# THE REDSHIFT-DISTANCE RELATION. V. GALAXY COLORS AS FUNCTIONS OF GALACTIC LATITUDE AND REDSHIFT: OBSERVED COLORS COMPARED WITH PREDICTED DISTRIBUTIONS FOR VARIOUS WORLD MODELS 

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#### Abstract

New $B V R$ photometry is listed for the brightest several members of clusters of galaxies, for all HMS groups, and for E galaxies that are identified with radio sources. Measurements are given for a set of 72 stars that define the S 20 (FW 130) $V-R$ system used here.

The galaxy colors, corrected for $K$ reddening, correlate with galactic latitude and provide a new solution for galactic reddening. A model for galactic absorption similar to that proposed by McClure and Crawford is supported. The reddening variation with galactic latitude is smaller by a factor of 2 than traditional values.

Redshifts of E galaxies in the range $0<z<0.9$ can be uniquely determined by a combination of four broad-band colors alone near effective wavelengths of $B, V, R$, and $\langle\lambda\rangle=8500 \AA$. The color-redshift relations are tight, and agree with predicted distributions obtained by redshifting a standard spectrum through the measuring bands. The data admit color evolution for E and S 0 galaxies at rates that must be smaller than $\Delta(B-V) / \Delta \log t \simeq 0.3$ mag color change between ages of $10^{9}$ and $10^{10}$ years, i.e., $0.3 \mathrm{mag} \mathrm{dex}{ }^{-1}$. This observational limit agrees with the color-decay rates that are expected to lie between 0.3 and $0.1 \mathrm{mag} \mathrm{dex}^{-1}$ from empirical considerations of C-M diagrams. Although small, such rates are large enough to encourage an observational search for color evolution in E galaxies where the look-back time is significantly long $(z>0.4)$.

Color distributions are calculated for various models for the production of galaxies as a function of time. Comparison with the observed distribution shows that (1) there are no young ( $t \leq 10^{9}$ years) E galaxies in our sample, and (2) the formation epoch for $E$ galaxies was peaked toward the old age side of the distribution. The data are consistent with the view that E galaxies formed at some definite time in the past. Continuous $E$ galaxy formation, as required in the steady-state model, is not verified here. The data do not contradict the hypothesis that giant E galaxies are the same age as the elliptical galaxy component of the Local Group ( $t \simeq 10^{10}$ years).


Subject headings: cosmology - galaxies, photometry of - interstellar reddening - redshifts

## I. INTRODUCTION

The energy distribution (per unit wavelength) for elliptical and S 0 galaxies is a steeply rising function from $\lambda_{0} 3300 \AA$ to $\lambda_{0} 4500 \AA$, after which it flattens redward (Oke and Sandage 1968; Whitford 1971; Schild and Oke 1971). Because of this, photometry of galaxies with redshifts in the interval $0 \leq z \leq 0.4$ will show the smallest effects of the spectrum shift ( $K$ dimming) if the effective wavelength of the detector is put longward of $\lambda_{0} \simeq 6500 \AA$.

Beginning in 1960, the commercial availability of multialkali S20 or S25 photomultipliers with their extended red response made possible routine observations of $V-R$ colors (Sandage and Smith 1963; Mendoza 1967) where the effective wavelength of $R$ is $\langle\lambda\rangle \simeq 6700 \AA$. Such photomultipliers (ITT-FW 130) were put into routine use at both Mount Wilson and Palomar in 1967. Observations of galaxies in clusters, in groups, and identified as radio sources were begun with the new equipment
in that year to supplement similar data obtained earlier with 1P21-S4 photomultipliers (Sandage 1972c, $d$, hereafter called Papers II and III).

The goals were fourfold. (1) The red observations, with their long baseline in $B-R$, give good leverage for the color-redshift relation (Whitford 1954). The extended baseline should permit a close upper limit to be put on the color evolution in the look-back time. (2) The Hubble diagram in $R$ magnitudes should be well defined because the $K_{R}$ correction for redshift is scarcely more than the precisely known bandwidth term $2.5 \log (1+z)$ for the interval $0<z<0.2$; the spectrum is nearly flat redward of $\lambda_{0} 5600 \AA$. (3) Independent photometry from that reported in Papers II and III permits an estimate of the external errors of the earlier material. Comparison gives some basis to judge the confidence on solutions for $q_{0}$ from the reliability of the photometric data alone. (4) Extend the discussion of Paper II on the richness correction by observing brightest E and S 0 galaxies in the small groups listed by Humason, Mayall, and Sandage ( 1956 [hereafter called HMS], table 11).

This fourfold program is the subject of this and the next two papers. In this report we set out the S20 photometric system in § II, list the observed data on galaxies in § III, and discuss information obtained from the color data alone in §§ IV-VII. Reddening due to the galactic system is rederived in §IV; the color-redshift relations and the conditions to be met by a four-color system that will uniquely determine redshifts in the range $0<z \leq 0.9$ are set out in $\S \mathrm{V}$; observational limits on the permitted color evolution in the look-back time are given in § VI; and predictions of color distributions for several world models are compared with the observations in § VII.

## II. THE PHOTOMETRIC SYSTEM

The photometry was done with the same filters used previously (Sandage and Smith 1963), including the $\mathrm{CuSO}_{4}$ liquid cell for red-leak suppression in $U$ and $B$. The $R$ filter was unblocked on the red side. Although tight transformations of $V-r$ to the standard $V-R$ system were obtained, it will be prudent in the future to define the longward response of the $R$ system with a filter, rather than by the cathode itself.

All data in this paper, and those parts of the data that pertain to the S20 photometry in Papers VI, VII, and VIII, are on a system defined by stars listed in table 1. These were used as night-to-night standards for the $100-\mathrm{inch}(254-\mathrm{cm})$ and $200-\mathrm{inch}(508-\mathrm{cm})$ reflectors. The system was carried progressively around the sky, and was closed upon itself between 1967 June and 1968 April.

The listed $(V-R)_{20}$ colors are not on the $(V-r)_{\mathrm{ss}}$ system of Sandage and Smith, but are, in the mean, on the system of Mendoza (1967) as determined from our overlap in M67. In turn, Mendoza's system $(V-R)_{\mathrm{M}}$ is reported to be that of the alphabet photometry of Iriarte et al. (1965) and of Johnson et al. (1966). With as many data for bright stars as are now measured therein, and with a difference only of definition via a color equation between $(V-r)_{\mathrm{SS}}$ and $(V-R)_{\mathrm{M}}$, it clearly serves no purpose to maintain separate systems that have the same information content. We therefore choose to transform our data to the Mendoza system. The approximate transformation is $(V-R)_{20}=0.94(V-r)_{\mathrm{ss}}+0.04$, where $(V-R)_{20}$ is the color listed in table 1 , and $(V-r)_{\text {ss }}$ is the system of Sandage and Smith.
$U B V$ values, obtained from 1P21 photometry, are known for many stars in table 1. A number have a long photometric history because they have been used as secondary standards (unpublished) with the 20-, $60-, 100$-, and 200 -inch Hale telescopes between 1954 and 1967, as tied to the fundamental standards of Johnson and Morgan (1953, table 3), and Johnson and Harris (1954, table 1). The $V_{20}$, and $(B-V)_{20}$ values in table 1 are close approximations to these $\mathrm{S} 4 U B V$ data, and there are no systematic differences from table 1 taken as a whole. However, any individual star listed there

