

## NEUTRAL OXYGEN AND NITROGEN LINES IN PLANETARY NEBULAE

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## ABSTRACT

The strength of the  $\lambda 6300$  [O I] line in optically thin planetary nebulae is linearly dependent on that of  $\lambda 3727$  [O II] over a range of more than two orders of magnitude, and averages 6.7% of  $I(\lambda 3727)$ . In optically thick planetaries with cool central stars ( $\log T_* < 4.65$ ), which contain significant neutral helium, the strengths of [O I] and [O II] are also correlated, but for these the strength of  $I(\lambda 6300)$  is suppressed and averages only 1.9% of  $I(\lambda 3727)$ . The  $\lambda 5199$  [N I] and  $\lambda 6584$  [N II] lines show an analogous but not quite so linear behavior;  $I(\lambda 5199)$  averages 1.3% of  $I(\lambda 6584)$  for thin nebulae, and only 0.3% for above thick nebulae. Real differences from the averages are indicated for specific nebulae. The lower relative strengths of the neutral lines in optically thick nebulae with cooler central stars imply that these nebulae exhibit less clumping than do the optically thin planetaries.

*Subject headings:* nebulae: planetary

## I. INTRODUCTION

The [O I] lines of planetary nebulae have historically presented a problem that now is considered largely solved. When compared with homogeneous photo-ionization models, the forbidden lines at  $\lambda 6300$  and  $\lambda 6363$  are far too strong (Campbell 1968; Flower 1969; Williams 1970). Williams (1973), however, showed that with the inclusion of condensations and charge exchange, theory can be made consistent with the observations.

The most extensive studies of [O I] to date are by Khromov (1965), who used line intensities measured in 30 objects by Vorontsov-Vel'yaminov *et al.* (1964), and by Campbell (1968), who determined the  $I(\lambda 6300)/I(H\alpha)$  ratio in 16 nebulae. Both authors show that the strength of [O I] correlates with that of [N II], and Khromov suggests a correlation with  $\lambda 5876$  He I, implying that the abundances of low-ionization species are also correlated. Khromov also indicated a possible positive relation between the strength of [O I] and electron density.

There is now so much more data available, however, particularly from the investigations of Aller and Czyzak (1979), Barker (1978), Kondratyeva (1978), and Torres-Peimbert and Peimbert (1977), that it is possible to examine [O I] line-strength variations among nebulae in much greater detail. In addition, the work of Hawley and Miller (1977, 1978a) and Hawley allows point-to-point examination of [O I] in single nebulae.

The [N I] line at  $\lambda 5199$  is more difficult to study because it is weaker than [O I]. But extensive observations, principally by Aller and Czyzak (1979) and Aller and Walker (1970), now allow the detailed study of the behavior of this line, and its relation to that of the [O I] line strengths.

## II. THE BEHAVIOR OF [O I] AND [N I]

The relation between [O I] and [N II] discovered by Khromov (1965) and Campbell (1968) suggests that [O I] and [N I] should be examined with respect to [O II] and [N II], respectively. The oxygen-line data are presented in Table 1. Column (1) gives the nebula's name, and columns (2) and (3) the logarithms of  $I(\lambda 6300)$  and  $I(\lambda 3727)$ , respectively, on the scale  $I(H\beta) = 100$ , corrected for interstellar extinction. In cases where  $\lambda 6300$  was blended with  $\lambda 6312$  [S III],  $I(\lambda 6300)$  was set equal to 3.14  $I(\lambda 6363)$  (Garstang 1968). Column (4) gives  $r(O) = I(\lambda 6300)/I(\lambda 3727)$ , and column (5) the reference to the  $\lambda 6300$  observation. Precedence was given to studies in which both  $\lambda 3727$  and  $\lambda 6300$  were observed. When  $\lambda 3727$  was not observed by the given reference, a mean  $I(\lambda 3727)$  was taken from Kaler (1976) plus the references of column (5). References to notes are shown in column (6). Averages of  $r(O)$  for nebulae with multiple observations are also placed in parentheses in column (6). Points are considered "half-weight" if the [O I] and [O II] data were taken from different references, if there was serious disagreement between two observations, or if  $I(\lambda 6363)$  and  $I(\lambda 6300)$  were discrepant. If  $I(\lambda 6300)$  as computed from  $I(\lambda 6363)$  was significantly larger than the [O I] and [S III] blend,  $I(\lambda 6300)$  was set equal to the blend intensity, and given half-weight.

