccccccccccc$	N67-326			
$\begin{array}{cccccccccccccccccccccccccccccccccccc$				
$\begin{array}{cccccccccccccccccccccccccccccccccccc$				
$ \begin{array}{cccccccccccccccccccccccccccccccccccc$	N56-395			
$ \begin{array}{cccccccccccccccccccccccccccccccccccc$	N56-381			
N34-175 0.0283 64.5 ± 6.9 7.0 ± 1.0 N34-173 0.0266 16.5 ± 1.8 3.5 ± 0.5				
N34-173 0.0266 16.5 ± 1.8 3.5 ± 0.5				
	S49-146	0.0250	10.2 ± 2.2	2.8 ± 0.5
S49-144 0.0146 <2.0 <1.2	S49-144	0.0146	<2.0	<1.2

Note.—Positions can be found in Ledlow et al. 1996.

The fact that the volume density of our poor clusters is constant within this restricted redshift range (Ledlow et al. 1996) indicates that our sample is basically volume limited. Our apparent magnitude cutoff reduces the volume density of observable galaxies by a factor of ~5 across our volume, whereas the projected area within which the candidate group galaxies must lie increases by a factor of 9 because of the surface overdensity condition. Therefore, it appears plausible that these two effects cancel, at least to within a factor of 2, in

We cross-correlated this sample of poor clusters with the RASS (Voges 1992). ROSAT was nearly ideal for imaging groups since the PSPC detector was most sensitive to cooler (1 keV) clusters. From the RASS sky scans, we assembled

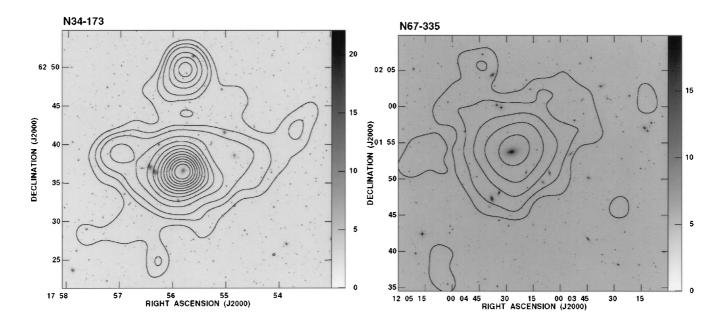
 $2^{\circ}.1 \times 2^{\circ}.1$ images from 0.5 to 2.0 keV with average exposure times of ≈550 s, corresponding to an X-ray flux limit of $\approx 4 \times 10^{-13}$ ergs cm⁻² s⁻¹. After exposure and background corrections, the images were smoothed with 204" FWHM Gaussians ($\approx 60 \ h^{-1}$ kpc at z = 0.02) to accentuate the extended intracluster mediums. Optical identifications of the X-ray sources were made by overlaying the RASS maps onto digital POSS-I images. Examples of RASS X-ray emission for four poor clusters are shown in Figure 1 (Plate L3).

For those clusters with X-ray detections, we summed the X-ray counts within a circle of linear radius $190 h^{-1}$ kpc about the centroid. We fit the surface brightness profiles of the seven best cluster detections with a $\beta = 2/3$ King model convolved with the smoothed point-spread function (PSF) determined from point sources in the fields. We measured an average $r_c = 65 h^{-1}$ kpc, which was used in correcting the flux within 190 h⁻¹ kpc to a total flux (following Briel & Henry 1993); this correction was 67%. The resulting RASS 0.5-2.0 keV luminosities along with 3 σ upper limits for undetected clusters are given in Table 1. We include a 10% calibration error in L_x added in quadrature with the Poisson error in the X-ray counts. The conversion between count rate and luminosity was performed using XSPEC, assuming a Raymond-Smith thermal spectrum with T = 1 keV, 0.3 solar abundance, and H_I column densities from Stark et al. (1992).

A total of 22/49 (45%) of the clusters in the sample were detected by the RASS. This X-ray detection rate is higher than that recently reported for other galaxy groups (e.g., Ebeling et al. 1994; Mulchaey et al. 1996). These detections represent some of the X-ray brightest, optically selected poor clusters in the northern sky. The X-ray luminosities span the range $(1.7-65) \times 10^{41} h^{-2}$ ergs s⁻¹, which run the gamut from that expected for bright, individual galaxies to richness class ≥0 Abell clusters. We do not include unresolved X-ray sources associated with individual galaxies as cluster detections. However, it is possible that some fraction of the total X-ray emission for these extended sources is due to emission from galaxies (e.g., active galactic nuclei, interstellar medium, or interacting galaxies; see e.g., Ebeling et al. 1994). From the $\log N$ - $\log S$ relationship for the RASS, we expect a total of only six random source projections within the search radii of these 49 clusters; but, this is an upper limit since we eliminated compact and noncluster sources using optical IDs.