TABLE 1
STARS THAT DEFINE THE S20 PHOTOMETRIC SYSTEM FOR DATA IN TABLES 2, 3, AND 4

| Object | $\alpha$ (1950) | $\delta$ (1950) | $V_{20}$ | $(U-B)_{20}$ | $(B-V)_{20}$ | $(V-R){ }_{20}$ | ${ }^{\mathrm{n}} 100$ | $\mathrm{n}_{200}$ | Chart | Remarks |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| Field Stars |  |  |  |  |  |  |  |  |  |  |
| $+30^{\circ} 50 \ldots \ldots$. | 02042 | +3103.7 | 10.18 | 1.29 | 1.23 | 0.95 | 3 | 0 |  | Peterson std: |
| F22.......... | 22742 | + 503 | 12. 74 | -0. 84 | -0. 06 | -0.07 | 1 | 2 | 1 |  |
| F24.......... | 23230 | + 331 | 12.37 | -1. 20 | -0. 22 | 0.21 |  | 0 | 1 |  |
| G5-28........ | 31541 | +1502.3 | 15.46 | 0.42 | 1.33 | 1.41 | 0 | 1 | 3 |  |
| + $30^{\circ} 663 . \ldots \ldots$. | 42240 | +3015.9 | 10.29 | 0.27 | 0.53 | 0.47 | 1 | 0 | ... | Petersonstd. |
| +30 ${ }^{\circ} 1705 \ldots$. | 82124 | + 3016.7 | 9.98 | 0.24 | 0.71 | 0.64 | 1 | 0 |  | Petersonstd. |
| F34.......... | 103642 | +4321 | 11. 23 | -1.36 | -0. 30 | -0.08 | 1 | 0 | 1 |  |
| F81........... | 132924 | +1556 | 13.48 | -1.00 | -0. 20 | -0.06 | 0 | 2 | 1 |  |
| F84.......... | 133248 | +1343 | 11.84 | -0. 74 | -0. 17 | 0.00 | 0 | 2 | 1 |  |
| F96.......... | 142800 | +2130 | 13.10 | 0.08 | 0.07 | 0.16 | 0 | 2 | 1 |  |
| F97.......... | 143348 | +3020 | 12. 35 | -0. 04 | 0.07 | 0.13 | 1 | 0 | 1 |  |
| F98.......... | 143612 | +2743 | 11.82 | -0. 21 | 0.00 | 0.07 | 1 | 0 | 1 |  |
| F105........ | 145306 | +1257 | 13.01 | -0. 05 | 0.01 | 0.03 | 1 | 0 | 1 |  |
| +300 $2726 . .$. | 155517 | +2955.1 | 10.03 | 0.64 | 0.89 | 0.74 | 1 | 0 |  | Peterson std. |
| W olf 672A.... | 171605 | +200 | 14. 35 | -0. 57 | 0.11 | 0.11 | 0 | 1 | 2,3 | G19-20 |
| Wolf 672B.... | 171606 | +200 | 14.00 | 1. 15 | 1.53 | 1. 51 | 0 | 1 | 2,3 | G19-21 |
| +300 3075.... | 175045 | +3007.1 | 10.00 | 1.56 | 1.33 | 1.02 | 2 | 0 | ... | Peterson std. |
| +260 3578.... | 193024 | +2615 | 9.31 | -0.17 | 0.42 | 0.38 | 1 | 0 | ... | In SS list |
| L997-21..... | 195400 | +109 | 13.66 | -0. 59 | 0.29 | 0.32 | 0 | 1 | 2 |  |
| +280 3590..... | 200012 | +2831.5 | 10. 10 | -0. 20 | 0.12 | 0.18 | 2 | 0 | . $\cdot$ | Peterson std. |
| +50 $5^{\circ} 4481 \ldots .$. | 201754 | + 553 | 10.11 | -0.03 | 0.61 | 0.58 | 7 | 0 | . . | In SS list |
| W 12754....... | 202236 | +2454 | 10.79 | -0. 25 | 0.40 | 0. 43 | 6 | 0 |  | In SS list |
| W olf 1346..... | 203213 | +2454 | 11.47 | -0.86 | -0.08 | -0. 13 | 1 | 0 | 2,3 | G186-31 |
| LDS 749B..... | 212937 | 000 | 14.70 | -1.00 | -0.02 | 0.10 | 1 | 0 | 2,3 | G26-9 |
| L930-80...... | 214454 | -758 | 14.75 | -0. 99 | -0.07 | 0.03 | 1 | 2 | 2 |  |
| $+^{28} 8^{\circ} 4211 \ldots$. | 214700 | +2827 | 10.53 | -1. 26 | -0.34 | -0. 14 | 1 | 0 |  | S4 UBV values |
| G93-48...... | 214953 | + 209.4 | 12. 71 | -0.75 | -0.02 | -0. 06 | 3 | 2 | 3 |  |
| + $30^{\circ}$ 4564.... | 215459 | +30 15.3 | 10.29 | 0.20 | 0.68 | 0.59 | 1 | 0 |  |  |
| G18-39....... | 221606 | + 811.7 | 10.36 | -0. 20 | 0.48 | 0.51 | 3 | 0 | 3 |  |
| G67-23... | 224639 | +22 20.5 | 14. 34 | -0. 68 | 0.21 | 0. 14 | 1 | 3 | 3 |  |
| -150 6290.... | 225036 | -1431 | 10.16 | +1. 15 | +1.60 | 1.46 | 1 | 0 | ... | S4 UBV values (JM) |
| +380 4955.... | 231118 | +3909 | 11.00 | -0.11 | 0.70 | 0.63 | 3 | 0 | . |  |
| F108.. | 231336 | - 207 | 12.96 | -1.08 | -0. 24 | -0.05 | 0 | 2 | 1 |  |
| G29-38....... | 232616 | $+458.5$ | 13.02 | -0. 57 | 0.17 | 0.08 | 0 | 1 | 3 |  |
| Stars in Clusters or Regions |  |  |  |  |  |  |  |  |  |  |
| 3C48A. | 13448 | +32 54 | 10.95 | 0.14 | 0.58 | 0.53 | 1 | 0 | 4 |  |
| B. | 11 | " | 13.52 | -0. 01 | 0.50 | 0.48 | 1 | 5 | " |  |
|  | " | " | 13. 59 | 0.05 | 0.50 | 0.52 | 1 | 0 | " |  |
| $\mathrm{C}^{1}$. | " | " | 12. 89 | 0.14 | 0.54 | 0.49 | 3 | 4 | " |  |
| D........ | " | " | 14.56 | 0.05 | 0.66 | 0.60 | 3 | 5 | " |  |
| M67-81....... | 84827 | +1200 | 10.04 | -0.42 | -0.06 | 0.02 | 10 | 0 | 5 |  |
| 83...... |  | " | 13. 20 | 0.10 | 0.60 | 0.55 | 3 | 3 | " |  |
| 84....... | "' | " | 10.54 | 0.97 | 1. 10 | 0.82 | 4 | 0 | " |  |
| 105....... | "' | " | 10.31 9.68 | 1. 26 | 1.25 1.37 | 0.95 1.06 | 7 1 | 0 | "' |  |
| 131........ | " | " | 11. 21 | 0.05 | 0.42 | 0.43 | 3 | 0 | " |  |
| 136... | " | " | 11.31 | 0.18 | 0.64 | 0.57 | 4 | 0 | " |  |
| 141....... | " | " | 10.49 | 0.99 | 1.10 | 0.84 | 1 | 0 | " |  |
| 151....... | " | " | 10.55 |  | 1. 10 | 0.84 | 1 | 0 | " |  |
| 156. | " | " | 10.96 | -0.01 | 0.11 | 0.14 | 3 | 0 | " |  |
| 157. | " | " | 12.78 | 0.11 | 0.60 | 0.52 | 2 | 0 | " |  |
| 193.. | " | " | 12.28 | 0.76 | 0.99 | 0.79 | 1 | 0 | " |  |
| 231.... | " | " | 11.49 | 0.87 | 1.06 | 0.78 |  | 0 | " |  |
| M67-III-34.... | " | " | 11.26 | 0.96 | 1.08 | 0.81 | 1 | 1 | 6 |  |
| III-35.. | " | " | 12.14 | 0.78 | 1.01 | 0.78 | 1 | 2 | 6 |  |
| III-53... | " | 2 | 12. 51 | 0.18 | 0.67 | 0.57 |  | 1 | 6 |  |
| 3C273B*..... | 122636 | + 220 | 14.92 | 1.02 | 1. 10 | 0.98 | 4 | 2 | 4 |  |
| E........ | "1" | 11 | 12. 71 | 0.13 | 0.69 | 0.61 | 4 | 1 | " |  |
| G3-III-9.... | 133954 | 11 +2838 | 13.51 | 0. 10 | 0.63 0.47 | 0.55 0.54 | 6 | 2 | 7 |  |
| M3-III-9..... | 133954 | +28 ${ }_{11} 38$ | 15.63 14.74 | 0.100 0.37 | 0.47 0.89 | 0.54 0.77 | 1 | 0 1 | 7 |  |
| 193...... | " | " | 14.74 14.04 | 0.37 0.62 | 0.89 1.01 | 0.77 0.86 | 0 0 | 1 | " |  |
| 297...... | " | " | 12. 89 | 1. 40 | 1.41 | 1. 06 | 1 | 2 | " |  |
| 1397. . . . . | " | " | 12.75 | (1.63) | 1.58 | 1.16 | 4 | 2 | " | $\mathrm{U}-\mathrm{B}$ is 1P21 value |
| 1402..... | " | " | 12. 69 | 0.12 | 0.65 | 0.56 | 4 | 2 | " |  |
| M92-III-13. . . | 171536 | +4312 | 12. 20: | 1.13 | 1. 35 | 1.06 | 0 | 2 | 8 | .Variable? |
| IV-17.... | " | " | 15. 56 | 0.00 | 0.03 | 0.12 | 0 | 2 | " |  |
| IV-114... | " | " | 13.90 | 0.38 | 0.88 | 0.78 | 0 | 2 | " |  |
| NGC 7790-C.. | 235556 | +60 57 | 12. 49 | 0.19 | 0.55 | 0.51 | 0 | 5 | 9 |  |
| E. . | " | " | 12. 75 | -0.11 | 0.40 | 0.44 | 0 |  | " |  |
| J.. | " | " | 13. 23 | 0.34 | 0.53 | 0.47 | 0 | 3 | " |  |
| K . . | " | " | 13.25 | 1.28 | 1.48 | 1. 14 | 0 | 1 | " |  |
| I . | " | " | 13.16 | 0.79 | 1.22 | 0.99 | 0 | 5 | 11 |  |