The nebulae of Table 1 are divided into three groups, depending upon the excitation of the nebular spectrum and the central star temperature ( $T_*$ ), which indicates whether a nebula can be called optically thick or thin in the Lyman continuum. The nebulae in Table 1a all have more than a trace of He II  $\lambda 4686$ , and from Hummer and Seaton (1964) are likely to be optically thin. Those in Table 1c all have  $\log T_* < 4.65$ . These nebulae contain neutral helium (Kaler 1978), and must

TABLE 1  
[O I] AND [O II] LINE DATA

Nebula (1)	log $I(6300)$ (2)	log $I(3727)$ (3)	$r(O)$ (4)	Ref. (5)	Notes (6)	Nebula (1)	log $I(6300)$ (2)	log $I(3727)$ (3)	$r(O)$ (4)	Ref. (5)	Notes (6)
a) $I(\lambda 4686) \geq 5\% I(H\beta)$ (thin)											
NGC 650	+1.60	2.90	0.050	A		NGC 6905	+0.20	1.28	0.083	A	$\frac{1}{2}$
NGC 2346	+1.19	2.56	0.043	A		NGC 7009	-0.32	1.24	0.026	CA, P	
NGC 2371	+0.62	1.57	0.113	T		NGC 7009A <sup>a</sup>	+1.41	2.11	0.20	CA	$\frac{1}{2}$
NGC 2392	+0.35	1.99	0.025	A, Z		NGC 7026	+0.63	1.74	0.078	A	d
NGC 2438	+1.25	2.46	0.061	T		NGC 7027	+1.15	1.52	0.425	KA, O, P	D
NGC 2440	+1.46	2.28	0.151	T		NGC 7293-A1	+1.60	2.37	0.171	H	
NGC 2452	+0.77	1.90	0.081	A		NGC 7293-A2	+1.43	2.42	0.102		
NGC 3132a	+1.18	2.51	0.046	T		NGC 7293-A3	+1.45	2.48	0.092		
NGC 3132b	+1.47	2.58	0.078			NGC 7293-A4	+1.68	2.66	0.106		(0.117)
NGC 3132c	+1.32	2.65	0.047		(0.057)	NGC 7293-B1	+1.69	2.71	0.097		
NGC 3132d	+1.65	2.96	0.049			NGC 7293-B2	+1.70	2.68	0.103		
NGC 3132e	+1.77	2.95	0.067			NGC 7293-B3	+1.80	2.65	0.142		
NGC 3242	+0.00	1.16	0.069	C, BR	$\frac{1}{2}$	NGC 7293-B4	+1.83	2.74	0.123		
NGC 3587	+1.06	2.54	0.033	BR	$\frac{1}{2}$	NGC 7662	-0.24	1.14	0.041	C	$\frac{1}{2}$
NGC 3918	+0.59	1.79	0.063	T	D	IC 351	-0.14	1.04	0.066	BA	
NGC 6644	+0.96	1.76	0.156	BA	d	IC 1747	+0.28	1.62	0.046	A	
NGC 6720-1	+0.36	1.90	0.029	H1		IC 2003	+0.54	1.70	0.069	A	
NGC 6720-2	+0.63	2.24	0.024			IC 2165	+0.37	1.68	0.046	A, T	