Next, we ask how many of the poor clusters in our sample may be artificial systems produced by projection effects? Given the complex observational selection effects inherent in our sample, we addressed this question using a new N-body + Eulerian Hydro large-scale structure simulation (Bryan & Norman 1996; Loken et al. 1996). The simulation was done for a mixed Cold + Hot (two neutrino) dark matter (CHDM) model (Primack et al. 1995) in a $(50 h^{-1} \text{ Mpc})^3$ box, with a 512^3 mesh (98 h^{-1} kpc zone⁻¹), and 3×256^3 particles. The model assumed h = 0.5, $\Omega_{CDM} = 0.725$, $\Omega_{\nu} = 0.20$, $\Omega_b = 0.075$, and $Q_2 = 19 \mu K$, and was run on the CM-5 at NCSA. Local maxima in the cold dark matter distribution were identified as "galaxy halos," and periodic boundary conditions were used to replicate the galaxies into a larger volume (see Loken et al. 1996 for details). A vantage point was chosen for an observer, and a Schechter selection function (e.g. Klypin et al. 1990) was applied to statistically select a catalog of Zwicky-like galaxies. Then we identified groups in the same manner as discussed at the beginning of this section. We found ≈58 clusters (there is some variation with vantage

^a Units in h^{-2} 10⁴¹ ergs s⁻¹. ^b Units in $h^{-5/2}$ 10¹² M_{\odot} .



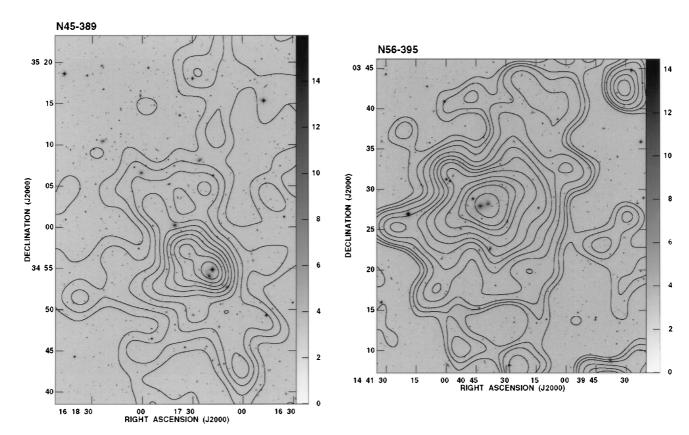


Fig. 1.—Overlays of RASS X-ray contours onto optical images from the digital POSS-I. Contour levels are as follows: 2, 3, 4, 6, 8, 10, 12, 14, 16, 18, 20, 22, $24 \times 6 \times 10^{-6}$ counts s⁻¹ (15" × 15" pixel)⁻¹ for N34-173; 2, 4, 8, 16, $32 \times 1.7 \times 10^{-5}$ counts s⁻¹ pixel⁻¹ for N67-335; 2, 4, 6, 8, 10, 12, 14, 16, 18, $20 \times 6.2 \times 10^{-6}$ counts s⁻¹ pixel⁻¹ for N45-389; 2, 3, 4, 6, 8, 10, 12, 15, 20, 25, 30, $40 \times 8.5 \times 10^{-6}$ counts s⁻¹ pixel⁻¹ for N56-395.

Burns et al. (see 467, L50)

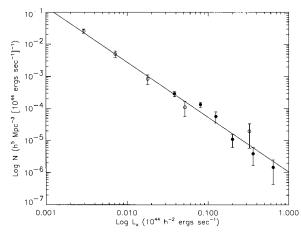


Fig. 2.—Differential cluster X-ray luminosity function. Open circles are from our statistical sample of poor clusters (also includes eight Abell clusters), with $0.01 \le z \le 0.03$. Filled circles are from a complete sample of Abell clusters, with $z \le 0.15$. Error bars are 1σ . The line is the best fit to a power law

point) in a volume with 0.01 < z < 0.03. Roughly 80%-86% of the groups had four or more galaxies within $2\ h^{-1}$ Mpc and virtually all (>96%) had at least two such members. Moreover, 88%-98% of the clusters selected from this two-dimensional percolation algorithm were spatially coincident with real clusters, which we had previously identified in the three-dimensional volume. We conclude that our poor clusters are nearly all likely to be real and not produced by projection effects.