Chart References to Table 1:
í. Feige, J. 1959.
Luyten, W. J. 1949.
Penston, M. J. et al. 1971
Johns on, H. L. , and Sandage, A. 1955.
Eggen, O. J. , and Sandage, A. 1964.
Sandage, A. 1953; 1970 (Fig. 1).
Sandage, A. and Walker, M. F. 196
Sandage, A. 1958.
Sandage A.
should not be used alone as a secondary standard for $U B V$, because of the few observations per star in the present S20 system.

It is noteworthy that for the galaxies themselves, the S 20 photometry is in systematic agreement to within 0.02 mag of the S4 system of Papers II and III (see § V here, and Paper VI that follows).

## III. THE PHOTOMETRIC DATA

Three photometric programs were carried simultaneously.
a) New data were obtained for the brightest several galaxies in the clusters listed in Paper II and for the added clusters Abell 98, 119, 274, 1213, 2029, and 2670. The last two contain good examples of "cD" galaxies; they are illustrated in Matthews, Morgan, and Schmidt (1964), and are important for the calibration of absolute magnitudes of cluster members as a function of Bautz-Morgan (1970) cluster class (Paper VII that follows).

The photometric data are set out in table 2 . The brighter galaxies were measured generally at Mount Wilson with the Hooker reflector, while the faint galaxies of large redshift were measured at Palomar with the Hale telescope. Pulse counting was used throughout, with data-readout using equipment built by the Astroelectronics Laboratory. The observations were made through a number of aperture sizes whose diameters are listed in the table.
b) Photometry was obtained for the brighter galaxies in all the E and S0 groups listed by HMS (their table 11) to test for a dependence of absolute magnitude of firstranked galaxies on the richness of aggregates. The data, set out in table 3, were obtained largely with the 100 -inch on Mount Wilson in 1967-1968. The galaxies in these HMS groups are all nearby, bright, and very large, requiring large apertures for adequate aperture corrections by the method of Paper I (Sandage 1972a). Because low-speed ( 5 MHz ) pulse amplifiers were then in use, appreciable coin idence corrections were required at the larger apertures, making the subtraction of two large numbers (galaxy plus sky count minus sky count alone) relatively inaccurate. Errors that amount to 0.1 mag in the intensities may be present for apertures larger than $\sim 60^{\prime \prime}$ diameter in table 3, but data for the smaller apertures will be more accurate. The colors will be affected to a lesser degree. Color values for the largest apertures, where the coincidence corrections are particularly uncertain, are omitted from the table.
c) Many of the radio sources measured in Paper III were remeasured in the present S20 program, and the data are listed in table 4. Tabulated in the "Remarks" column is an estimate of the emission line strength $S_{6}$ (defined by Schmidt 1965) as an aid to judge how much the measured colors are affected by line emission. The words yes and $n o$, based on the size of $S_{6}$, or on other knowledge of the spectrum, indicate that the colors were, or were not, used to determine the latitude and redshift effects of §§ IV and $V$.

## Iv. SOLUTION FOR GALACTIC REDDENING: NEW ABSORPTION CORRECTIONS

Cosmological information is contained in the data through the correlation of color and redshift. Of principal interest is the second-order color change due to evolution of the stellar content in the look-back time. This is known to be very small. Extraneous color effects must first be removed before there is hope even of its detection.

Galactic reddening is the primary correction. Recent studies have shown that the reddening at high latitudes is exceptionally small, in contrast to the pronounced latitude effect that is usually inferred from the galaxy counts (Hubble 1934; Shane and Wirtanen 1954, 1967; Noonan 1971). Hubble's half-thickness absorption of $A_{B} \simeq 0.25 \mathrm{mag}$, and Shane and Wirtanen's largest values of $A_{B} \simeq 0.5 \mathrm{mag}$, bracket the values by Noonan, de Vaucouleurs and Malik (1969), and others. But all are much

TABLE 2
S20 PHOTOMETRY OF BRIGHT GALAXIES IN CLUSTERS




TABLE 4
S20 PHOTOMETRY OF RADIO GALAXIES

larger than would be predicted from the observed polar cap reddening from many recent determinations (summarized in Paper II, § III $b$, with references).

In view of the problem, we have made yet another solution for systematic reddening as a function of galactic latitude from the new color data in tables 2, 3, and 4. Mean $B-V, V-R$, and $B-R$ values were obtained by averaging the observed colors of all measured galaxies in each group. The measured colors were corrected to zero redshift using the $K$ values listed in table 5. The corrected colors are, to high order, independent of the redshift ${ }^{1}$ because the $K$ reddenings are well known for small values of the redshift $z \equiv \Delta \lambda / \lambda_{0}$. The colors, so corrected, are correlated with the path length through our Galaxy in figure 1. The dependence on galactic latitude is small but definite, in agreement with the prior conclusion of Peterson (1970) from different data using the same method.