NGC 6720-3	+0.85	2.22	0.043			IC 4406	+1.50	2.65	0.072	T	
NGC 6720-4	+1.41	2.83	0.038		(0.042)	IC 5117	+0.70	1.46	0.174	A	D
NGC 6720-5	+2.09	3.25	0.069			IC 5217	+0.28	1.41	0.073	A	d
NGC 6720-6	+1.99	3.30	0.049			BB 1	+0.02	0.86	0.144	H3	
NGC 6741	+0.30	2.09	0.016	C	$\frac{1}{2}$	Ha 4-1	+0.86	2.27	0.039	T1, H3	
NGC 6778	+0.70	1.78	0.083	A		Hu 1-1	+1.45	2.61	0.070	A, BA, K	
NGC 6790	+0.52	1.36	0.145	A	D	Hu 1-2	+0.80	1.76	0.125	A	d
NGC 6818	+0.27	1.72	0.036	A		J900	+0.55	1.84	0.051	A	
NGC 6853-2	+1.10	2.54	0.036	H2		M1-1	-0.02	1.20	0.070	A, BA	$\frac{1}{2}$
NGC 6853-5	+1.89	3.22	0.047		(0.047)	M1-17	+1.34	1.58	0.57	K	D
NGC 6853-6	+1.96	3.21	0.057			M1-33	+1.01	1.78	0.17	K	D
NGC 6884	+0.03	1.59	0.028	A	d	Vy 1-2	+0.51	1.60	0.081	BA	$\frac{1}{2}$ , d
NGC 6886	+1.00	2.20	0.063	A	d						
b) $\log T_* \geq 4.65$ , $I(4686) \leq 5\% I(H\beta)$											
NGC 5315	+0.76	1.66	0.127	T	D	IC 2501	+0.65	1.81	0.070	T	D
NGC 6210	+0.35	1.63	0.054	BA		IC 4846	+0.51	1.52	0.094	A, BA	D
NGC 6543	+0.11	1.49	0.035	A, BC, C	d	J320	-0.10	1.29	0.041	BA	
NGC 6567	+0.35	1.36	0.100	BA		M1-9	+0.49	1.54	0.089	K	D
NGC 6572	+0.79	1.67	0.131	P	d	M1-74	+0.55	1.53	0.105	A, BA	D
NGC 6751	+0.54	2.55	0.0098	A	d	M2-50	+0.76	1.31	0.279	BA	D, $\frac{1}{2}$
NGC 6826	-0.85	1.41	0.0054	BR	$\frac{1}{2}$	Me 2-2	+0.29	1.48	0.065	BA	d
NGC 6833	+0.17	1.20	0.093	BA	d	Sn 1	+0.69	1.28	0.253	BA	$\frac{1}{2}$
NGC 6891	-0.07	1.37	0.037	C	$\frac{1}{2}$						
c) $\log T_* < 4.65$ (thick)											
NGC 40	+0.47	2.65	0.0066	A		He2-131	+0.31	1.88	0.027	T	D
IC 418	+0.36	2.23	0.013	A, O, T	D	He2-138	+0.05	1.56	0.031	T	
IC 2149	-0.09	1.93	0.0096	BA		Hu 2-1	+0.18	1.95	0.018	A, BA	d
IC 4593	+0.23	1.80	0.025	BA, C		M1-5	+0.47	2.06	0.026	BA	d
BD +30	+0.41	1.99	0.026	O, P	d	M1-14	+0.55	2.27	0.019	T	
Cn 3-1	+0.25	2.28	0.0097	A, BA		M1-65	+0.11	1.90	0.016	K1	D
Hb -12	-0.86	1.17	0.0095	A	$\frac{1}{2}$	M2-9	+1.08	2.33	0.056	BA	D, $\frac{1}{2}$