3. THE CLUSTER X-RAY LUMINOSITY FUNCTION

The differential XLF for our poor clusters is represented by the open circles in Figure 2. Poisson error bars are plotted, and these points include the eight Abell clusters found in the same volume as our poor clusters. The X-ray observations are surface-brightness-limited, and they contain a small Malmquist-type bias; so, the first point in the XLF had to be computed using the smaller volume to which the first bin is complete to compensate for the bias.

We also plot the XLF for rich clusters as filled circles in Figure 2 using a subsample of Abell clusters observed with the RASS by Briel & Henry (1993). We used only those clusters with $z \le 0.15$, which is a relatively complete sample of richness class ≥ 0 clusters (e.g., Ebeling et al. 1996). Luminosities were corrected to our cosmology, and converted to our observed energy band, assuming a free-free spectrum with T=3 keV and $r_c=95$ h^{-1} kpc for richness class 0 clusters, and T=6 keV and $r_c=125$ h^{-1} kpc for all others. Our subsample contains 67 Abell clusters of which 33 (49%) were detected with luminosities ranging from 0.3 to 8.1×10^{43} h^{-2} ergs s⁻¹. The first two rich cluster bins also suffer from a Malmquist bias and were corrected in the same manner as the poor clusters.

In computing the XLF, we used both a direct binning approach (treating upper limits as detections) and a Kaplan-Meier (e.g., Feigelson & Nelson 1985) estimator for detections and upper limits. The two methods produced identical slopes for the differential and cumulative XLFs.

As shown in Figure 2, the XLF is well fit to a power law in the following form: $-5.98(\pm0.22)$ – $1.71(\pm0.19) \times \log (L_X/10^{44})$. Within a 90% confidence level, our XLF agrees with those of Edge et al. (1990) and Henry et al. (1995) in regions where they overlap. It appears that the XLF for our optically selected

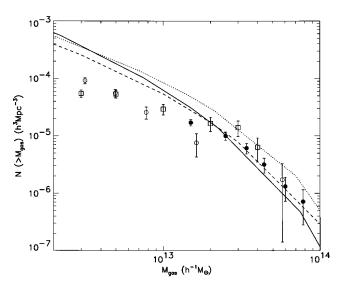


Fig. 3.—Integral cluster gas mass function for the observed poor clusters (open circles), the Abell clusters (filled circles), and the numerically "observed" clusters from the CHDM simulation (squares). We have assumed h=0.5 in order to plot the observationally determined masses. Using the same mass bins each time, 500 realizations of the observed GMF were made with randomly applied errors of up to 50% on the masses; however, we found that the Poisson error bars (shown here) dominate. Press-Schechter predictions for the GMF are shown for the CHDM model (solid line) and flat, Λ + CDM universes with $\Omega=0.5$ (dotted line) and $\Omega=0.3$ (dashed line).

poor clusters is a smooth extrapolation of the XLF for nearby Abell clusters. This further suggests that the bulk X-ray properties of poor clusters are not fundamentally different from rich clusters.

4. THE INTRACLUSTER MEDIUM MASS FUNCTION

We estimated the gas masses for the poor clusters in our sample using the measured $L_{\rm X}$ values and again assuming that X-ray surface brightness follows a β -model. The X-ray mass within a radius r is given by $M_{\rm gas}(\le r) = 4\epsilon^{-1/2}L_{\rm X}^{1/2}$ $r_c^{3/2}[r/r_c-\arctan(r/r_c)]$ for $\beta=2/3$, where ϵ is the volume emissivity (calculated for a Raymond-Smith plasma using XSPEC). In Table 1, we list the gas masses within 1.5 h^{-1} Mpc of the cluster centers assuming $r_c=65$ h^{-1} kpc and T=1 keV

In an effort to assess the reliability of our observational gas mass estimates, we calculated the gas masses of our three-dimensional clusters in the CHDM volume, using the above approach based on their luminosity, and we compared this with their known gas masses. We find that the observational mass determination does a reasonable job of estimating the true masses although there is substantial scatter ($\approx 50\%$ standard deviation).