Open circles in figure 1 are radio sources marked yes in table 4. Closed circles are the clusters and groups in tables 2 and 3. Least-squares solutions with their correlation coefficients are

$$
\begin{align*}
B-V & =0.976+0.033(\csc b-1) ; & & |r|=0.56  \tag{1}\\
& \pm 0.006 \pm 0.005 & & \\
V-R & =0.861+0.026(\csc b-1) ; & & |r|=0.55  \tag{2}\\
& \pm 0.005 \pm 0.005 & &
\end{align*}
$$

and

$$
\begin{array}{rlr}
B-R & =1.837+0.059(\csc b-1) ; \quad|r|=0.62  \tag{3}\\
& \pm 0.009 \pm 0.010
\end{array}
$$

where the quoted errors are 1 sigma values.
There is no evidence from these data for a general half-layer absorption anywhere near as strong as suggested by the absorption values at the poles inferred from galaxy counts. Our data force the same dilemma as others before them: either the galaxy counts above $|b|>30^{\circ}$ have little to do with those at lower latitudes (Noonan 1971),

TABLE 5
ADOPTED K CORRECTIONS*

| $z$ | $K(B-V)$ | $K(V-R)$ | $K(B-R)$ | $K_{B}$ | $K_{V}$ | $K_{R}$ |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| 0.00.... | 0.00 | 0.00 | 0.00 | 0.00 | 0.00 | 0.00 |
| 0.04..... | 0.13 | 0.03 | 0.16 | 0.19 | 0.06 | 0.03 |
| 0.08.... | 0.27 | 0.05 | 0.32 | 0. 40 | 0.13 | 0.08 |
| 0.12..... | 0.42 | 0.08 | 0.50 | 0.62 | 0.20 | 0.12 |
| 0.16..... | 0. 55 | 0.11 | 0.66 | 0.82 | 0.27 | 0.16 |
| 0.20..... | 0. 64 | 0.17 | 0.81 | 1.03 | 0.39 | 0. 22 |
| 0.24..... | 0.69 | 0.26 | 0.95 | 1.22 | 0.53 | 0.27 |
| 0.28..... | 0.71 | 0.39 | 1. 10 | 1. 40 | 0. 70 | 0.31 |
| 0.32.... | $\ldots$ | 0.53 | $\cdots$ | $\cdots$ | 0.88 | 0.35 |
| 0.36.... | $\cdots$ | 0.68 | $\ldots$ | $\cdots$ | 1.08 | 0.40 |
| 0.40.... | $\ldots$ | 0.77 | $\ldots$ | $\cdots$ | 1. 25 | 0.48 |
| 0.44.... | $\ldots$ | 0.84 | $\ldots$ | $\cdots$ | 1.41 | 0.57 |
| 0.48..... |  | 0. 88 | $\cdots$ | $\ldots$ | 1. 54 | 0.66 |

*Combination of Whitford and Schild and Oke's values in
Ap. J. 169, 209 and 215, 1971.

[^0]or the ratio of absorption to reddening is very large. The present data shed no light on the alternatives, but the result of Freeman (1973) for the highly obscured galaxy Lyngå 1 show that a high ratio is inadmissible. McClure and Crawford's (1971) model gives an acceptable solution.

The ratio of the latitude coefficients in equations (1) and (2) is 1.27. Although this number has appreciable error due to the high scatter of figure 1, it is the same as the ratio of reddenings for normal galactic dust of $E(B-V) / E(V-R)=1.25$ from Whitford (1958) and Nandy (1964, table 13). This satisfactory agreement leads us to believe that the correlations in figure 1 actually represent true reddening.

In view of these results, we have reconsidered the absorption corrections of Paper II (§ III $b$ ), and adopt the following modified equations to be used in the next several papers of the series:

$$
\begin{array}{ll}
A_{B}=0.132(\csc b-1) & \text { for }|b| \leq 50^{\circ}, \\
A_{B}=0 & \text { for }|b| \geq 50^{\circ}, \tag{4}
\end{array}
$$

but with equation (4) modified slightly between $|b|=40^{\circ}$ and $50^{\circ}$ so that the approach to zero at $50^{\circ}$ is smooth with no discontinuity:

$$
\begin{equation*}
A_{V}=0.75 A_{B}, \quad A_{R}=0.72 A_{V} \tag{5}
\end{equation*}
$$

## V. DETERMINATION OF REDSHIFTS FROM COLORS ALONE

The observed mean colors were corrected for latitude effect by equations (1)-(3) and are listed in columns (4), (6), and (8) of table 6 . Colors in these columns should then depend only on the intrinsic properties of the galaxies themselves, and a function of redshift alone.


Fig. 1.-Correlation of colors, corrected to zero redshift, with path length through the galaxy. Closed circles, clusters and groups from tables 2 and 3. Open circles, radio galaxies in which the effect of emission lines on their colors is small or negligible.