<sup>a</sup> NGC 7009A, east ansa.

NOTES.— $\frac{1}{2}$  = half weight; D =  $\log N_e > 4$ ; d =  $3.8 < \log N_e < 4$ .

REFERENCES.—A: Aller and Czyzak 1979; BA: Barker 1978; BC: Boeshaar, Czyzak and Aller 1975; BR: Boeshaar 1974; C: Campbell 1968; CA: Czyzak and Aller 1979; H: Hawley 1978; H1: Hawley and Miller 1977; H2: Hawley and Miller 1978a; H3: Hawley and Miller 1978b; K: Kondratyeva 1978; K1: Kondratyeva 1979; KA: Kaler *et al.* 1976; O: O'Dell 1963; P: Peimbert and Torres-Peimbert 1971; T: Torres-Peimbert and Peimbert 1977; T1: Torres-Peimbert and Peimbert 1979; Z: Zipoy 1976.

TABLE 2  
 [N I] AND [N II] LINE DATA

Nebula (1)	log $I(\lambda 5199)$ (2)	log $I(\lambda 6584)$ (3)	$10^2$ $r(N)$ (4)	Ref. (5)	Nebula (1)	log $I(\lambda 5199)$ (2)	log $I(\lambda 6584)$ (3)	$10^2$ $r(N)$ (4)	Ref. (5)
a) $I(\lambda 4686) \geq 5\%$ , $I(H\beta)$ (thin)									
NGC 650	+1.1	2.88	1.66	A	NGC 6818	-0.19	1.40	2.57	A
NGC 2371	-0.16	1.98	0.72	A	NGC 6884	-0.60	1.33	1.17	A
NGC 2392	-0.22	1.93	0.71	T	NGC 6886	+0.06	2.24	0.66	A
NGC 2440	+1.11	3.07	1.10	AW	NGC 7009	-0.44	1.10	0.36	CA, P
NGC 2452	+0.19	2.19	1.00	A	NGC 7009A <sup>a</sup>	+0.56	2.56	1.00	CA
NGC 3242	-0.55	0.49	9.1	T	NGC 7026	-0.04	2.09	0.74	A
NGC 3587	+0.08	2.10	0.96	BR	NGC 7027	-0.18	1.96	0.74	KA, P
NGC 3918	-0.26	1.83	0.81	T	NGC 7293S <sup>b</sup>	+1.43	2.94	3.18	W
NGC 5882	-0.64	1.14	1.66	T	NGC 7293C	+1.06	2.11	8.91	W
NGC 6302	+1.10	2.70	2.51	AW	NGC 7662	-0.97	0.71	2.09	A
NGC 6309	-0.08	2.15	0.59	AW	IC 351	-0.85	0.29	7.24	A
NGC 6720-1	-0.23	1.70	1.17	H	IC 2003	+0.09	1.49	3.98	A
NGC 6720-2	+0.20	2.02	1.51		IC 2165	-0.83	1.30	0.74	A
NGC 6720-3	+0.07	2.35	0.53		IC 5117	-0.53	1.58	0.78	A
NGC 6720-4	+0.75	2.71	1.10		IC 5217	-0.90	1.39	0.51	A
NGC 6720-5	+1.49	2.97	3.31		Ha 4-1	-0.01	1.99	1.0	T1
NGC 6720-6	+1.29	3.04	1.78		Hu 1-1	+0.37	2.35	1.05	A
NGC 6741	+0.32	2.31	1.02	AW	Hu 1-2	+0.16	2.27	0.78	A
NGC 6778	+0.38	2.53	0.71	A	J900	-0.55	1.55	0.79	A
NGC 6790	-0.38	1.16	2.88	A					
b) $\log T_* \geq 4.65$ , $I(\lambda 4686) \leq 5\%$ , $I(H\beta)$									
NGC 5315	-0.09	2.30	0.41	T	NGC 6826	-1.40	1.00	0.40	AW
NGC 6543	-0.67	1.21	1.32	A	NGC 6891	-0.72	1.02	1.83	AW
NGC 6572	-0.16	1.74	1.26	A, P	IC 4997	-0.72	1.29	0.98	AW
c) $\log T_* < 4.65$ (thick)									
NGC 40	-0.08	2.42	0.31	A	Cn 3-1	-0.65	2.29	0.12	A
IC 418	-0.53	2.21	0.18	A, T	Hb 12	-0.55	1.96	0.31	A
IC 2149	-0.36	1.62	1.05	T	Hu 2-1	-1.77	1.68	0.04	A
BD +30	-0.18	2.46	0.23	P					

<sup>a</sup> NGC 7009A, east ansa.

<sup>b</sup> NGC 7293S, southwest part of ring; NGC 7293C, central region.

REFERENCES.—A: Aller and Czyzak 1979; AW: Aller and Walker 1970; BR: Boeshaar 1974; CA: Czyzak and Aller 1979; H: Hawley and Miller 1977; KA: Kaler *et al.* 1976; P: Peimbert and Torres-Peimbert 1971; T: Torres-Peimbert and Peimbert 1977; T1: Torres-Peimbert and Peimbert 1979; W: Warner and Rubin 1975.

all be optically thick. The objects in Table 1b are intermediate between the above two groups; they are usually considered as thick since there is little or no He II in their spectra, but they have  $\log T_* > 4.65$ .