In Figure 3, we show the integral gas mass function (GMF) for the cluster samples. We also calculated the GMF that would be observed in the CHDM simulation volume. The GMF is somewhat more reliable in grid-based numerical simulations than the XLF (Anninos & Norman 1996). Since we had previously compiled a list of three-dimensional X-ray clusters within the volume, we were able to correlate their projected positions on the sky with those of the optical groups. We found that 77%–85% of the optical groups had projected X-ray emission within 20" of the optical group centroid. The GMF for those three-dimensional clusters that coincided with a complete, volume-limited subsample of our optical groups,

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using the same selection function as the observed groups, is shown in Figure 3. The fairly good agreement between the model and the observations for the poor clusters again suggests that our observed clusters in Table 1 form a nearly complete subsample of all groups with the selection criteria given in § 2.

The curves in Figure 3 correspond to Press-Schechter (PS) predictions for the GMF in three different cosmologies. The PS GMF for the CHDM model agrees well with the GMF for the Abell clusters but lies above that of the poor clusters. This is not unexpected because of the particular selection criteria used in constructing the sample in Table 1; we have selected only a subset (albeit complete) of all possible groups within our observed volume. The remaining two curves correspond to flat, COBE-normalized, Λ + CDM models ($Q_2 = 21.8 \mu K$): one with $\Omega = 0.5$, $\Omega_b = 0.035$, $\sigma_8 = 1.22$, h = 0.6, and the other with $\Omega=0.3$, $\Omega_b=0.03$, $\sigma_8=1.09$, and h=0.7. We see that the GMF amplitude drops as the total mass density of the universe decreases, and that the GMF can potentially be used to constrain cosmological models. The Λ + CDM model with $\Omega = 0.3$ also fits the rich cluster observational data very well, although the model with $\Omega = 0.5$ clearly does not (even decreasing Q_2 by the maximum allowable 10% results in only marginal improvement). Since we do not observe all possible low-mass groups, the conclusions are less clear for poorer clusters. Numerical simulations of these models are needed to include observational selection effects. Nevertheless, the fact that both the CHDM and the Λ + CDM models give similar results for the GMF indicates that our modeling of the selection effects is not strongly dependent on the assumed CHDM cosmology.

5. CONCLUSIONS

We have constructed a sample of optically selected poor galaxy clusters and measured their X-ray properties from the ROSAT all-sky survey. This is the most extensive sample of poor cluster X-ray properties published to date. We combined this sample of poor clusters with a complete subsample of $z \le 0.15$ Abell clusters with similar RASS data to form a list of clusters with a wide range of richnesses, masses, and ICM properties. In an effort to understand the observational selection effects, we used a CHDM numerical simulation to model

the cluster Gas Mass Function. Both the observational results (volume density vs. z) and the numerical model indicate that our sample is approximately volume limited for groups with more than four galaxies brighter than M^* , within a radius of $180 h^{-1}$ kpc. The difference between the Press-Schechter and simulated "observed" groups at low M_{gas} is due to the lack of brighter galaxies in some groups. The observed GMF agrees well with a simple CHDM, flat Universe model with a constant baryon fraction of 7.5%, although we do not claim that it is a unique fit to the data.

From our optically selected samples, we have produced the first X-ray luminosity function that includes both rich and poor clusters and, thus, spans a wide range in luminosities. We find that both the XLF and the intracluster gas mass function are smoothly continuous curves from poor to rich clusters. These functions are broadly consistent with a hierarchical clustering model in which rich clusters are formed via mergers of poorer clusters, and this process is ongoing at the present epoch. Renzini et al. (1993) have argued that strong stellar winds early in the life history of ellipticals may result in the ejection of much of the original gas from galaxy groups. However, the continuity of the XLF and GMF over a large range in cluster richness is not consistent with the expulsion of a significant quantity of gas from poor groups.

Our new XLF, in particular, should be useful for studies of the evolution of optically selected clusters. Our rich and poor cluster XLF will provide a low-z baseline, against which higher z cluster XLFs can be compared. There has been tantalizing evidence presented for evolution in the cluster XLF for some samples of z > 0.3 clusters (e.g., Henry et al. 1992; Castander et al. 1994). The abundance of poor clusters relative to rich clusters will provide important constraints on detailed models for the formation of rich clusters (e.g., Castander et al. 1995).

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