TABLE 6
CORRECTED COLORS FOR CLUSTERS, GROUPS, AND RADIO GALAXIES

| Object <br> (1) | (2) | (3) | $\overline{\langle B-V\rangle}$ <br> (4) | $\begin{gathered} \langle B-V\rangle \\ -E-K \\ (5) \end{gathered}$ | $\begin{gathered} \langle V-R\rangle \\ -E \\ (6) \end{gathered}$ | $\begin{gathered} \langle V-R\rangle \\ -E-K \\ (7) \end{gathered}$ | $\begin{gathered} \langle B-R\rangle \\ -E \\ (8) \end{gathered}$ | $\begin{gathered} \langle B-R\rangle \\ -E-K \\ (9) \end{gathered}$ | Galaxies |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| Clusters |  |  |  |  |  |  |  |  |  |
| Virgo. | +70 | 0.00381 | 1.028 | 1.028 | 0.848 | 0.848 | 1. 876 | 1.876 | N4486 |
| Peg I | -48 | 0.01279 | 1. 066 | 1.036 | 0.894 | 0.884 | 1. 960 | 1.920 | N7619, 7626, 7562 |
| 0122+3305... | -29 | 0.0170 | 1. 020 | 0.970 | 0.847 | 0.837 | 1.867 | 1. 807 | N507, 499, 508 |
| Per (N1265) .. | -13 | 0.0181 | 1.036 | 0.986 | 0.880 | 0. 860 | 1. 916 | 1.846 | N 1265 (poor obs) |
| A1213........ | +69 | 0.0287 | 1.093 | 1.003 | 0.853 | 0.833 | 1. 946 | 1.836 | N1, N 2 |
| A119.. | -64 | 0.0387 | 1. 124 | 1. 004 | 0.880 | 0.850 | 2. 004 | 1. 854 | N 1, N2 |
| Peg II ....... | -48 | 0.0428 | 1. 162 | 1.022 | 0.937 | 0.897 | 2. 099 | 1.919 | N7501, 7503 |
| $2322+1425 \ldots$ | -43 | 0.0439 | 1. 078 | 0.938 | 0.906 | 0.866 | 1. 984 | 1. 804 | N7649, N1, N7 |
| 0106-1536... | -78 | 0.0526 | 1. 149 | 0.979 | 0.896 | 0.856 | 2. 045 | 1. 835 | N1, N2, N10 |
| $1239+1852 \ldots$ | +81 | 0.0718 | 1. 247 | 1. 007 | 0.922 | 0.872 | 2. 169 | 1.879 | N1, N2, N4 |
| $1520+2754 .$. . | +57 | 0.0722 | 1. 150 | 0.910 | 0.944 | 0.894 | 2. 094 | 1. 804 | N2, N3 |
| A2670........ | -69 | 0.0774 | 1. 244 | 0.984 | 0.920 | 0. 860 | 2. 164 | 1. 844 | N 1 , NA |
| A2029. | +50 | 0.0777 | 1. 220 | 0.960 | 0.912 | 0.852 | 2. 132 | 1. 812 | N1 |
| $0705+3506 \ldots$ | +18 | 0.0779 | 1. 176 | 0.916 | 0.872 | 0.812 | 2. 048 | 1. 728 | N1, N2, N1A |
| 1513 + 0433... | +49 | 0.0944 | 1. 239 | 0.919 | 0.897 | 0.837 | 2. 136 | 1. 756 |  |
| A98..... | -42 | 0.1028 | 1. 344 | 0.984 | 0.978 | 0.908 | 2. 322 | 1. 892 | N1, NA, NE |
| A 274. | -64 | 0.1289 | 1. 434 | 0.984 | 0.985 | 0.895 | 2. 419 | 1. 879 | N1, NA, NE |
| $0025+2223 \ldots$ | -40 | 0.1594 | 1.480 | 0.930 | 0.980 | 0.860 | 2. 460 | 1. 790 | N1, N2, N4 |
| $0138+1832 \ldots$ | +61 | 0.1917 | 1. 560 | 0.940 | 1. 106 | 0.946 | 2. 666 | 1.886 | N1, N1A |
| $0855+0327 \ldots$ | +29 | 0.2018 | 1. 644 | 1. 004 | 1. 064 | 0. 894 | 2. 708 | 1. 898 | N1, N2, N2A, N5 |
| $1447+2617$. . | +63 | 0.36 | 1. 486 |  | 1. 597 |  | 3. 083 |  | N4 |
| 3C295.... | +61 | 0. 461 | 1. 415 |  | 1.731 |  | 3. 146 |  | N1 |
| HMS Groups |  |  |  |  |  |  |  |  |  |
| G68. ......... | -32 | 0.0226 | 1.011 | 0.941 | 0.855 | 0.835 | 1.866 | 1.776 | N68, 70, 71 |
| G80.. | -40 | 0.0209 | 1.090 | 1. 030 | 0.876 | 0.856 | 1. 966 | 1.886 | N80, 83 |
| G128. | -60 | 0.0155 | 0.955 | 0.915 | 0.859 | 0.839 | 1. 814 | 1. 754 | N 128 |
| G194. | -60 | 0.0177 | 0.990 | 0.940 | 0.854 | 0.834 | 1. 844 | 1. 774 | N193, 194 |
| G383. | -30 | 0.0175 | 1. 062 | 1. 012 | 0.879 | 0.859 | 1. 941 | 1.871 | N383 |
| G507. | -29 | 0, 0155 | 1. 005 | 0.965 | 0.853 | 0.833 | 1. 858 | 1. 798 | N507, 499, 508 |
| G547. | -63 | 0.0180 | 1. 050 | I. 000 | 0.881 | 0.861 | 1. 931 | 1. 861 | N547, 545, 543, 564 |
| G741. | -54 | 0.0188 | 1. 072 | 1. 022 | 0.854 | 0. 874 | 1. 946 | 1.876 | N741 |
| G1600. | -33 | 0.0160 | 1. 012 | 0.962 | 0.834 | 0.814 | 1. 846 | 1.776 | N 1600 |
| G2563. | +29 | 0.0159 | 1. 007 | 0.957 | 0.872 | 0.852 | 1. 879 | 1. 809 | N2562, 2563 |
| G2832. | +44 | 0.0200 | 1.022 | 0.962 | 0.859 | 0.839 | 1. 881 | 1. 801 | N2832, 2831 |
| G3158. | +55 | 0.0234 | 1. 019 | 0. 949 | 0.885 | 0.865 | 1. 904 | 1. 814 | N3158, 3159, 3163 |
| G3193. | +55 | 0.0040 | 0.973 | 0.973 | 0.849 | 0.849 | 1. 822 | 1.822 | N3193 |
| Leo.. | +64 | 0.0026 | 1. 030 | 1.030 | 0.889 | 0.889 | 1. 919 | 1.919 | N3368, 3379 |
| G5044. | +46 | 0.0087 | 1.021 | 0.991 | 0.880 | 0.870 | 1. 901 | 1. 861 | N 5044 |
| G5077. | +50 | 0.0084 | 1. 034 | 1. 004 | 0.858 | 0.848 | 1. 892 | 1.852 | N5077 |
| G5371. | +71 | 0.0076 | 1. 007 | 0. 987 | 0.866 | 0.856 | 1.873 | 1. 843 | N5353, 5354 |
| G5846. | +49 | 0.0060 | 1. 024 | 1. 004 | 0.858 | 0.848 | 1. 882 | 1.852 | N5846, 5813, 5831 |
| G7242. | -16 | 0.0204 | 1. 070 | 1. 000 | 0.868 | 0.848 | 1. 938 | 1. 848 | N7242 |
| G7385. | -41 | 0.0258 | 1. 103 | 1. 023 | 0.916 | 0.886 | 2. 019 | 1.909 | N7385, 7386 |
| G7619. | -48 | 0.0128 | 1. 058 | 1. 018 | 0.907 | 0.897 | 1.965 | 1.915 | N7619, 7626, 7562 |
| vV172 | +45 | 0.0526 | 1. 136 | 0. 966 | 0.931 | 0.891 | 2. 067 | 1. 857 | A, C, D, E |
| Radio Sources With $\mathrm{S}_{6} \leq 4$ |  |  |  |  |  |  |  |  |  |
|  |  |  |  |  |  |  | $S_{6}$ |  |  |
| 3 C 15. | -64 | 0.0733 | 1. 146 | 0.896 | 0.887 | 0.827 | 2. 033 | 1. 723 | $\sim 1$ |
| 3 C 29. | -64 | 0.0450 | 1. 059 | 0.919 | 0. 890 | 0.850 | 1.949 | 1. 769 | 0 |
| 3 C 76.1. | -36 | 0.0328 | 1. 082 | 0.982 | 0.912 | 0.882 | 1. 994 | 1. 864 | 0 |
| 3 C 78. | -45 | 0.0289 | 1. 134 | 1. 044 | 0.973 | 0.953 | 2. 107 | 1. 997 | 1 : |
| 3 C 88. | -42 | 0.0302 | 1. 104 | 1. 004 | 0.938 | 0.908 | 2. 042 | 1. 912 | 2 |
| 3C264. | +73 | 0.0206 | 1. 023 | 0.963 | 0.879 | 0.859 | 1. 902 | 1. 822 | 0 |
| 3 C 270. | +67 | 0.0070 | 0.990 | 0.970 | 0.838 | 0.828 | 1. 828 | 1. 798 | 0 |
| 3 C 272.1. | +74 | 0.0029 | 1. 022 | 1. 012 | 0.842 | 0.842 | 1. 864 | 1. 854 | < 3 |
| $3 \mathrm{C278}$. | +50 | 0.0143 | 1. 040 | 1. 000 | 0. 884 | 0.874 | 1. 924 | 1.874 | 0 |
| 3 C 293. | +76 | 0.0454 | 0.990 | 0.850 | 0.879 | 0.839 | 1. 869 | 1. 689 | $\sim 2$ |
| 3C296. | +62 | 0.0237 | 1. 031 | 0.961 | 0.857 | 0.837 | 1. 888 | 1. 798 | 0 |
| 3C317. | +50 | 0.0351 | 1. 095 | 0.985 | 0.912 | 0.872 | 2. 007 | 1. 857 | 3 |
| 3C318.1...... | +49 | 0.0457 | 1. 121 | 0.981 | 0.942 | 0.902 | 2. 063 | 1. 833 | 0 |
| 3C338. | +44 | 0.0303 | 1. 046 | 0.946 | 0.832 | 0.802 | 1. 878 | 1. 748 | 0 |
| 3C388. | +20 | 0.0917 | 1. 335 | 1. 024 | 0.925 | 0.865 | 2. 260 | 1. 889 | 1 |
| 3C402. | +13 | 0.0259 | 1. 042 | 0.962 | 0.925 | 0.905 | 1. 967 | 1. 867 | 0 |
| 3C442. | -34 | 0.0270 | 1. 038 | 0.958 | 0.832 | 0.802 | 1. 870 | 1. 760 | $\sim 1$ |
| 3C449. | -16 | 0.0181 | 1.010 | 0.960 | 0.855 | 0.835 | 1. 865 | 1. 795 | 0 |
| 3C455........ | -41 | 0.0331 | 1. 040 | 0. 940 | 0.853 | 0.823 | 1. 893 | 1. 763 | 0 |
| 3C465. | -33 | 0.0301 | 1. 060 | 0.960 | 0.914 | 0.894 | 1. 974 | 1. 854 | , |
| Radio Sources With $\mathrm{S}_{6}>4$ |  |  |  |  |  |  |  |  |  |
| 3C17......... | -65 | 0.2201 | 0.752 | 0.082 | 0.852 | 0.642 | 1. 604 | 0.724 | 7 |
| 3C26.......... | -67 | 0.2106 | 1. 137 | 0.487 | 0.858 | 0.668 | 1. 995 | 1. 155 | 11 |
| 3C28......... | -37 | 0. 1959 | 1. 390 | 0.760 | 1. 128 | 0.968 | 2. 518 | 1. 728 | 2 |
| 3C33......... | -49 | 0.0600 | 1. 152 | 0.952 | 0.835 | 0.785 | 1. 987 | 1. 737 | 12 |
| 3C66......... | -17 | 0.0215 | 0.943 | 0.883 | 0.832 | 0.812 | 1. 775 | 1. 695 | ? |
| 3C75......... | -45 | 0.0241 | 1. 213 | 1. 133 | 1. 022 | 1.002 | 2. 235 | 2. 135 | ? |
| 3C98......... | -31 | 0.0306 | 1. 171 | 1.071 | 0.948 | 0.918 | 2. 119 | 1. 989 | 8 |
| 3 C 192. | +26 | 0.0596 | 1. 163 | 0.963 | 0.832 | 0.782 | 1. 995 | 1. 745 | $\sim 5$ ? |
| 3C198. | +23 | 0.0809 | 0.849 | 0.578 | 0.749 | 0.689 | 1. 598 | 1. 267 | 7 |
| 3C219........ | +45 | 0. 1745 | 1. 469 | 0. 889 | 0.929 | 0.799 | 2. 398 | 1. 688 | 8 |
| 3C223.1...... | +49 | 0.1075 | 1. 279 | 0. 909 | 0.912 | 0.832 | 2. 191 | 1. 741 | $\sim 6$ : |
| 3C381........ | +23 | 0.1614 | 1. 369 | 0.818 | 0.884 | 0.764 | 2. 253 | 1. 582 | $>8$ |
| 3C382........ | +17 | 0.0586 | 1. 063 | 0.873 | 0.867 | 0.817 | 1. 930 | 1. 690 | 5 |
| 3C436........ | -19 | 0.2154 | 1. 382 | 0.722 | 1.116 | 0.916 | 2. 498 | 1. 638 | 5 |
| 3C452. | -17 | 0.0820 | 1. 215 | 0.935 | 0. 862 | 0.802 | 2. 077 | 1. 737 | 5 |
| 3C456........ | -46 | 0.2337 | 0.777 | 0.097 | 1. 125 | 0.885 | 1. 902 | 0.982 | 7: |
| MSH 23-112.. | -64 | 0.0825 | 1. 049 | 0.769 | 0.934 | 0.874 | 1. 983 | 1. 643 | 7 |