Table 2 is entirely analogous to Table 1. Columns (2), (3), and (4) give  $\log I(\lambda 5199)$  [N I],  $\log I(\lambda 6584)$  [N II], and  $r(N) = I(\lambda 5199)/I(\lambda 6584)$ , respectively. Since the data on the [N I] lines are intrinsically poorer than those on [O I], "half-weight" nebulae are not singled out.

The ratios  $r(O)$  and  $r(N)$  will show some dependence on electron density. From the target areas of Saraph and Seaton (1974), Pradhan (1976) and Henry and Williams (1968), and the transition probabilities of Garstang (1968) and Nussbaumer (1971), the critical densities at which radiative de-excitation of the nebular lines equals collisional de-excitation (for  $T_e = 10^4$ ) for [O I], [N I], [O II], and [N II] is  $10^6$ , 2

$\times 10^3$ ,  $8 \times 10^3$ , and  $7 \times 10^4 \text{ cm}^{-3}$ . Electron densities were calculated from [O II], [Cl III] or [S II] data provided by the references in Table 1, and Saraph and Seaton (1970), and nebulae with  $\log N_e \geq 4.0$  or with  $3.8 \leq \log N_e < 4.0$  are indicated as such in column (6) of Table 1. Collisional de-excitation is not significant for [O I], but it is important in suppressing  $\lambda 3727$  of [O II] for higher densities. Consistent with this expectation, Table 1 shows that the nebulae with the highest density have the highest values of  $r(O)$ . A plot of  $r(O)$  against  $N_e$  shows that  $r(O)$  tends to be larger for densities above about  $6 \times 10^3$ .

The effect of density on  $r(N)$  should be the opposite of that for  $r(O)$ ; [N I] will show a larger relative collisional de-excitation rate than will [N II]. The observed effect is not so obvious as it is for oxygen. The suppression of [N I] can be seen in that for optically thin nebulae with  $\log N_e \leq 3.8$ , 39% of the nebulae

have  $r(\text{N}) > 0.015$ , whereas for  $\log N_e > 3.8$ , the figure is only 10%. Yet a large number of low-density nebulae also have low  $r(\text{N})$ . In general, the density effects cannot be considered quantitatively for either [O I] or [N I] since the densities in the regions where the neutral lines are formed may be different from those which produced the ionized lines.

The values of  $\log I(\lambda 6300)$  in Table 1 are plotted against those of  $\log I(\lambda 3727)$  in Figure 1. The nebulae of Tables 1a, 1b, and 1c are plotted with different symbols, as shown in the caption. The objects in Tables 1a and 1b with  $\log N_e > 3.8$  are not plotted; those with  $\log N_e > 3.8$  from Table 1c are indicated by a horizontal bar. The "half-weight" points are denoted by smaller symbols.

This figure displays a number of interesting features. First is the extreme linearity of curve defined by the optically thin nebulae. A quadratic least-squares fit is almost exactly a straight line, with a slope of 1. For these nebulae,  $I(\lambda 6300)$  is directly proportional to  $I(\lambda 3727)$ . If we consider the optically thin points of Figure 1, excluding those of half-weight and first averaging the nebulae with multiple observations,  $I(\lambda 6300)/I(\lambda 3727)$  averages 0.067. The standard deviation is large, here  $\pm 0.035$ . Much of this scatter is due to observational error, but some is clearly intrinsic. Note, for example, that all the individual points of NGC 7293 (see the legend to Fig. 1) are above the mean line, and that all but one of the points for NGC 6720 are below it.

Second, the nebulae of Table 1c, those thick objects with the coolest central stars, and which contain

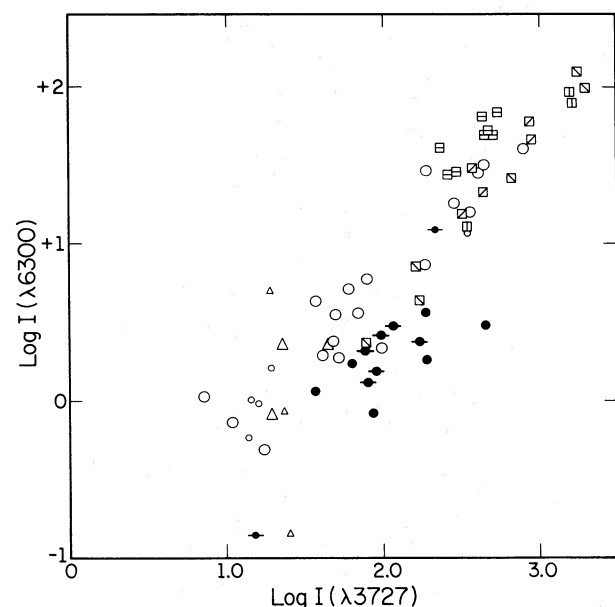


FIG. 1.— $\log I(\lambda 6300)$  [O I] plotted against  $\log I(\lambda 3727)$  [O I].  $\circ$ : optically thin;  $\square$ : optically thin nebulae for which point-to-point observations are available;  $\boxtimes$ : NGC 3132;  $\boxminus$ : NGC 6720;  $\boxplus$ : NGC 6853;  $\boxdot$ : NGC 7293;  $\triangle$ : optically thick nebulae with  $\log T_* \geq 4.65$  (weak or absent He II);  $\bullet$ : optically thick nebulae with  $\log T_* < 4.65$ ;  $\ominus$ : same as above but with  $\log N_e > 3.8$ ; half-weight points are denoted by small symbols.

neutral helium, are clearly below the line indicated by the thin nebulae. Note that for this set,  $r(\text{O})$  is *not* significantly higher for the dense nebulae:

$$\frac{\langle r(\text{O}) \rangle (\log N_e \geq 3.8)}{\langle r(\text{O}) \rangle (\log N_e < 3.8)} = 1.26 \pm 0.33.$$

For these thick nebulae,  $I(\lambda 6300)$  averages only 1.9% of  $I(\lambda 3727)$ , where the dense nebulae are included. The difference between the thin nebulae of Table 1a and the thick objects of Table 1c is highly significant:

$$\frac{\langle r(\text{O}) \rangle (1a)}{\langle r(\text{O}) \rangle (1c)} = 3.5 \pm 0.6.$$

For these thick nebulae,  $I(\lambda 6300)$  is suppressed relative to  $I(\lambda 3727)$  by over a factor of 3.

Third, note that  $\lambda 6300$  bears about the same relation to  $\lambda 3727$  within a given nebula as it does from one nebula to another. For example, the curve defined by NGC 6720 is similar to that defined by the general run of thin objects.

Finally, the thick nebulae with hotter central stars, those of Table 1b, fit well with the thin nebulae of Table 1a, and not with the thick objects of 1c. The strength of  $\lambda 6300$  is suppressed only when a nebula is very thick, as defined by the facts that the central star temperature is low and that helium is neutral in the outer zone.

The relation between [O I] and electron density found by Khromov (1965) is not confirmed. In fact, the effect is just the opposite; the highest [O I] strengths are found in the large, low-density nebulae such as NGC 7293.

The [N I] and [N II] data are similarly plotted in Figure 2. However since the effects of density on  $r(\text{N})$

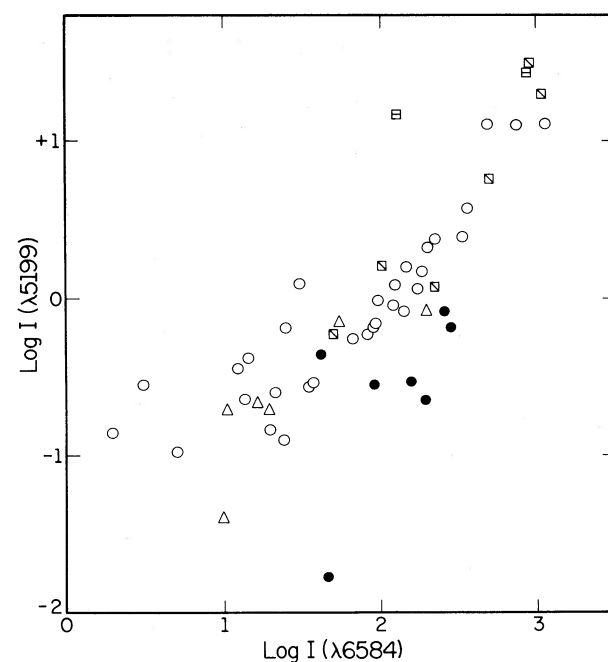


FIG. 2.— $\log I([\text{N I}]\lambda 5199)$  plotted against  $\log I([\text{N II}]\lambda 6584)$ . See the caption to Fig. 1 for an explanation of the symbols.