Fig. 2.-Correlation of the $B-V$ color, corrected for galactic reddening by eq. (1), with redshift for groups and clusters from tables 2 and 3. Closed circles, the present S20 photometry; open circles, independent S4 photometry reported in Paper II. The line is the calculated relation obtained by shifting the standard energy distribution of giant E galaxies through the measuring bands.

The color-redshift effect in $B-V$ is shown in figure 2 . Closed circles are the present data, open circles are the similarly re-reduced colors from Paper II (table 1) discussed by Oke and Sandage (1968), but there without latitude correction.

A small difference exists between the open and closed circles for $z \geqslant 0.16$ : the S4 colors are slightly bluer than the S20 values. Comparison of the overlap between the S20 data in tables 2, 3, and 4 here with the 1P21 data of Papers II and III give $\langle\Delta(B-V)\rangle=0.016 \pm 0.006(\mathrm{rms})$ in the sense of the S 4 photometry minus this paper. The sigma of the distribution of color differences between the two sets is $\sigma=0.038$ mag. If equal errors are present in both investigations, each will have a distribution of internal errors of $\sigma E(B-V) \simeq 0.027 \mathrm{mag} .{ }^{2}$ The source of the slight systematic difference between the open and the closed dots at $z \equiv \Delta \lambda / \lambda_{0} \simeq 0.16$ in figure 2 is not known.

The solid line in figure 2 is the predicted run of color with redshift if all E and S 0 galaxies have an identical spectrum equal to the mean of the standard galaxy of Whitford (1971) and Schild and Oke (1971). These calculated $K(B-V)$ values are listed in table 5. The curve has been normalized to $(B-V)_{0}=0.97$ at $z=0$, which is that calculated by Schild and Oke from their energy distribution, using the precepts of Matthews and Sandage (1963). The earlier calculation by Oke and Sandage (1968) gave $B-V=1.04$ for the centralmost region of M31 alone ( $\theta \simeq 14 \mathrm{arcsec}$ diameter). This was redder than the present value by 0.07 mag due to the radial color gradient, which explains entirely the color shift of 0.07 mag required by these authors to make their calculations coincide with the present observations in zero point.

It is evident from figure 2 that the $B-V$ colors become bluer again as the redshift becomes larger than $z \simeq 0.3$. The behavior is expected because the ultraviolet energy

[^1]distribution flattens shortward of $\lambda_{0} \simeq 3600 \AA$ relative to its precipitous drop in the range $3600 \AA<\lambda<4200 \AA$. The $B$ and $V$ filters fall in this region of changing gradient in the proper galaxy spectrum if the redshift is larger than $z \simeq 0.3$. The bluing occurs in $B-V$ at this redshift for the same reason that $U-B$ becomes bluer for increasing redshift in the interval $0<z \leqslant 0.04$ (Sandage 1972b, table 3).

Figure 3 shows the color-redshift effect in $V-R$, again compared with the predicted variation. The agreement is good. The zero-point of the line has been put at $(V-R)_{0}$ $=0.86$, which is the empirical value for $z=0$. The main point to note in figure 3 is the shallow gradient for redshifts smaller than $\sim 0.16$, and the rapid steepening thereafter. Note also that the shape of the curve for $z \geqslant 0.22$ is similar to that in figure 2 for $z \geq 0$ as expected because the proper wavelengths of the $B$ and $V$ filters at $z=0$ are nearly the same as those for the $V$ and $R$ filters for $z \simeq 0.2$.

An important practical application of these color-redshift effects is evident from figure 3. The variation of color with $z$ is well enough defined, and is steep enough in the interesting region beyond $z \geqslant 0.2$ that moderately accurate redshifts can be determined from the $V-R$ broad-band colors alone. From the horizontal scatter in this diagram, and from figure 4 for $B-R$, the error for such measurements will be of order $\sigma(\Delta z) \simeq 0.01$, which is highly satisfactory for many purposes. The technique is the same as used by Baum (1962) except that it can evidently be simplified to two, or at most three, colors, for $z \leqslant 0.5$.