are not very clear, all the nebulae are considered, regardless of  $N_e$ . The conclusions for nitrogen are very similar to those for oxygen. The curve for nitrogen is not quite linear;  $\langle r(\text{N}) \rangle$  for  $I(\lambda 6584) > 2.4$ ,  $1.8 < I(\lambda 6584) < 2.4$ , and  $1.0 < I(\lambda 6584) < 1.8$  is 0.016, 0.008, and 0.016, respectively. However, the effects of observational selection, especially at the faint end, cannot be evaluated. A least-squares linear fit to all the points for  $\log I(\lambda 6584) > 1$  (where NGC 6720 is averaged) gives an overall slope of 1.03. The average value of  $r(\text{N})$  for these thin nebulae is 0.013. The  $\lambda 5199$  [N I] line is suppressed by a factor of  $4.1 \pm 1.8$  for the optically thick nebulae of Table 2c, which agrees within the errors with the suppression factor for  $\lambda 6300$  [O I]. If we exclude the nebulae of Tables 1c and 2c, the values of  $r(\text{O})$  and  $r(\text{N})$  do not correlate.

The empirical results for  $r(\text{O})$  and  $r(\text{N})$  are summarized in Table 3. The half-weight points are not used for oxygen, the data for the four objects with multiple point-to-point observations are averaged and treated as a single data point, and the three nebulae with the weakest [N II] are not included in  $r(\text{N})$ . High-density nebulae are excluded for oxygen from the set of thin nebulae and the thick nebulae with hotter central stars, but not from the thick nebulae with cool central stars, and not from the nitrogen data. Column (2) gives the mean  $r$  for both  $I(\lambda 6300)/I(\lambda 3727)$  and  $I(\lambda 5199)/I(\lambda 6584)$ . Columns (3) and (4) give the standard deviations and the mean errors of the means, respectively, and columns (5) and (6) show the slopes and correlation coefficients for linear least-squares fits to Figures 1 and 2, under the assumption that all the error is in the  $y$ -axis.

### III. DISCUSSION

The linear relationship between [O I] and [O II] and between [N I] and [N II] are simple to understand, since  $\text{O}^0$  and  $\text{N}^0$  must come from  $\text{O}^+$  and  $\text{N}^+$ , whether by recombination or charge exchange. The surprise in this study is that [O I] and [N I]

are suppressed, for a given  $I(\lambda 3727)$  or  $I(\lambda 6584)$ , for the thick nebulae which have central stars cool enough such that a significant amount of helium is neutral.

The neutral oxygen and nitrogen lines must be produced in transition zones between ionized and neutral matter. As pointed out by Williams (1970, 1973) condensations are necessary in order to produce the large transition zone area necessary for the production of the high observed line strengths. The data imply that the set of nebulae with the coolest central stars ( $\log T_* < 4.65$ ) are more homogeneous than are the nebulae with higher temperature stars. The objects with low  $T_*$  are small and presumably relatively young. Possibly widespread filamentation has not yet had time to develop. The dependence of line strength on filamentary structure is supported by the data on NGC 7293, which exhibits some of the highest values of  $r(\text{O})$  and  $r(\text{N})$ , and which exhibits some of the most extreme filamentary structure known (see the widely reproduced photograph taken by Baade in, e.g., Perek and Kohoutek 1967). The fact that  $I(\lambda 6300)/I(\lambda 3727)$  is not increased for the dense, thick nebulae of Table 1c implies that for these objects the increased density correlates with a further decrease in the existence of filamentary structure.

### IV. SUMMARY

The results presented here confirm the relation found by Khromov (1965) and Campbell (1968) between neutral and singly ionized species. The detailed analysis of the [O I] and [N I] lines in the spectra of planetary nebulae reveals the following:

1. The strength of  $\lambda 6300$  [O I] in optically thin planetaries is directly proportional to that of  $\lambda 3727$  [O II] over a range of more than two orders of magnitude.
2.  $I(\lambda 6300)$  averages 6.7% of  $I(\lambda 3727)$  for optically thin, lower density planetaries.
3. High density ( $\log N_e > 3.8$ ) increases the value of  $r(\text{O}) = I(\lambda 6300)/I(\lambda 3727)$  in optically thin nebulae because of collisional de-excitation of  $\text{O}^+$  ( $^2\text{D}$ ).

TABLE 3  
LINE STRENGTH RATIOS,  $r(\text{O})$  AND  $r(\text{N})$

Type (1)	$\bar{r}(\text{O}) = \langle I(6300)/I(3727) \rangle$ (2)	Standard Deviation (3)	m. e. of Mean (4)	Slope (5)	Correlation Coefficient (6)
Thin.....	0.067	0.035	0.007	0.97	0.95
Thick, $\log T_* \geq 4.65$ .....	0.065	0.017	0.010	...	...
Thick, $\log T_* < 4.65$ .....	0.019	0.008	0.002	0.43 0.93 <sup>a</sup>	0.63 0.80 <sup>a</sup>

Type	$\bar{r}(\text{N}) = \langle I(5199)/I(6584) \rangle$	Standard Deviation	m. e. of Mean	Slope	Correlation Coefficient
Thin.....	0.0126	0.0089	0.0016	1.03	0.90
Thick, $\log T_* \geq 4.65$ .....	0.0103	0.0056	0.0023	...	...
Thick, $\log T_* < 4.65$ .....	0.0032	0.0034	0.0013	0.97	0.59

NOTE.— $\bar{r}(\text{O})(\text{thin})/\bar{r}(\text{O})(\text{thick}, \log T_* < 4.65) = 3.5 \pm 0.6$ ;  $\bar{r}(\text{N})(\text{thin})/\bar{r}(\text{N})(\text{thick}, \log T_* < 4.65) = 4.1 \pm 1.8$ .

<sup>a</sup> With two half-weight points.

4. Planetary with central stars which have temperatures less than 44,000 K ( $\log T_* < 4.65$ ), which are optically thick and contain neutral helium, have relatively lower  $\lambda 6300$  strengths;  $I(\lambda 6300)$  averages only 1.9% of  $I(\lambda 3727)$ . The [O I] and [O II] line strengths are again proportional, but the increase of  $r(O)$  for high-density nebulae is not seen.

5. The thick planetaries with higher central star temperatures show [O I] line strengths in agreement with those of the thin nebulae.

6. The [N I] and [N II] lines bear the same relations to one another as do the [O I] and [O II] lines, except that the curve defined by  $I(\lambda 5199)$  versus  $I(\lambda 6584)$  may not be quite linear, and density effects [here the reduction of  $r(N) = I(\lambda 5199)/I(\lambda 6584)$  for high density] are not clearly seen.

7.  $I(\lambda 5199)$  averages 1.35% of  $I(\lambda 6584)$  for thin nebulae and 0.32% for thick nebulae with cool central stars.

8. Real differences in [O I]/[O II] (and probably

[N I]/[N II]) are seen among nebulae; for example, NGC 7293 and NGC 6720 show [O I] strengths consistently larger and smaller than the average, respectively.

9. The relation between [O I] and [O II], and between [N I] and [N II], is about the same within a particular nebula as it is from one nebula to another.

Since condensations are needed in order to produce the high neutral line strengths found in planetaries, the above results indicate that optically thin nebulae have a greater degree of filamentation than do the thick nebulae with cool central stars, in which the neutral lines are suppressed. There is also some evidence that high density in the thick, low  $T_*$  objects further suppresses the incidence of condensations and filamentary structure.

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*Note added in proof.*—L. H. Aller informs me that  $\log I(5199)$  [N I] for Hu 2-1 in Aller and Czyzak (1979) should read  $-0.77$ , not  $-1.77$ , which raises the low point in Figure 2 into good accord with the other optically thick points. With this change, columns (2)–(6) in the last row of Table 3 should read 0.0037, 0.0031, 0.0012, 0.39, and 0.53, respectively, and  $\bar{r}(N)(\text{thin})/\bar{r}(N)(\text{thick}, \log T_* < 4.65) = 3.5 \pm 1.2$ , the same suppression factor as for oxygen.

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