The information in figures 2-4 allows us to optimize a broad-band system for the problem. The diagrams show that (1) $\Delta(B-V)$ is most sensitive to redshift in the low- $z$ range $0<z \leqslant 0.2$; it becomes insensitive and multivalued for larger $z$. (2) $\Delta(V-R)$ is sensitive to redshift in the range $0.2 \leq z \leqslant 0.6$; it becomes insensitive and multivalued for larger $z$ for the same reason as $B-V$. (3) For redshifts larger than $z=0.6$, and in the absence of strong color-evolution in the look-back time, a band near $\langle\lambda\rangle=8500 \AA$, combined with $R$ will reproduce the crucial steep part of the curves in figures 2 and 3 in the redshift range $0.9>z \geqslant 0.6$.

Redshifts by color can therefore be obtained in the range $0.2<z \leqslant 0.9$ by combining measurements made at $\langle\lambda\rangle \simeq 5500 \AA, 6700 \AA$, and $8500 \AA$ only. To cover the entire range $0<z \leqslant 0.9$ requires only the addition of the $B$ filter near $\langle\lambda\rangle \simeq 4400 \AA$.


Fig. 3.-Same as fig. 2 for reddening corrected $V-R$ colors


Fig. 4.-Same as fig. 2 for reddening corrected $B-R$ colors

These color properties of E and S0 galaxies are the basis upon which a redshift program for faint clusters has been started by Kristian, Westphal, and the writer using a silicon vidicon area photometer similar to that developed by McCord and Westphal (1972).

The similarity of the energy distributions of the galaxies in the sample is shown in figure 5, which is the histogram of measured colors corrected for galactic and $K$ reddening. The rms spread is $\sigma \simeq 0.03 \mathrm{mag}$ for $V-R$, and $\sigma \simeq 0.04 \mathrm{mag}$ for $B-V$. The external errors in the photometry, estimated earlier in this section, are distributed with $\sigma_{E} \simeq 0.027 \mathrm{mag}$, giving a true cosmic dispersion in $B-V$ of $\sigma(B-V) \simeq$ 0.029 mag. The homogeneity of the colors for giant E and S0 galaxies implied by this small spread confirms earlier work of Stebbins and Whitford $(1937,1952)$ where the similarity was first demonstrated, by Lasker (1966), Oke and Sandage (1968), and others.


Fig. 5.-Distribution of observed colors for the galaxies in clusters, groups, and radio sources with $S_{6} \leq 4$ that are listed in table 6 .

## VI. LIMITS ON COLOR EVOLUTION IN THE LOOK-BACK TIME

The generally good agreement of the predictions with the observations in figures 2, 3 , and 4 precludes any gross color evolution in the look-back time. The same conclusion follows from the recent observation of 3C 295 with higher spectral resolution (Oke 1971). But some color evolution must, of course, occur, and it is of interest to put observational limits on the effect, and to estimate reasonable expectation values from current knowledge of stellar evolution.

Empirical data are available on the correlation of integrated colors and age for simple stellar aggregates such as star clusters, where the cluster members are nearly coeval (Sandage 1963, fig. 7; Gray 1965, fig. 4). In the age range that is relevant here, $8<\log t \leqslant 10$ where $t$ is in years, the color rate from both quoted studies is $\Delta(B-V)$ / $\Delta \log t \simeq 0.3 \mathrm{mag}$ per factor of 10 in age, i.e., $0.3 \mathrm{mag} \mathrm{dex}^{-1}$. The stated time interval covers the range of main-sequence termination points from about 1 mag brighter than M11 at $M_{V} \simeq-2$ to NGC 188 at $M_{V} \simeq+4$ (see, for example, the composite C-M diagram of Sandage and Eggen 1969, fig. 6). The colors of the termination point itself vary between $(B-V)_{0}=-0.1$ and $(B-V)_{0}=+0.6$ over the stated time interval, and the colors of the integrated aggregates change from +0.2 to +0.8 , according to the color-age correlations previously cited.

But it is expected that E and S 0 galaxies will show smaller color variations because more light is locked in the unevolved faint main sequence. However, the variation clearly cannot be zero because the main-sequence termination point will have changed by 6 magnitudes in $V$, and 0.6 mag in $B-V$, for the factor of 100 in age. In the following we carry through two calculations based on color-evolutionary rates of $0.3 \mathrm{mag}_{\mathrm{dex}}{ }^{-1}$ and $0.1 \mathrm{mag} \mathrm{dex}^{-1}$, which seem likely to bracket the true range.

Consider first the bulk of the points in figures 2 and 3 . What is the expected color change at $z=0.2$ for the two rates? Let the age of the oldest galaxies be $13 \times 10^{9}$ years. For $q_{0}=0$, the look-back time is $t_{\mathrm{LB}}=[z /(1+z)] H_{0}{ }^{-1}$ which, for $z=0.2$ and $H_{0}=55$ (Sandage and Tammann 1973) is $t_{\mathrm{LB}}=3 \times 10^{9}$ years. The time is essentially the same at $2.7 \times 10^{9}$ years if $q_{0}=+1$ (Sandage 1961, table 3). Therefore, the ratio of the age of galaxies observed at a redshift of 0.2 to that of galaxies seen nearby is $(13-3) / 13$ or $\Delta \log t=0.11$. Hence, the expected color change in proper wavelengths will be $\Delta(B-V)=0.03 \mathrm{mag}$ for the rate of $0.3 \mathrm{mag} \mathrm{dex}^{-1}$, and $\Delta(B-V)=0.01$ mag for the smaller rate of $0.1 \mathrm{mag}_{\mathrm{dex}}{ }^{-1}$. Both predictions are compatible with the observations, which show no systematic color difference among galaxies as a function of redshift at the 0.05 mag level ${ }^{3}$ for $0 \leq z \leq 0.2$.

For 3C 295 at $z=0.461$, the look-back time for $q_{0}=0$ is $5.6 \times 10^{9}$ years, or the age is $(13-5.6) \times 10^{9}=7.4 \times 10^{9}$ years, or $\Delta \log t=0.24$. In this time interval the expected color changes are $\Delta(B-V)=0.07 \mathrm{mag}$ for the $0.3 \mathrm{mag} \mathrm{dex}^{-1}$ rate, and 0.02 mag for the slower rate. Oke's (1971) observations may just preclude the 0.07 mag rate, but do not at all eliminate the smaller rate.

Constraints put by the present observations on the color evolution are therefore not severe and do not yet positively eliminate even the fast rates predicted from open clusters. The estimated rates are, however, fast enough so that they can be seriously looked for observationally at $z>0.4$, and detailed scanner observations are clearly the way to do it.

The main conclusion from this section is that color-evolutionary rates as high as $\Delta(B-V) / \Delta \log t=0.2 \mathrm{mag} \mathrm{dex}^{-1}$ would be undetected for $z \leq 0.4$ in any data now available. Very slow rates such as $\Delta(B-V) / \Delta \log t \leqslant 0.05 \mathrm{mag} \mathrm{dex}^{-1}$ appear to be unrealistic from considerations of the stellar content of E galaxies in terms of a composite C-M diagram.

[^2]
## VII. THEORETICAL COLOR DISTRIBUTIONS FOR VARIOUS

 WORLD MODELS COMPARED WITH OBSERVATIONSa) Steady-State Prediction

On a steady-state hypothesis, all structures must be renewed continuously (statistically uniformly) in time, and any sufficiently large region of space will contain galaxies of all ages. However, the probability of finding an extremely old galaxy is small because the expansion will remove most such systems out of the volume. McCrea (1950) has given the distribution of ages such that $f(t) d t$ is the number of galaxies whose ages are between $t$ and $t+d t$, by considering a uniform production rate and removal by expansion. With assumptions on the color change with time, $f(t)$ can be transformed into the distribution $g(B-V)$ per unit $B-V$ interval.

McCrea shows that.

$$
\begin{equation*}
f(t) d t=3 n H \exp (-3 H t) d t \tag{6}
\end{equation*}
$$

where $H$ is the expansion rate and $n$ is the number of galaxies in the sample. Over the time interval discussed in the last section, the color change with time will be closely approximated by

$$
\begin{equation*}
B-V=a \log t+b \tag{7}
\end{equation*}
$$

which gives $t$ for any $B-V$ once $a$ and $b$ are known.
Equation (6) can be transformed into the distribution of colors by differentiating equation (7) and substituting into equation (6) giving

$$
\begin{equation*}
g(B-V) \Delta(B-V)=6.91 n a^{-1} H t \exp (-3 H t) \Delta(B-V), \tag{8}
\end{equation*}
$$

which, parametrically with equation (7), is a function of $B-V$ alone.
Is equation (8) applicable to our present sample of cluster galaxies? It might be argued that clusters form in the steady state only after the galaxies destined to become members are already older than a certain age, the process of clustering itself taking an appreciable time. That this is not the case is shown directly by the observations of cluster and field E galaxies themselves (McClure and van den Bergh 1968) in which the color properties of both samples are the same. Further evidence against such an argument is the similarity of colors of E galaxies in very dense clusters and in sparse groups such as G194, G5077, and G5353 in our sample, where the colors are significantly the same as for first-ranked members of the very rich clusters. This similarity of colors argues against a detectable age difference in groups of different density. Hence groups and clusters probably do not form by the slow process of accumulation, where the formation time itself would depend on the density of the aggregate.

The conclusion is that our sample of brightest galaxies in groups and clusters should not be biased toward old systems merely because they are in aggregates. Rather, the direct observations of McClure and van den Bergh show that cluster members are representative of field ellipticals.

Our observed distribution is compared with equation (8) in figure 6 for three colorevolution rates, using the color-age relations

$$
\begin{align*}
& B-V=0.3 \log t-2  \tag{9}\\
& B-V=0.1 \log t \tag{10}
\end{align*}
$$

and

$$
\begin{equation*}
B-V=0.05 \log t+0.53 \tag{11}
\end{equation*}
$$



Fig. 6.-Comparison of the observed color distribution, corrected for error (top panel), with predictions for the steady-state model (eq. [8]) for three-color decay rates. The rate in panel (b) at $\Delta(B-V) / \Delta \log t=0.05 \mathrm{mag} \mathrm{dex}^{-1}$ is unrealistically small.

These have been normalized to give $B-V \simeq 1.04$ at $t=13 \times 10^{9}$ years, in reasonable agreement with what we know of colors and ages of E galaxies deduced from data in the Local Group.

We have normalized equation (8) to $n=62$ to compare with the observed sample in figure 5. Shown in figure 6 are the three theoretical distributions of equation (8) with equations (9)-(11), using $H_{0}{ }^{-1}=17.7 \times 10^{9}$ years (Sandage and Tammann 1973). The observed distribution of figure 5, corrected for the observational errors of $\sigma_{E}=0.028 \mathrm{mag}$ from $\S \mathrm{V}$, is shown in the top panel.

The predicted distributions for the two color rates of 0.1 and 0.3 are in clear disagreement with the observations. Too many blue (young) galaxies are predicted, and the shape of the distribution is strongly non-Gaussian. The distribution for the slow rate (panel $b$ ) is approximately as narrow as the observations, but the required rate is too slow to be realistic. The color distributions predicted by a steady-state model appear then to disagree with observation by a wide margin.

## b) Continuous Galaxy Formation with an Upper Age Limit

Another simple model is continuous galaxy formation within a Friedmann universe, with an upper limit to the age of the oldest galaxy. If the formation rate is constani per unit time, then, by the same argument that transformed equation (6) to equation (8), it follows that the expected color distribution is

$$
\begin{equation*}
g(B-V) \Delta(B-V)=2.303 n t a^{-1} T_{0}^{-1} \Delta(B-V), \tag{12}
\end{equation*}
$$



Fig. 7.-Same as fig. 6 but for a model with a uniform E galaxy formation rate with time, and with an upper limit to the age of the oldest galaxy at $13 \times 10^{9}$ years.
where $T_{0}$ is the maximum age, and $a$ and $t$ are given by equation (7), or the explicit forms of equations (9) and (10). Equation (12) is a rising exponential in $B-V$, as seen by eliminating $t$ using equation (7).

The two predicted distributions are shown in figure 7 using equations (9) and (10) with $T_{0}=13 \times 10^{9}$ years and a lower age limit of $10^{8}$ years. Comparison with the observations given in the top panel shows that neither prediction agrees well with observation. From this we would conclude that the formation of giant E and S0 galaxies took place during a shorter time interval than $9<\log t<10$. This is the main result of the present paper. It agrees with the formation history that is required by the distribution of flattenings of elliptical galaxies (Sandage, Freeman, and Stokes 1970), where such galaxies must presumably form during the short time interval of about a free-fall time (generally less than $10^{9}$ years) by collapse from the pregalactic medium.

## c) Bias

The preceding arguments depend on two assumptions:

1. Ages of $E$ and $S 0$ galaxies in clusters and groups are representative of the $E$ galaxy population as a whole (§ VII $a$ ).
2. There has been no bias that has caused us to observe galaxies in a restricted age range only, rather than in the total relevant range. The only criterion used to choose the galaxies in tables 2 and 3 was that the aggregates must contain mostly E galaxies, not spirals. Bias would be present due to this section only if ellipticals were of systematically different age than spirals, such as would obtain if, say, ellipticals evolved from spirals, causing the early end of the age sequence to be eliminated artificially from the sample. However, that spirals and ellipticals are not connected evolutionarily is shown by the strong dichotomy in the flattenings (Hubble 1926; de Vaucouleurs 1959;

Sandage et al. 1970). And because color was not used as a selection criterion for tables $2-4$, there should be no age bias.

The conclusions from figures 6 and 7 are then (1) the absence of blue giant ellipticals is real; (2) no young (i.e., $t \leqslant 10^{9}$ years) giant E galaxies exist in our sample. ${ }^{4}$ Whether they exist in any sample requires more extensive data, but figures 6 and 7 show that such galaxies do not exist in the numbers required, either by the steady-state or by the continuous-production models. Our present data are consistent with the view that E galaxies are the same age to within close limits, that they formed at some definite time in the past, and that no substantial fraction of the class, if any, has formed more recently. It will be possible to put closer limits on the interval of galaxy birth when color observations of giant E galaxies are available that are accurate to 1 percent or better.

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[^0]:    ${ }^{1}$ Deviations from table 5 due to color evolution in the look-back time will not affect the problem discussed here because most of the galaxies that enter the latitude dependence have redshifts smaller than $z \simeq 0.16$, where the time variation of color is clearly negligible by the results of later sections.

[^1]:    ${ }^{2}$ This, then, is not particularly good photometry judged against what can be routinely achieved for stars. The large dispersion is undoubtedly due to the addition of the color data from Mount Wilson, where the sky is not dark. The uncertainty is due to the large coincidence correction as previously discussed. In view of the importance of the color distribution discussed in § VII, the present data should be repeated using the modern wide-bandwidth ( 33 MHz ) preamplifiers, now in use, that reduce the dead-time correction to negligible values.

[^2]:    ${ }^{3}$ It should be noted that the observations are not made at the proper wavelengths (in the rest frame of the galaxies). However, figures 2 and 3 together, which for various $z$-values bracket $B-V$ (proper), show that no color evolution in $B-V$ (proper) exists larger than the stated limits.

[^3]:    ${ }^{4}$ To be sure, blue E galaxies do exist, provided that they are of faint absolute magnitude, but they are blue for reasons other than age (Baum 1959; Sandage 1972b).

